

*Surface-integral formalism for transfer reactions:  
Applications and applicability for (d,p) reactions*

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TORUS Collaboration



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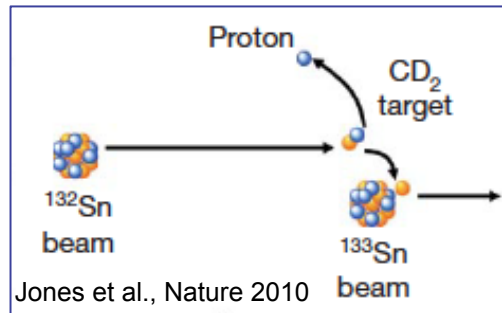
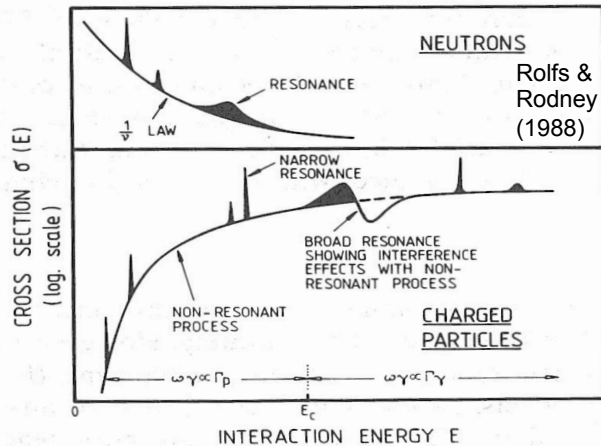
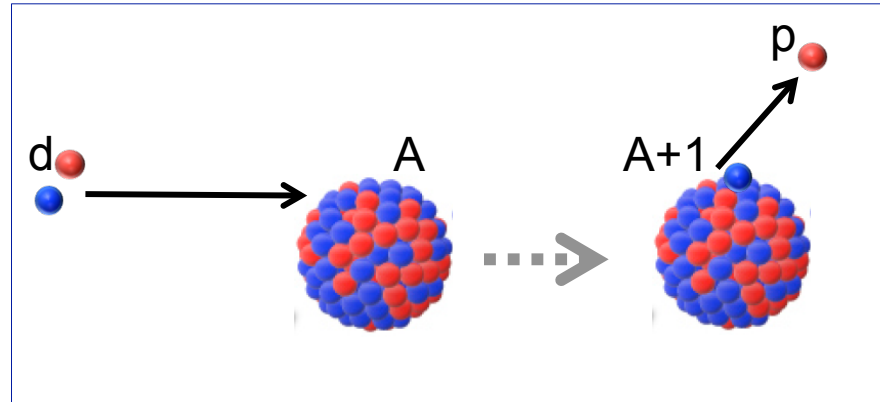
# Studying nuclear structure with (d,p) one-nucleon transfers

(d,p) reactions:

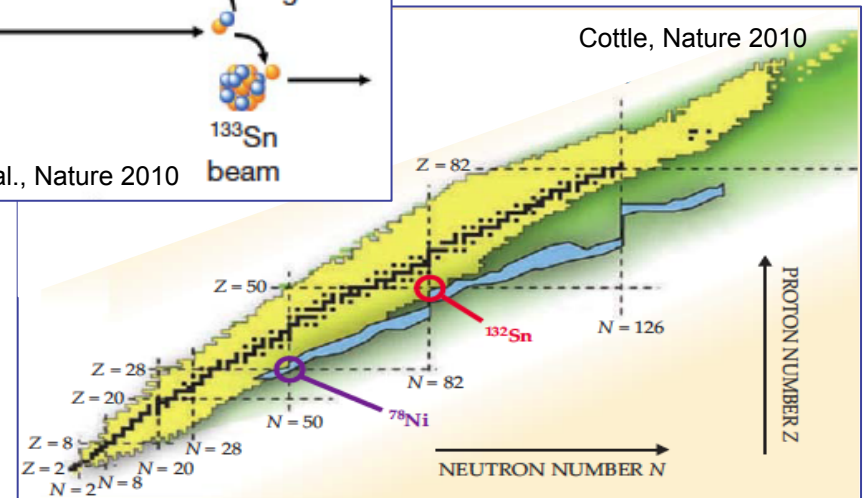
- Simplest mechanism for adding a neutron
- Traditionally used to study stable nuclei
- Used in inverse kinematics at RIB facilities, for studying weakly-bound systems

Theoretical descriptions of (d,p) reactions:

- Progress over the years: Plane-wave theory, DWBA (zero-range & finite-range), coupled-channels approach, breakup, etc.



Jones et al., Nature 2010



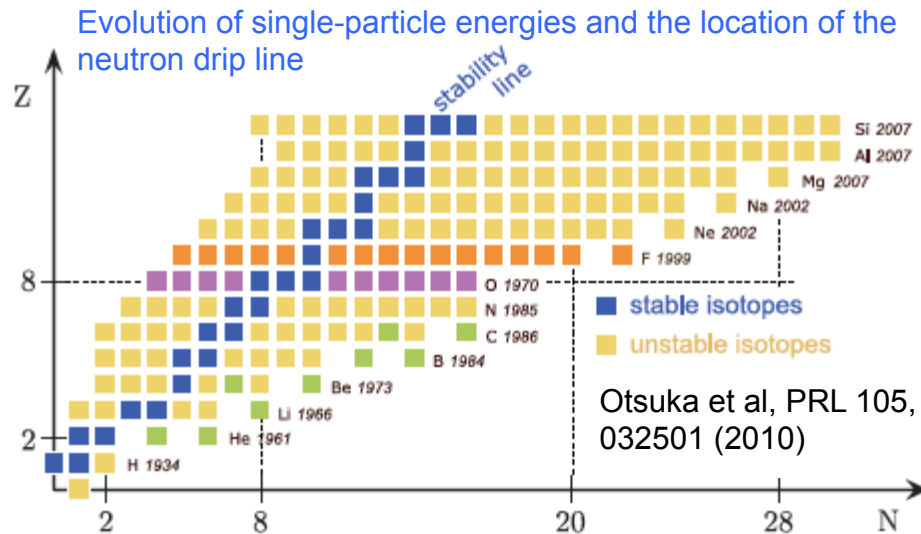
**But:** Current theories of (d,p) reactions not very useful for transfers to resonance states:

- Conceptual: extracting spectroscopic information
- Practical: convergence issues

# Resonances in low-energy nuclear physics

Resonances:

- Unstable quantum-mechanical states
- Occur in light, medium-mass, and heavy nuclei
- Crucially affect astrophysical reaction rates
- Abundant in weakly-bound nuclei

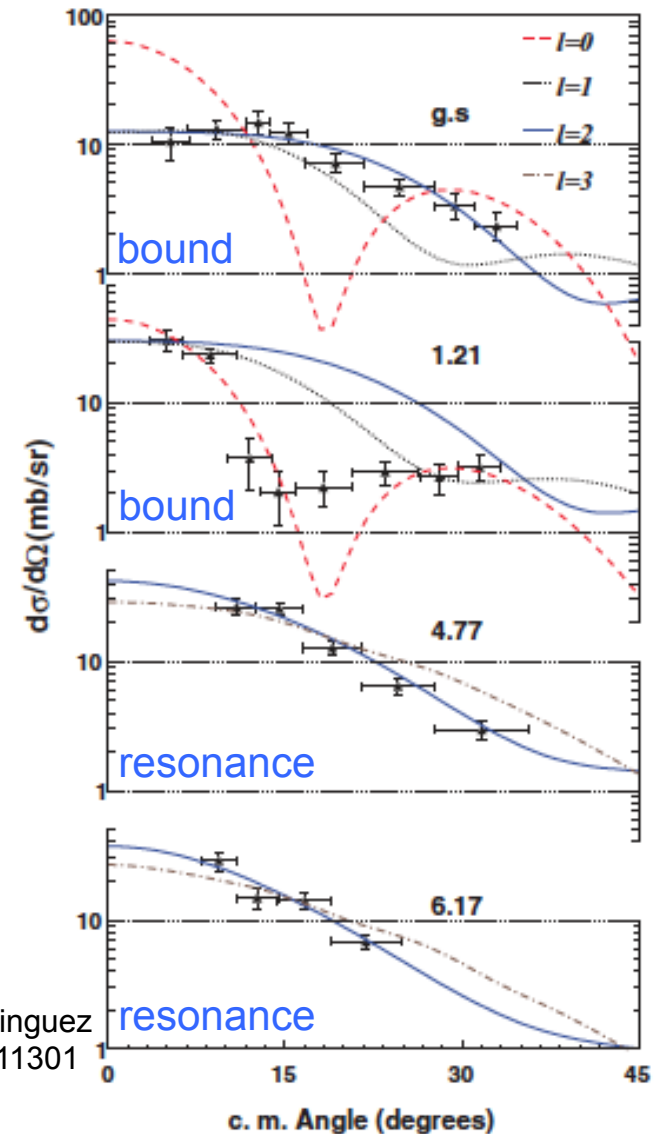


Current approach:

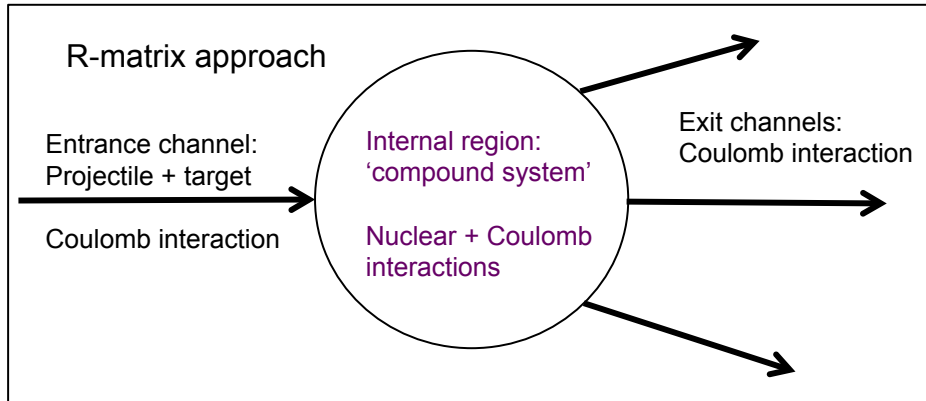
- Apply standard (d,p) descriptions to resonances
- Increase model space to achieve convergence

Fernandez-Dominguez et al, PRC 84, 011301 (R) (2011)

$^{20}\text{O}(d,p)^{21}\text{O}$  inverse-kinematics experiment at GANIL to determine N=16 shell gap



# Describing resonances in binary reactions

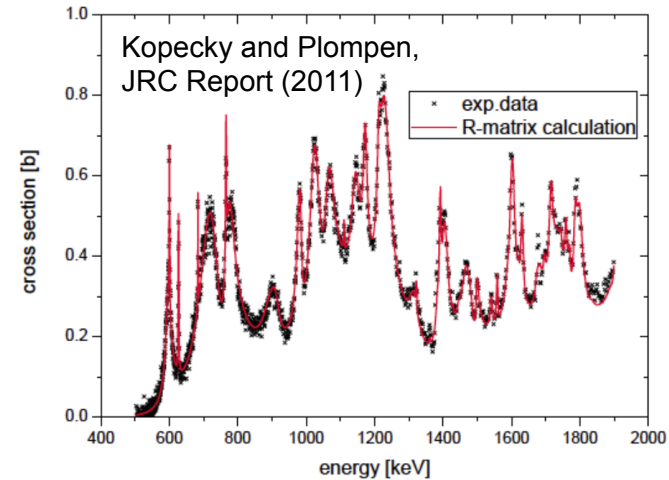


## R-matrix approach:

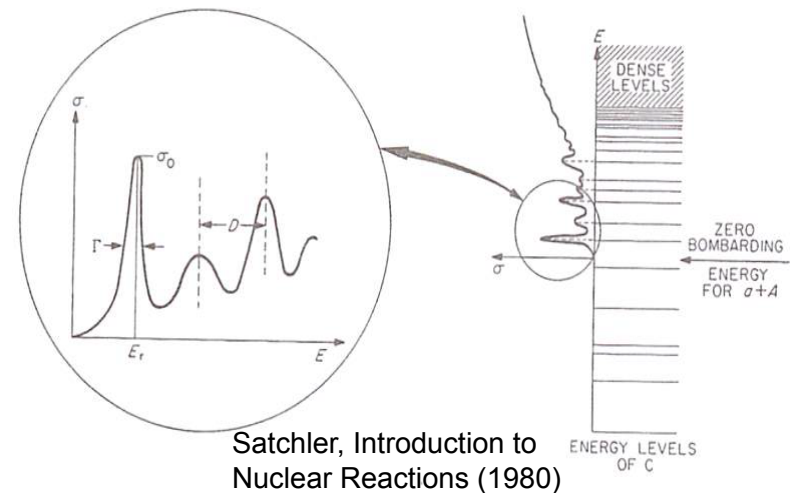
- Main idea: divide space into 2 regions:
  - $r \leq a$  -- interior: nuclear and Coulomb interactions
  - $r > a$  -- exterior: Coulomb only
- Formalism:
  - Interior: expand nuclear wave function in set of basis functions
  - Exterior: scattering wave function
  - Surface: matching conditions allow to parameterize collision matrix  $\rightarrow$  expressions for cross sections
- Connect observed parameters ( $E_R, \Gamma$ ) to formal parameters ( $\check{E}_R, \gamma^2$ )
- Typical applications adjust parameters to reproduce measured cross sections

## Experimental studies of resonances:

- Elastic & inelastic scattering, capture, etc.



- Characterization of resonances: position & widths



Satchler, Introduction to Nuclear Reactions (1980)

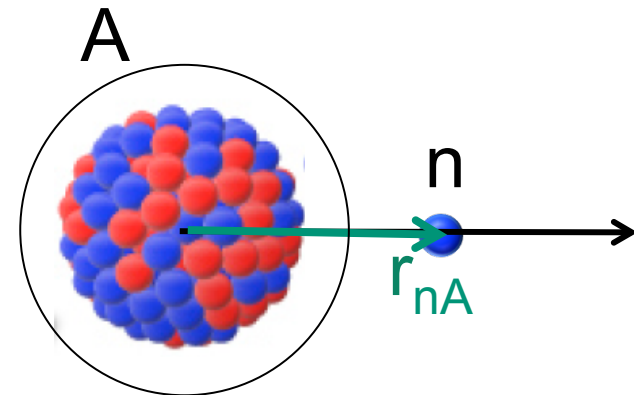
# Exploring R-matrix ideas for (d,p) one-nucleon transfers

## Proposed new formalism (Mukhamedzhanov, 2011):

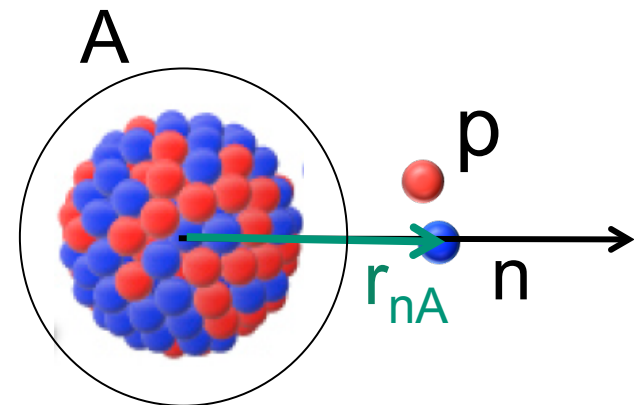
- R-matrix concepts:
  - surface separating internal and external regions
  - cross sections expressed in terms of reduced widths, logarithmic derivatives, surface radii
- Goals for (d,p):
  - useful for resonances
  - reduce dependence on model for interior
  - extract useful spectroscopic quantities from comparison to experiment (widths)
- Formalism:
  - applicable to stripping to bound and resonance states
  - general enough to include deuteron breakup contributions via CDCC (continuum-discretized coupled-channels method)
  - bonus: resolves practical issues related to numerical convergence

Formalism:  
Mukhamedzhanov, PRC 84, 044616 (2011)

## Binary system



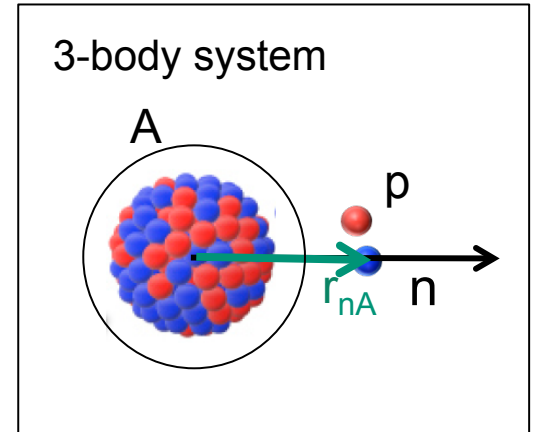
## 3-body system



# Exploring R-matrix ideas for (d,p) one-nucleon transfers II

Transition matrix element M:

- Connects initial to final wave function
- Cross section  $\sigma \sim M^2$



$$M^{(\text{post})} = \langle \Phi_f^{(-)} | \Delta V_{pF} | \Psi_i^{(+)} \rangle$$

DWBA

$$\langle \varphi_F \chi_{pF}^{(-)} | \Delta V_{pF} | \varphi_d \varphi_A \chi_{dA}^{(+)} \rangle$$

$$\langle I_A^F \chi_{pF}^{(-)} | \Delta V_{pF} | \varphi_d \chi_{dA}^{(+)} \rangle$$

3-body

$$\langle \varphi_F \chi_{pF}^{(-)} | \Delta V_{pF} | \varphi_A \Psi_i^{3B(+)} \rangle$$

$$\langle I_A^F \chi_{pF}^{(-)} | \Delta V_{pF} | \Psi_i^{3B(+)} \rangle$$

CDCC

$$\langle I_A^F \chi_{pF}^{(-)} | \Delta V_{pF} | \Psi_i^{\text{CDCC}(+)} \rangle$$

$\Psi_i^{(+)}$  : exact d+A scattering function  
 $\Phi_f^{(-)} = \varphi_F \chi_{pF}^{(-)}$  exit channel function  
 $\Delta V_{pF} = V_{pA} + V_{pn} - U_{pF}$   
 $I_A^F = \langle \varphi_A | \varphi_F \rangle$

- One-body overlap of A and A+1 systems
- carries structure information
  - typically approximated by single-particle function

$$M^{(\text{prior})} = \langle \Psi_f^{(-)} | \Delta V_{dA} | \Phi_i^{(+)} \rangle$$

$$\Delta V_{dA} = V_{pA} + V_{nA} - U_{dA}$$

# Generalized R-matrix formalism for (d,p) reactions I

Splitting the transition matrix element M:

- Interior and exterior with respect to  $r_{nA}$

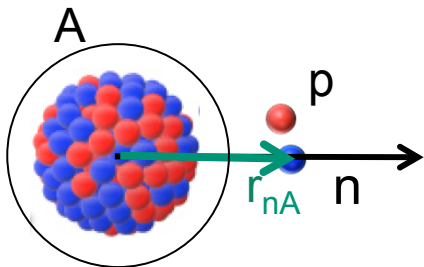
$$M^{(\text{post})} = \langle \Phi_f^{(-)} | \Delta V_{pF} | \Psi_i^{(+)} \rangle$$

DWBA

$$\langle \Phi_F \chi_{pF}^{(-)} | \Delta V_{pF} | \Phi_d \Phi_A \chi_{dA}^{(+)} \rangle$$

$$\langle I_A^F \chi_{pF}^{(-)} | \Delta V_{pF} | \Phi_d \chi_{dA}^{(+)} \rangle$$

3-body system



Interior + exterior

$$M^{(\text{post})} = M^{(\text{post})}(0, a) + M^{(\text{post})}(a, \infty)$$

$$I_A^F = \langle \Phi_A | \Phi_F \rangle = I_A^F(r_{nA})$$

Mukhamedzhanov

$$M^{(\text{post})}(a, \infty) = M_{\text{surf}}(a) + M^{(\text{prior})}(a, \infty)$$

$$M_{\text{surf}}(a) = \langle I_A^F \chi_{pF}^{(-)} | [\overleftarrow{T} - \overrightarrow{T}] | \Phi_d \chi_{dA}^{(+)} \rangle_{\text{ext}}$$

$$\int_{r \leq R} d\mathbf{r} f(\mathbf{r}) [\overleftarrow{T} - \overrightarrow{T}] g(\mathbf{r})$$

$$= -\frac{1}{2\mu} \oint_{r=R} dS [g(\mathbf{r}) \nabla_{\mathbf{r}} f(\mathbf{r}) - f(\mathbf{r}) \nabla_{\mathbf{r}} g(\mathbf{r})]$$

$$= -\frac{1}{2\mu} R^2 \int d\Omega_{\mathbf{r}} \left[ g(\mathbf{r}) \frac{\partial f(\mathbf{r})}{\partial r} - f(\mathbf{r}) \frac{\partial g(\mathbf{r})}{\partial r} \right]_{r=R}$$

Surface term

$$M_{\text{surf}}(a) = f(a, C_A^F, B_{nA})$$

$B_{nA}$  = log derivative of  $I_A^F$  at surface radius  $a$

ANC:  $C_A^F$  defined through:  $I_A^F(r_{nA}) \rightarrow C_A^F W(kr_{nA})$

related to reduced width amplitude  $C_A^F \sim \gamma_{nA}$

## Generalized R-matrix formalism for (d,p) reactions II

DWBA matrix element

$$M^{(\text{post})} = M^{(\text{post})}(0, \mathbf{a}) + M_{(\text{surf})}(\mathbf{a}) + M^{(\text{prior})}(\mathbf{a}, \infty)$$

$$M_{(\text{surf})}(\mathbf{a}) = \sqrt{\frac{R_{nA}}{2\mu_{nA}}} \sum_{j_{nA} m_{j_{nA}} m_{l_{nA}} M_n} \langle J_A M_A j_{nA} m_{j_{nA}} | J_F M_F \rangle \langle J_n M_n l_{nA} m_{l_{nA}} | j_{nA} m_{j_{nA}} \rangle \langle J_p M_p J_n M_n | J_d M_d \gamma_{nA} j_{nA} l_{nA} \rangle \\ \times \int d\mathbf{r}_{pF} \chi_{-\mathbf{k}_{pF}}^{(+)}(\mathbf{r}_{pF}) \int d\Omega_{\hat{\mathbf{r}}_{nA}} Y_{l_{nA} m_{l_{nA}}}^*(\hat{\mathbf{r}}_{nA}) \left[ \varphi_d(\mathbf{r}_{pn}) \chi_{\mathbf{k}_{dA}}^{(+)}(\mathbf{r}_{dA}) (B_{nA} - 1) - R_{nA} \frac{\partial \varphi_d(\mathbf{r}_{pn}) \chi_{\mathbf{k}_{dA}}^{(+)}(\mathbf{r}_{dA})}{\partial r_{nA}} \right] \Big|_{r_{nA}=R_{nA}}$$

### Assessing the approach:

- Internal – external separation sensible?
- Dominant surface term? Size of corrections?
- Study cross sections arising from different terms
- Start with DWBA and bound states
- Investigate resonances

### Cases considered so far:

- $^{90}\text{Zr}(\text{d,p})$  for  $E_d=11$  MeV
  - $^{91}\text{Zr}$  gs, 1<sup>st</sup> excited state,  $2f_{7/2}$  resonance
- $^{48}\text{Ca}(\text{d,p})$  for  $E_d=13$  MeV
  - $^{49}\text{Ca}$  gs, 1<sup>st</sup> excited state
- $^{20}\text{O}(\text{d,p})$  for  $E_d=21$  MeV
  - $^{21}\text{O}$  gs, 1<sup>st</sup> excited state,  $1d_{3/2}$  and  $1f_{7/2}$  resonances
- $^{12}\text{C}(\text{d,p})$  for  $E_d=30$  MeV
- $^{40}\text{Ca}(\text{d,p})$  for  $E_d=34.4$  MeV
- $^{209}\text{Pb}(\text{d,p})$  for  $E_d=52$  MeV
- Planned:  $^{48}\text{Ca}(\text{d,p})$  for  $E_d=19.3$  and  $56$  MeV



# Assessing the R-matrix ideas Ia

## 1. Interior vs exterior contributions

$$M = M(0,a) + M(a,\infty)$$

**This case:**

- $^{90}\text{Zr}(d,p)$  for  $E_d=11$  MeV  
 $^{91}\text{Zr}$  gs (5/2+)  
 1<sup>st</sup> excited state (1/2+)  
 2f<sub>7/2</sub> resonance

← bound

## Observations

- ‘action is in the nuclear surface’
- Post formalism more sensitive to larger radii than prior:

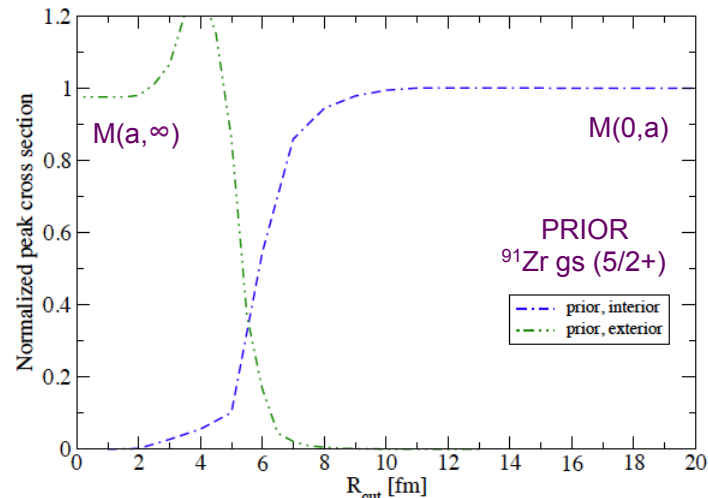
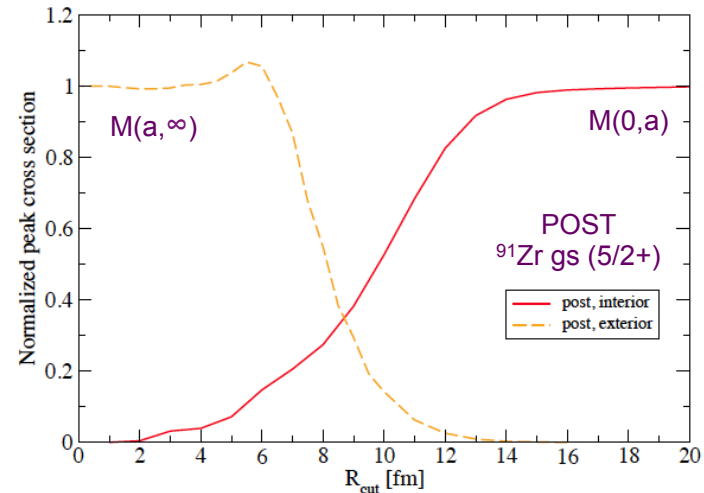
$$M^{(\text{post})} = \langle \Phi_f^{(-)} | \Delta V_{pF} | \Psi_i^{(+)} \rangle$$

$$\Delta V_{pF} = V_{pA} + V_{pn} - U_{pF}$$

$$M^{(\text{prior})} = \langle \Psi_f^{(-)} | \Delta V_{dA} | \Phi_i^{(+)} \rangle$$

$$\Delta V_{dA} = V_{pA} + V_{nA} - U_{dA}$$

Peak cross section relative to full calculation



← R<sub>cut</sub> [fm]

surface radius with respect to  $r_{nA}$

# Assessing the R-matrix ideas Ib

## 1. Interior vs exterior contributions

$$M = M(0,a) + M(a,\infty)$$

### This case:

- $^{90}\text{Zr}(d,p)$  for  $E_d=11$  MeV
- $^{91}\text{Zr}$  gs (5/2+)
- 1<sup>st</sup> excited state (1/2+) ← bound
- $2f_{7/2}$  resonance

### Observations

- ‘action is in the nuclear surface’
- Post formalism more sensitive to larger radii than prior:

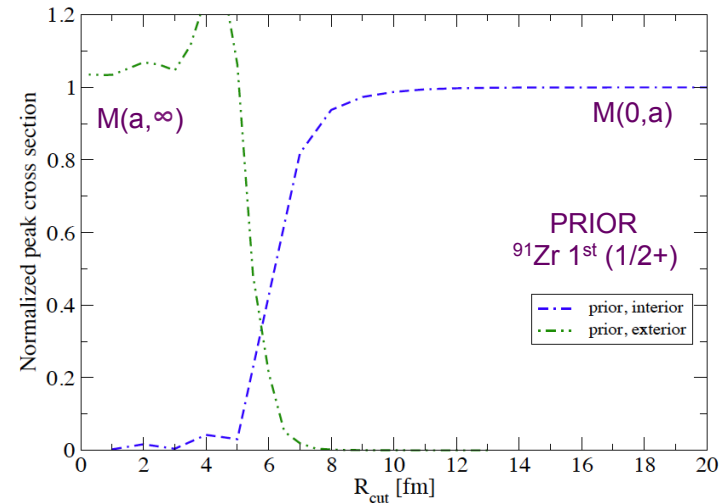
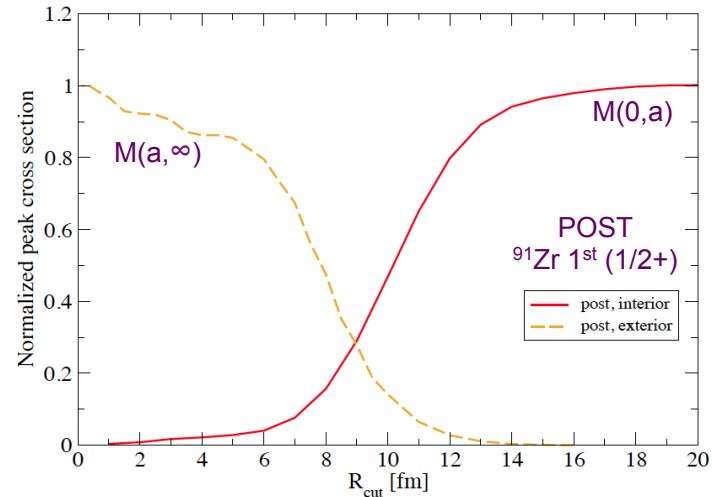
$$M^{(\text{post})} = \langle \Phi_f^{(-)} | \Delta V_{pF} | \Psi_i^{(+)} \rangle$$

$$\Delta V_{pF} = V_{pA} + V_{pn} - U_{pF}$$

$$M^{(\text{prior})} = \langle \Psi_f^{(-)} | \Delta V_{dA} | \Phi_i^{(+)} \rangle$$

$$\Delta V_{dA} = V_{pA} + V_{nA} - U_{dA}$$

Peak cross section relative to full calculation



# Assessing the R-matrix ideas Ic

## 1. Interior vs exterior contributions

$$M = M(0,a) + M(a,\infty)$$

### This case:

- $^{90}\text{Zr}(d,p)$  for  $E_d=11$  MeV
- $^{91}\text{Zr}$  gs ( $5/2+$ )
- 1<sup>st</sup> excited state ( $1/2+$ )
- $2f_{7/2}$  resonance

← resonance

### Observations

- 'action is in the nuclear surface'
- Post formalism more sensitive to larger radii than prior:

$$M^{(\text{post})} = \langle \Phi_f^{(-)} | \Delta V_{pF} | \Psi_i^{(+)} \rangle$$

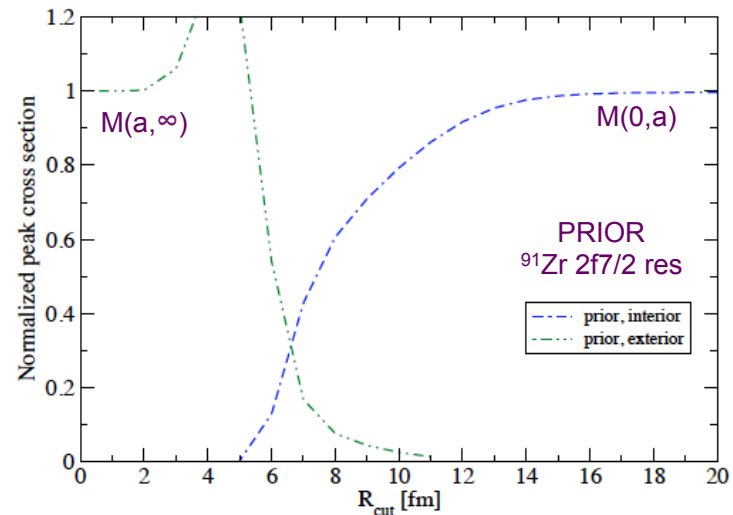
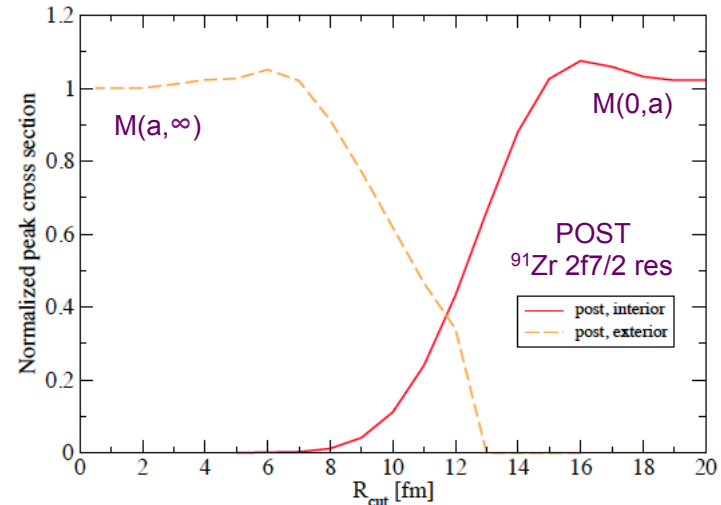
$$\Delta V_{pF} = V_{pA} + V_{pn} - U_{pF}$$

$$M^{(\text{prior})} = \langle \Psi_f^{(-)} | \Delta V_{dA} | \Phi_i^{(+)} \rangle$$

$$\Delta V_{dA} = V_{pA} + V_{nA} - U_{dA}$$

- Resonance: reduced contribution from interior, more pronounced surface effect

Peak cross section relative to full calculation



## Assessing the R-matrix ideas IIa

### 2. Surface contribution

$$M = M^{(\text{post})}(0, a) + M_{(\text{surf})}(a) + M^{(\text{prior})}(a, \infty)$$

#### This case:

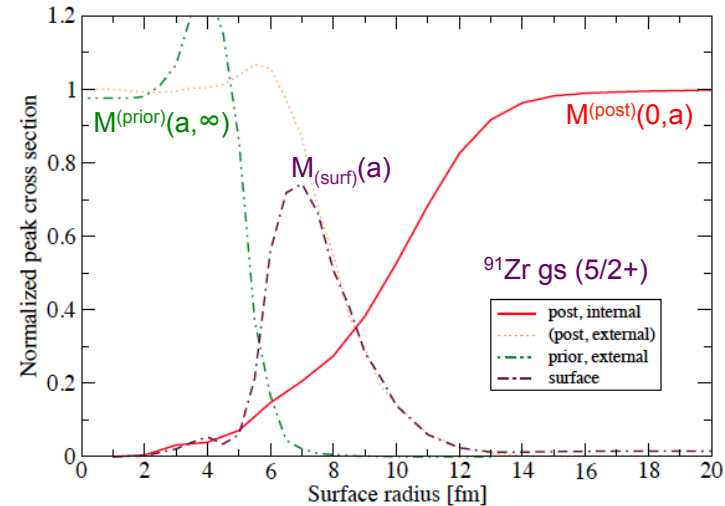
- $^{90}\text{Zr}(d, p)$  for  $E_d = 11$  MeV
- $^{91}\text{Zr}$  gs (5/2+)
- 1<sup>st</sup> excited state (1/2+)
- $2f_{7/2}$  resonance

← bound

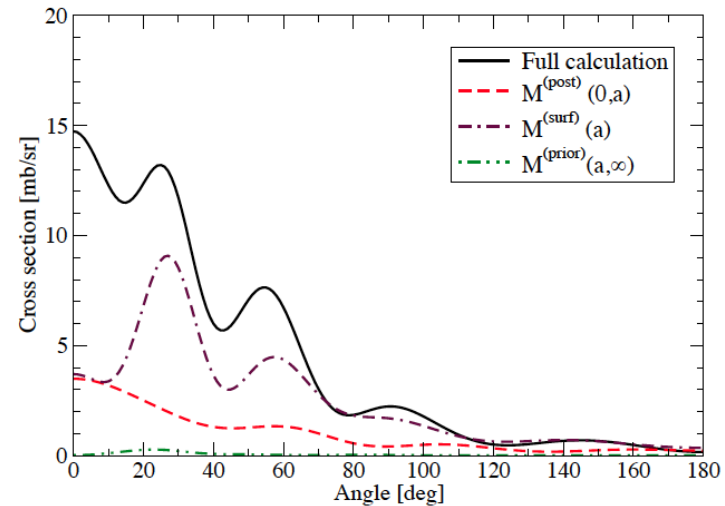
#### Observations

- Surface term indeed dominant 6-8 fm
- Small interior contributions → little dependence on model for interior
- Small exterior contributions → better convergence for resonance case
- Surface term does not produce the whole cross section, corrections required from internal/external

Peak cross section relative to full calculation



Angular cross sections for  $a=7$  fm



# Assessing the R-matrix ideas IIb

## 2. Surface contribution

$$M = M^{(post)}(0,a) + M_{(surf)}(a) + M^{(prior)}(a,\infty)$$

### This case:

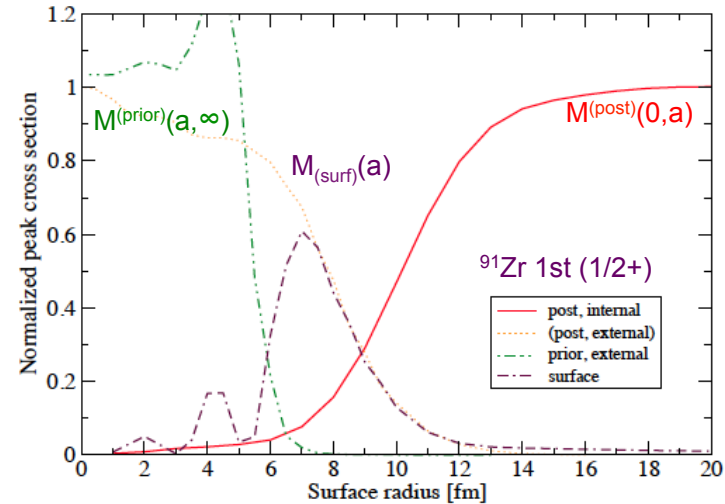
- $^{90}\text{Zr}(d,p)$  for  $E_d=11$  MeV
- $^{91}\text{Zr}$  gs (5/2+)
- 1<sup>st</sup> excited state (1/2+)
- $2f_{7/2}$  resonance

← bound

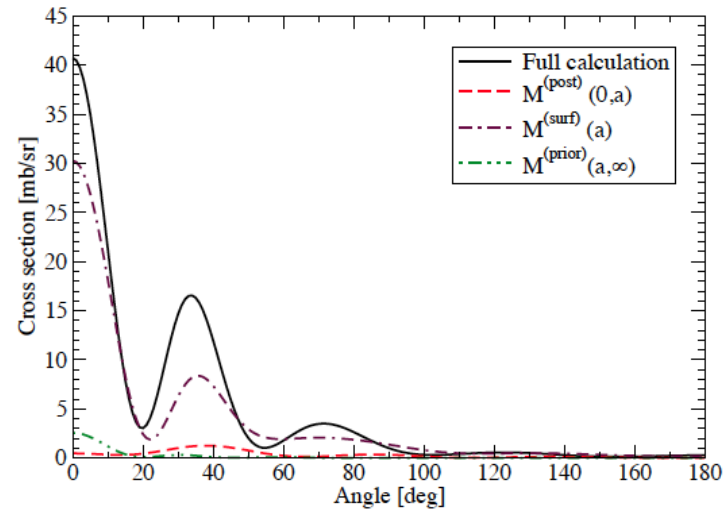
### Observations

- Surface term indeed dominant 6-8 fm
- Small interior contributions → little dependence on model for interior
- Small exterior contributions → better convergence for resonance case
- Surface term does not produce the whole cross section, corrections required from internal/external

Peak cross section relative to full calculation



Angular cross sections for a=7 fm



# Assessing the R-matrix ideas IIc

## 2. Surface contribution

$$M = M^{(\text{post})}(0, a) + M_{(\text{surf})}(a) + M^{(\text{prior})}(a, \infty)$$

### This case:

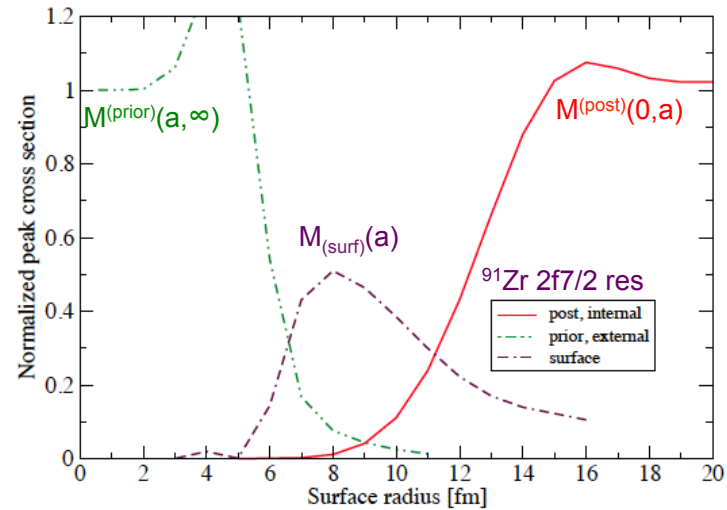
- $^{90}\text{Zr}(d, p)$  for  $E_d = 11$  MeV
- $^{91}\text{Zr}$  gs (5/2+)
- 1<sup>st</sup> excited state (1/2+)
- $2f_{7/2}$  resonance

← resonance

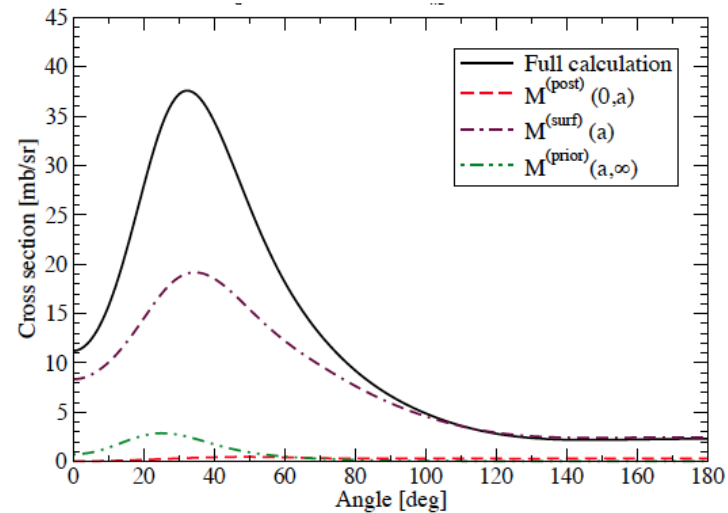
### Observations

- Surface term indeed dominant 6-8 fm
- Small interior contributions → little dependence on model for interior
- Small exterior contributions → better convergence for resonance case
- Surface term does not produce the whole cross section, corrections required from internal/external
- Reduced interior contribution at peak for surface term

Peak cross section relative to full calculation



Angular cross sections for a = 8 fm



# Assessing the R-matrix ideas - $^{48}\text{Ca}$

## 2. Surface contribution

$$M = M^{(\text{post})}(0, a) + M_{(\text{surf})}(a) + M^{(\text{prior})}(a, \infty)$$

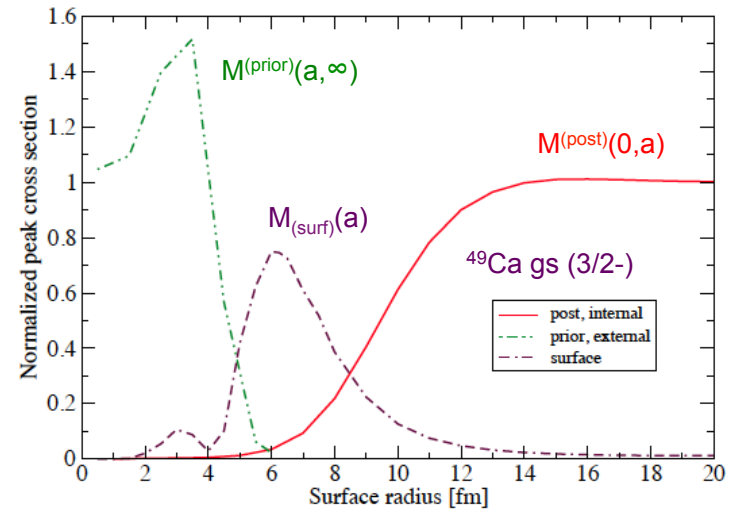
### This case:

- $^{48}\text{Ca}(d, p)$  for  $E_d = 13$  MeV  
 $^{49}\text{Ca}$  gs (3/2-) ← bound  
 1<sup>st</sup> excited state (1/2-)

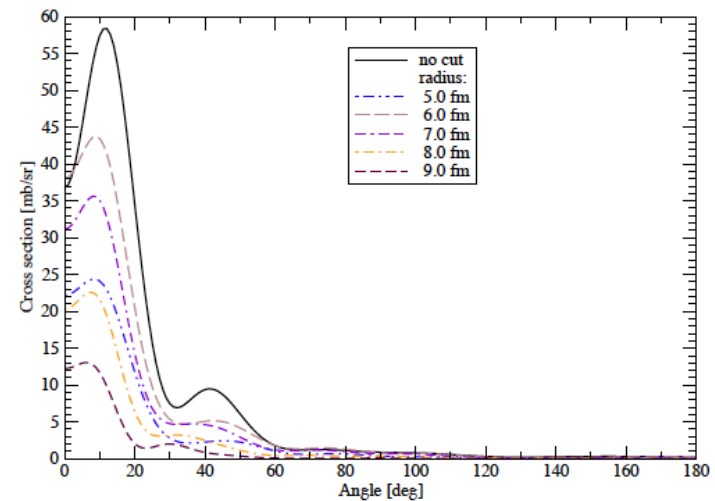
### Observations

- Surface term indeed dominant 5-8 fm
- Small interior contributions → little dependence on model for interior
- Small exterior contributions → better convergence for resonance case
- Surface term does not produce the whole cross section, corrections required from internal/external
- Cross section calculated from surface term sensitive to surface radius

Peak cross section relative to full calculation



Surface angular cross sections for various a values



# Assessing the R-matrix ideas - $^{48}\text{Ca}$

## 2. Surface contribution

$$M = M^{(\text{post})}(0, a) + M_{(\text{surf})}(a) + M^{(\text{prior})}(a, \infty)$$

### This case:

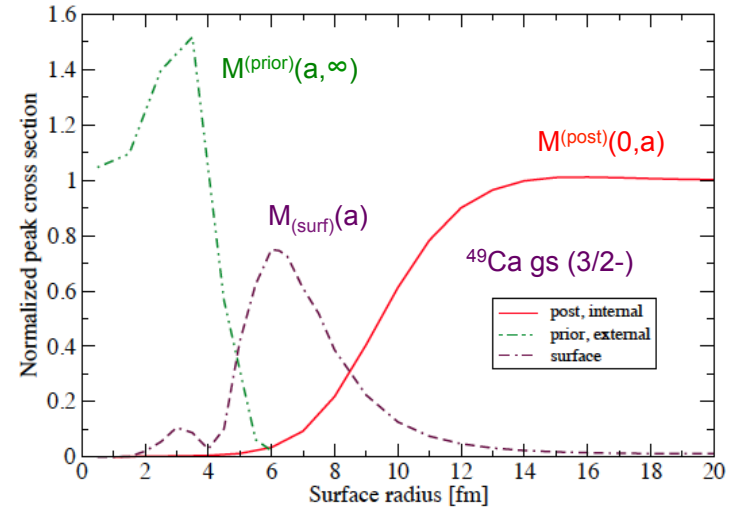
- $^{48}\text{Ca}(d, p)$  for  $E_d = 13$  MeV
- $^{49}\text{Ca}$  gs (3/2-)
- 1<sup>st</sup> excited state (1/2-)

← bound

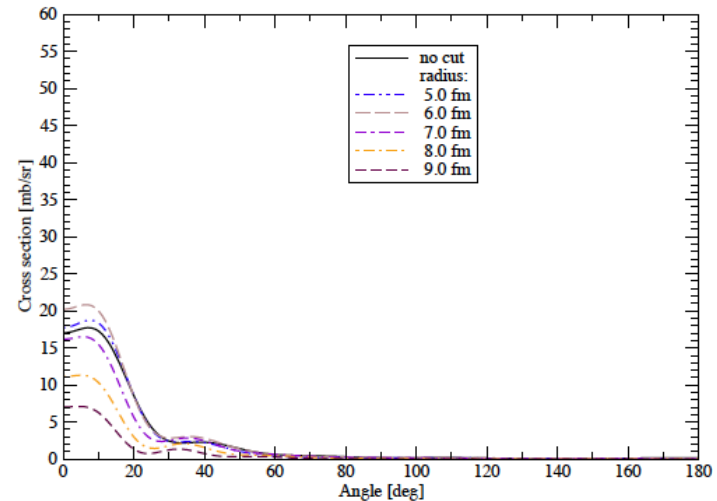
### Observations

- Surface term indeed dominant 5-7 fm
- Small interior contributions → little dependence on model for interior
- Small exterior contributions → better convergence for resonance case
- Surface term does not produce the whole cross section, corrections required from internal/external
- Cross section calculated from surface term sensitive to surface radius

Peak cross section relative to full calculation



Surface angular cross sections for various a values





# Assessing the R-matrix ideas – $^{20}\text{O}$

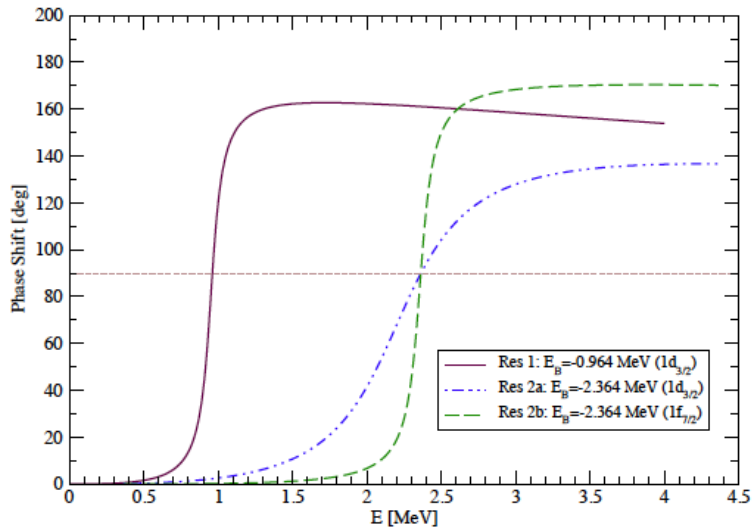
In progress

Data for both bound and resonance states available!

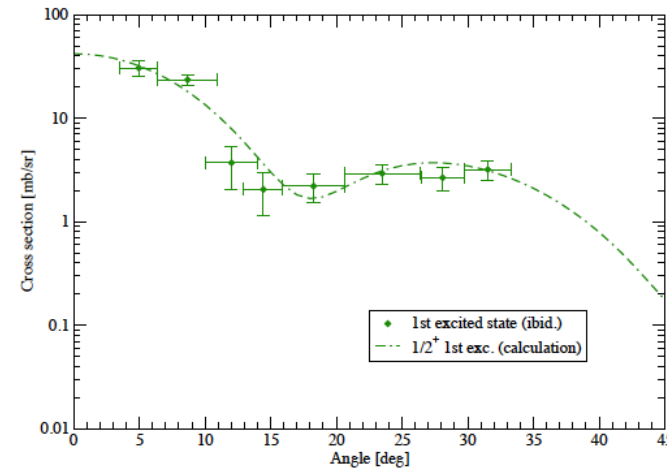
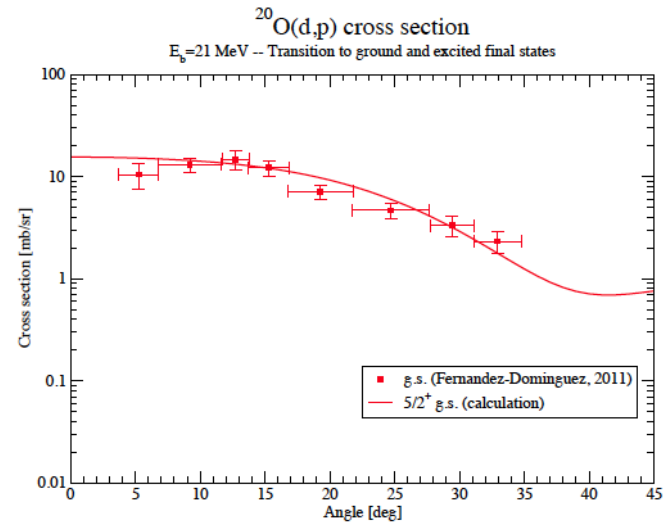
**This case:**

- $^{20}\text{O}(d,p)$  for  $E_d=21$  MeV
  - $^{20}\text{O}$  gs (3/2-) ← bound
  - 1<sup>st</sup> excited state (1/2-) ← bound
  - 1d<sub>3/2</sub> and 1f<sub>7/2</sub> resonances ← resonance

Phase shifts for resonances considered



Cross sections for bound states compared to data



## Next: Extension of the formalism to include breakup

### DWBA matrix element

$$M^{(\text{post})} = M^{(\text{post})}(0, a) + M_{(\text{surf})}(a) + M^{(\text{prior})}(a, \infty)$$

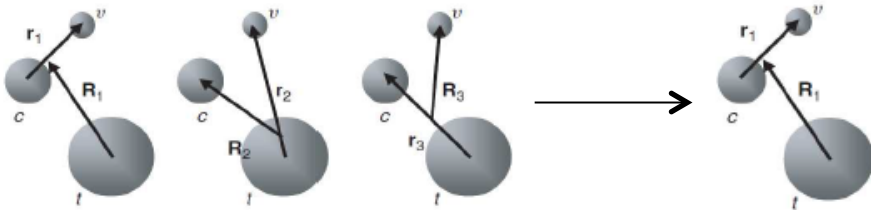
### CDCC matrix element

$$M^{(\text{post})} = M^{(\text{post})}(0, a) + M_{(\text{surf})}(a)$$

$$M^{(\text{prior})}(a, \infty) = 0 \text{ (is included in breakup)}$$

### CDCC (Continuum-discretized coupled channels)

- Approximate treatment of 3-body problem
- Describes breakup of deuteron



- Successfully used for describing data
- Currently revisited via comparison with Fadeev

### CDCC extension of R-matrix formalism

- Simultaneous calculation of breakup and transfer cross sections
- Exterior term included in breakup, convergence issues removed
- More peripheral, reduce interior contribution
- Surface term dominant

# Conclusions

Studying resonances with (d,p):

- Already underway at RIB facilities
- Conceptual and practical problems have to be overcome

## **New formalism:**

- Builds on ideas from successful R-matrix approach
- Separation into interior and exterior regions works formally well, surface term emerges as important contributor, can be expressed in terms of familiar R-matrix parameters -> meaningful spectroscopic information
- Test cases show that the surface term is dominant; other contributions may not be negligible, but resonances less affected by interior contributions
- Including breakup via CDCC removes exterior prior contribution, thus eliminates convergence problem for resonances

**Further studies will clarify conditions where the surface formalism will work well.**

**Promising approach for transfers to resonances.**

# TORUS Collaboration

# ReactionTheory.org

**TORUS: Theory of Reactions for Unstable iSotopes**  
**A Topical Collaboration for Nuclear Theory**

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[Research Proposal](#)

[Research Papers](#)

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[Workshops](#)

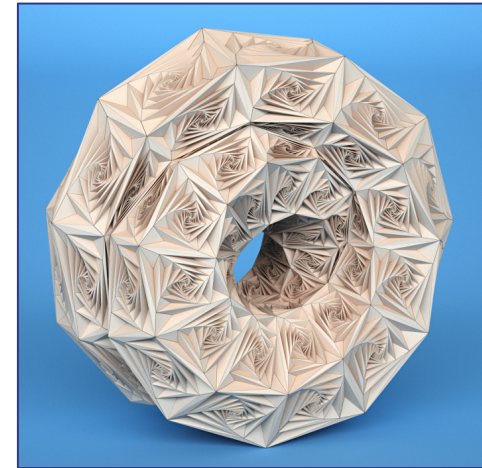
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## Theory of Reactions for Unstable iSotopes

A Topical Collaboration to develop new methods that will advance nuclear reaction theory for unstable isotopes by using three-body techniques to improve direct-reaction calculations and by developing a new partial-fusion theory to integrate descriptions of direct and compound-nucleus reactions. This multi-institution collaborative effort is directly relevant to three areas of interest identified in the solicitation: (b) properties of nuclei far from stability; (c) microscopic studies of nuclear input parameters for astrophysics and (e) microscopic nuclear reaction theory.



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# Appendix

# Investigating the role of the core-core interaction $V_{pA}$

## 3. $V_{pA}$ dependence of the various contributions

$$M = M^{(\text{post})}(0,a) + M_{(\text{surf})}(a) + M^{(\text{prior})}(a,\infty)$$

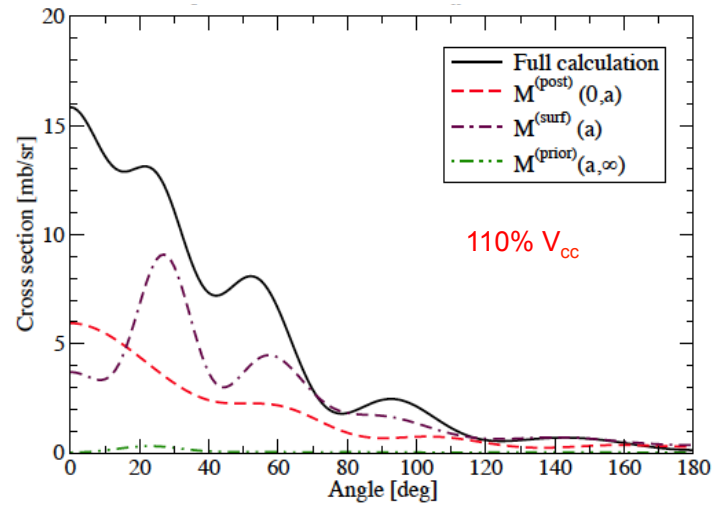
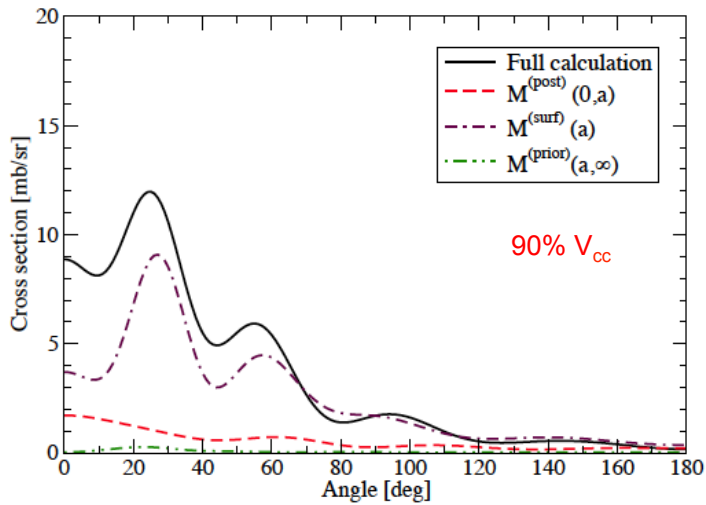
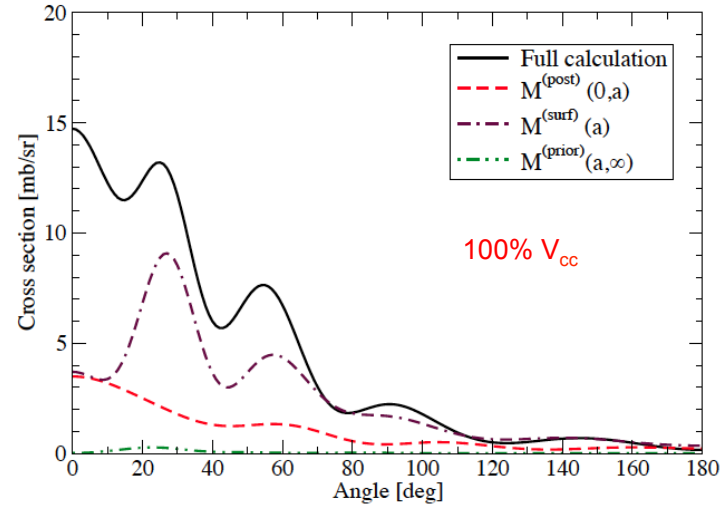
### This case:

- $^{90}\text{Zr}(d,p)$  for  $E_d=11$  MeV
- $^{91}\text{Zr}$  gs (5/2+) ← bound
- 1<sup>st</sup> excited state (1/2+)
- $2f_{7/2}$  resonance

### Observations

- Overall cross section and relative strength of contributions varies with the strength of the core-core interaction

Angular cross sections for  $a=7$  fm



# Investigating the role of the core-core interaction $V_{pA}$

## 3. $V_{pA}$ dependence of the various contributions

$$M = M^{(\text{post})}(0,a) + M^{(\text{surf})}(a) + M^{(\text{prior})}(a,\infty)$$

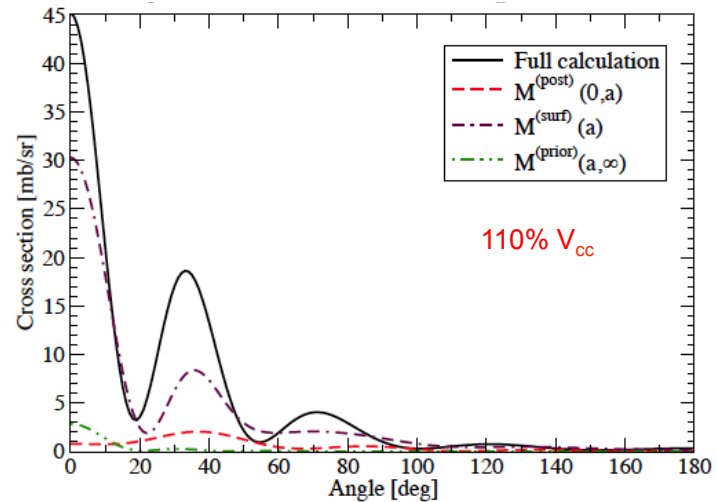
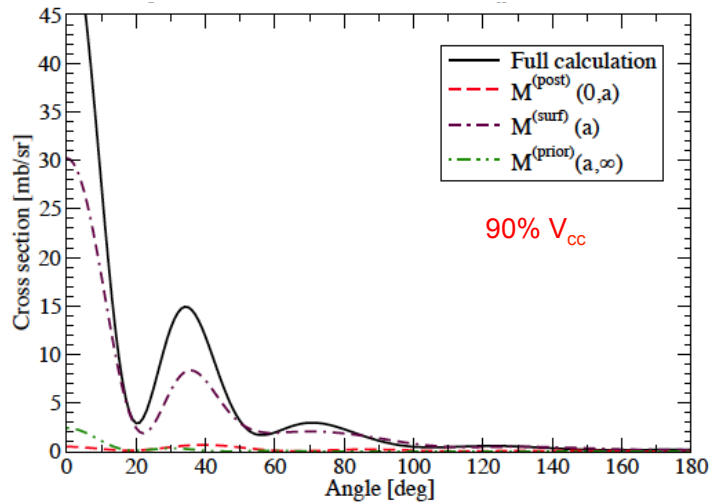
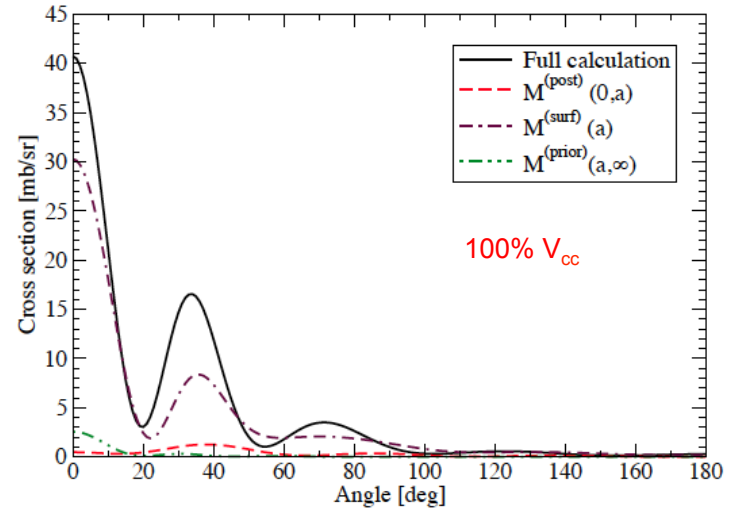
### This case:

- $^{90}\text{Zr}(d,p)$  for  $E_d=11$  MeV
- $^{91}\text{Zr}$  gs (5/2+)
- 1<sup>st</sup> excited state (1/2+) ← bound
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### Observations

- Overall cross section and relative strength of contributions varies with the strength of the core-core interaction

Angular cross sections for  $a=7$  fm



# Exploring R-matrix ideas for (d,p) one-nucleon transfers II

Transition matrix element M:

- Connects initial to final wave function
- Cross section  $\sigma \sim M^2$

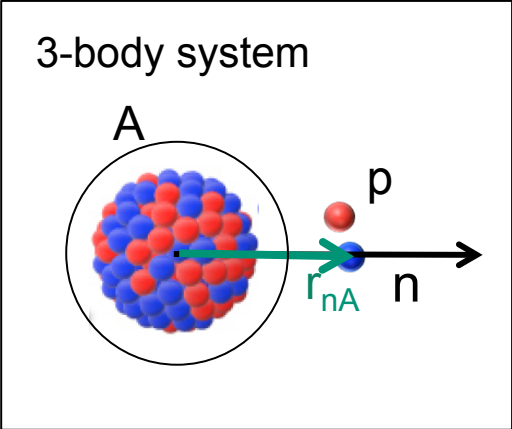
$$M^{(\text{post})} = \langle \Phi_f^{(-)} | \Delta V_{pF} | \Psi_i^{(+)} \rangle$$

DWBA

$$\langle \varphi_F \chi_{pF}^{(-)} | \Delta V_{pF} | \varphi_d \varphi_A \chi_{dA}^{(+)} \rangle$$

$$\langle I_A^F \chi_{pF}^{(-)} | \Delta V_{pF} | \varphi_d \chi_{dA}^{(+)} \rangle$$

3-body  
 $\approx$   
CDCC



$\Psi_i^{(+)}$  : exact d+A scattering function  
 $\Phi_f^{(-)} = \varphi_F \chi_{pF}^{(-)}$  exit channel function  
 $\Delta V_{pF} = V_{pA} + V_{pn} - U_{pF}$   
 $I_A^F = \langle \varphi_A | \varphi_F \rangle = I_A^F(r_{nA})$

One-body overlap of A and A+1 systems

- carries structure information
- typically approximated by single-particle function

$$M^{(\text{prior})} = \langle \Psi_f^{(-)} | \Delta V_{dA} | \Phi_i^{(+)} \rangle$$

$$\Delta V_{dA} = V_{pA} + V_{nA} - U_{dA}$$



# Assessing the R-matrix ideas - $^{48}\text{Ca}$

## 2. Surface contribution

$$M = M^{(\text{post})}(0, a) + M_{(\text{surf})}(a) + M^{(\text{prior})}(a, \infty)$$

### This case:

- $^{48}\text{Ca}(d, p)$  for  $E_d = 13$  MeV
  - $^{49}\text{Ca}$  gs (3/2-) ← bound
  - 1<sup>st</sup> excited state (1/2-) ← bound

### Observations

- Surface term indeed dominant 5-7 fm
- Small interior contributions → little dependence on model for interior
- Small exterior contributions → better convergence for resonance case
- Surface term does not produce the whole cross section, corrections required from internal/external

Peak cross section relative to full calculation

