

Exploring R-matrix ideas for the description of one-nucleon transfer reactions

NHEP Section talk
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TORUS Collaboration



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TORUS Collaboration

ReactionTheory.org

TORUS: Theory of Reactions for Unstable iSotopes
A Topical Collaboration for Nuclear Theory

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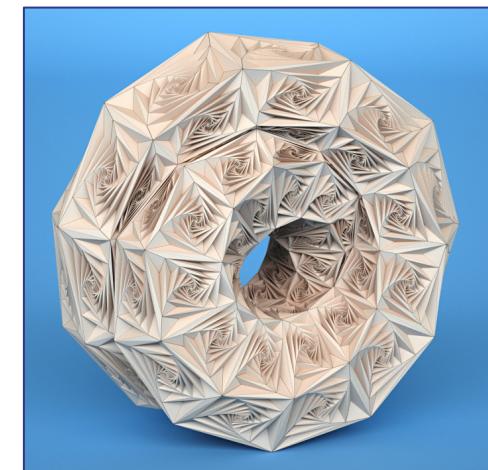
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Theory of Reactions for Unstable iSotopes

A Topical Collaboration to develop new methods that will advance nuclear reaction theory for unstable isotopes by using three-body techniques to improve direct-reaction calculations and by developing a new partial-fusion theory to integrate descriptions of direct and compound-nucleus reactions. This multi-institution collaborative effort is directly relevant to three areas of interest identified in the solicitation: (b) properties of nuclei far from stability; (c) microscopic studies of nuclear input parameters for astrophysics and (e) microscopic nuclear reaction theory.



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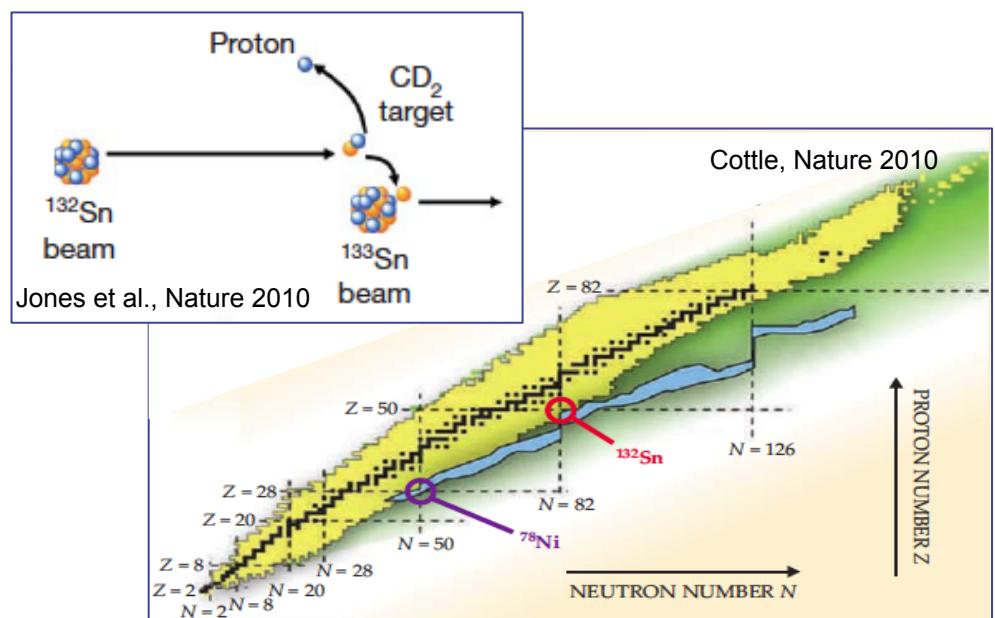
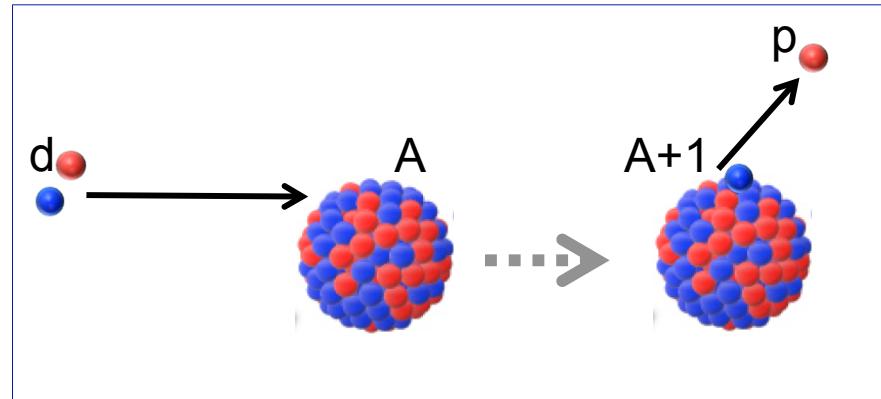
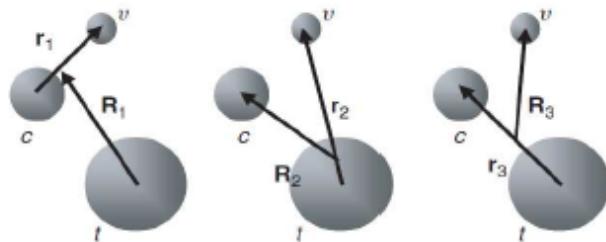
Studying nuclear structure with (d,p) one-nucleon transfers

(d,p) reactions:

- Simplest mechanism for adding a neutron
- Traditionally used to study stable nuclei
- Used in inverse kinematics at RIB facilities, for studying weakly-bound systems

Theoretical descriptions of (d,p) reactions:

- Progress over the years: Plane-wave theory, DWBA (zero-range & finite-range), coupled-channels approach, breakup, etc.



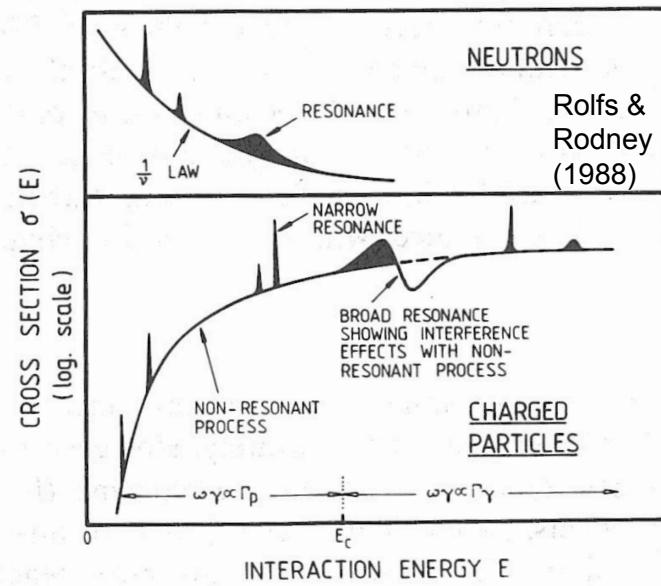
Current status of (d,p) direct-reaction theories:

- Developing Fadeev techniques to better account for 3-body effects (TORUS collab.)
- Conceptual work needed: rethinking spectroscopic factors
- Not very useful for transfers to resonance states.

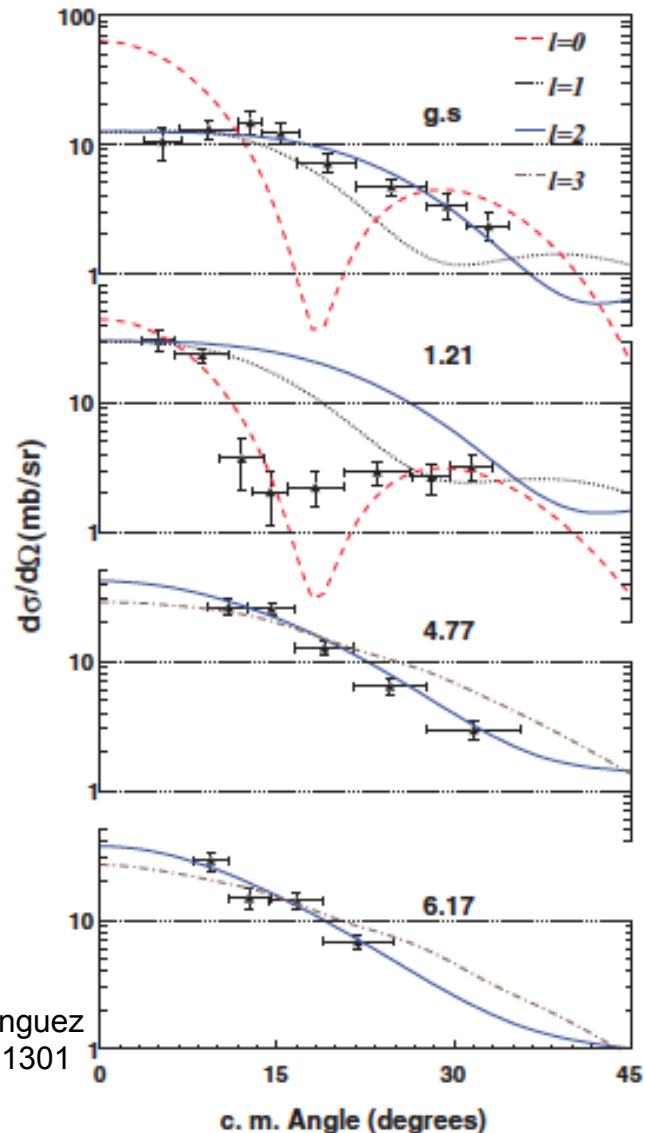
Resonances in low-energy nuclear physics

Resonances:

- Unstable quantum-mechanical states
- Occur in light, medium-mass, and heavy nuclei
- Crucially affect astrophysical reaction rates
- Abundant in weakly-bound nuclei



$^{20}\text{O}(\text{d},\text{p})^{21}\text{O}$ inverse-kinematics experiment at GANIL to determine N=16 shell gap



Problems in applying standard method to resonances:

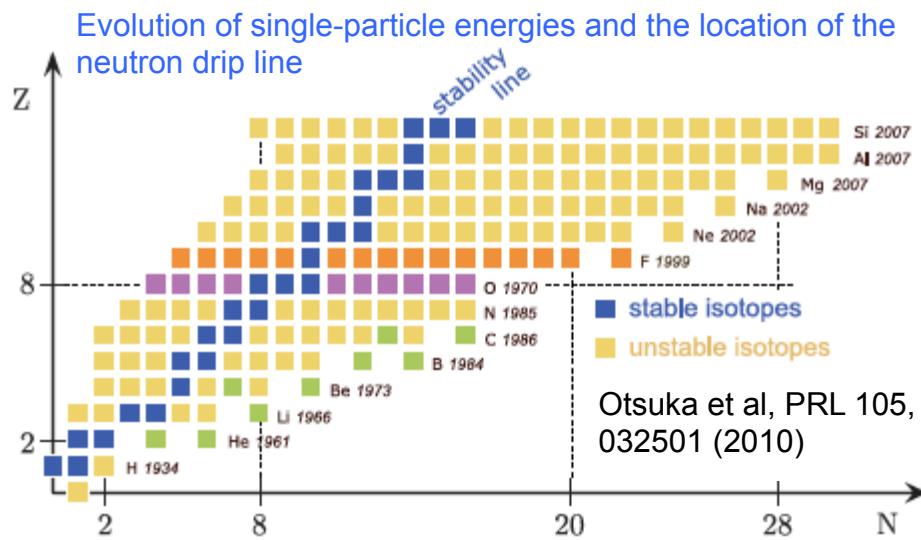
- Conceptual: meaning of spectroscopic factor?
- Practical: convergence issues

Fernandez-Dominguez et al, PRC 84, 011301 (R) (2011)

Resonances in low-energy nuclear physics

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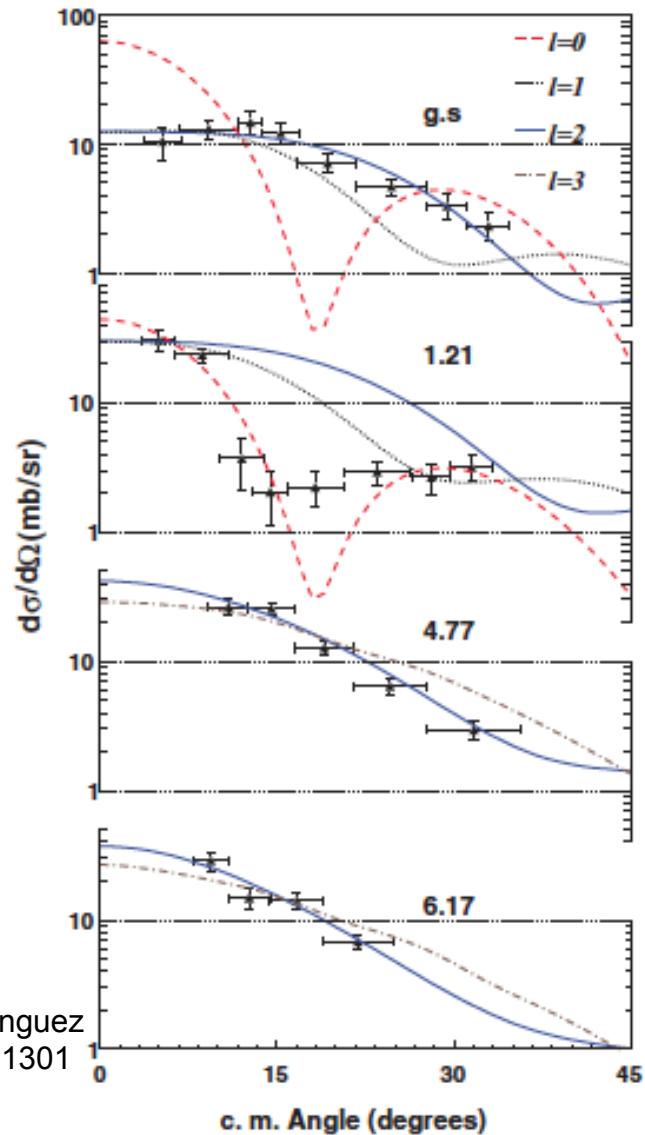


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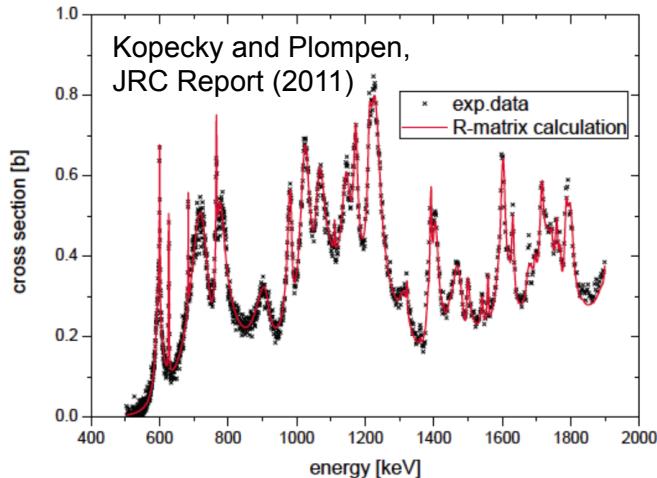
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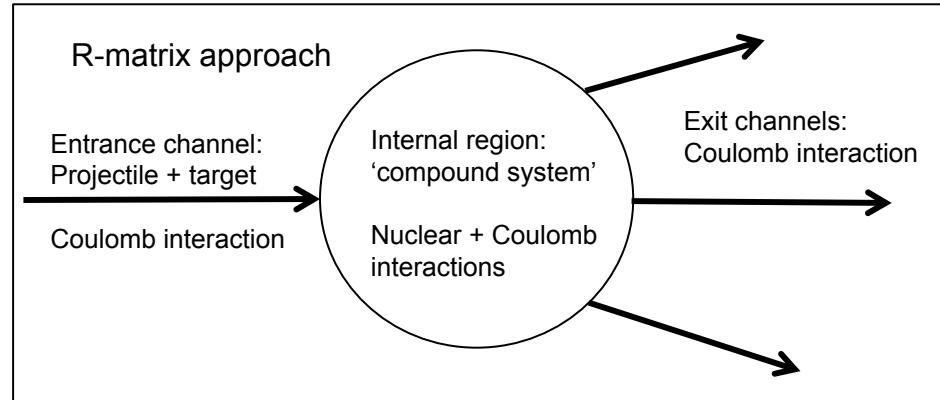
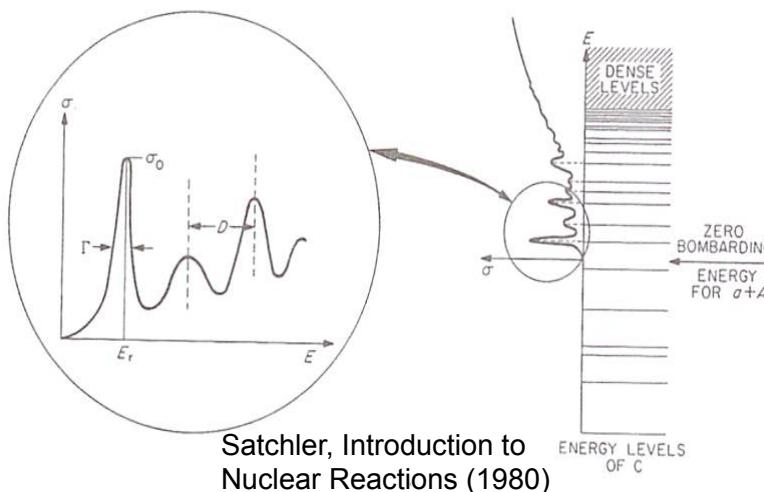
Describing resonances in binary reactions

Experimental studies of resonances:

- Elastic & inelastic scattering, capture, etc.



- Characterization of resonances: position & widths



R-matrix approach:

- Main idea: divide space into 2 regions:
 $r \leq a$ -- interior: nuclear and Coulomb interactions
 $r > a$ -- exterior: Coulomb only
- Formalism:
 Interior: set of basis functions to express nuclear wave function
 Exterior: scattering wave function
 Surface: matching conditions allow to parameterize collision matrix \rightarrow expressions for cross sections
- Connect observed parameters (E_R , Γ) to formal parameters (\tilde{E}_R , γ^2)
- Typical applications adjust parameters to reproduce measured cross sections

Exploring R-matrix ideas for (d,p) one-nucleon transfers

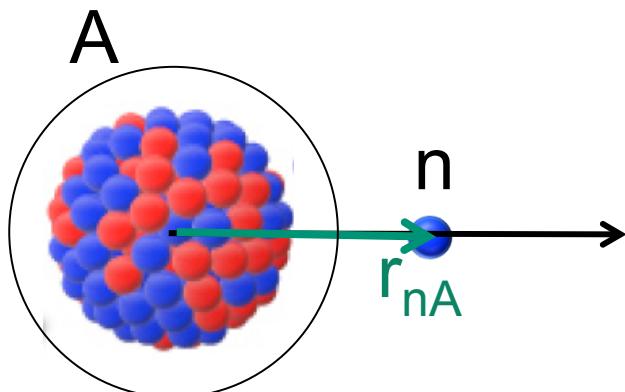
Proposed new formalism (Mukhamedzhanov, 2011):

- R-matrix concepts:
 - surface separating internal and external regions
 - cross sections expressed in terms of reduced widths, logarithmic derivatives, surface radii
- Applicable to stripping to bound and resonance states
- Provides conceptually improved way to describe (d,p) transfer reactions
- General enough to include deuteron breakup contributions via CDCC
- Resolves practical issues related to numerical convergence

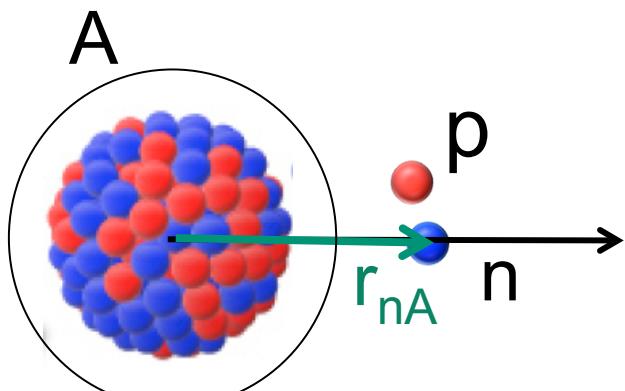
Formalism:

Mukhamedzhanov, PRC 84, 044616 (2011)

Binary system



3-body system



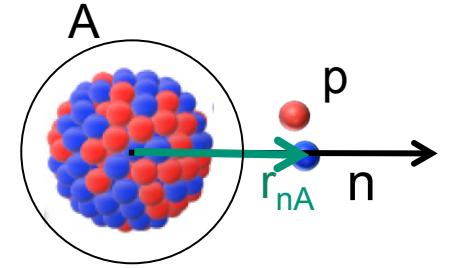
Exploring R-matrix ideas for (d,p) one-nucleon transfers II

Transition matrix element M:

- Connects initial to final wave function
- Cross section $\sigma \sim M^2$

$$M^{(\text{post})} = \langle \Phi_f^{(-)} | \Delta V_{pF} | \Psi_i^{(+)} \rangle$$

3-body system



DWBA

$$\langle \Phi_F \chi_{pF}^{(-)} | \Delta V_{pF} | \Phi_d \Psi_A \chi_{dA}^{(+)} \rangle$$

$$\langle I_A^F \chi_{pF}^{(-)} | \Delta V_{pF} | \Phi_d \chi_{dA}^{(+)} \rangle$$

3-body

$$\langle \Phi_F \chi_{pF}^{(-)} | \Delta V_{pF} | \Psi_A \Psi_i^{3B(+)} \rangle$$

$$\langle I_A^F \chi_{pF}^{(-)} | \Delta V_{pF} | \Psi_i^{3B(+)} \rangle$$

CDCC

$$\langle I_A^F \chi_{pF}^{(-)} | \Delta V_{pF} | \Psi_i^{\text{CDCC}(+)} \rangle$$

$\Psi_i^{(+)}$: exact d+A scattering function

$\Phi_f^{(-)}$ = $\Phi_F \chi_{pF}^{(-)}$ exit channel function

$\Delta V_{pF} = V_{pA} + V_{pn} - U_{pF}$

$I_A^F = \langle \Phi_A | \Phi_F \rangle$ one-body overlap

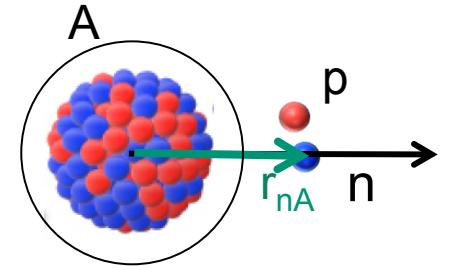
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DWBA

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$$\langle I_A^F \chi_{pF}^{(-)} | \Delta V_{pF} | \Phi_d \chi_{dA}^{(+)} \rangle$$

3-body

$$\langle \Phi_F \chi_{pF}^{(-)} | \Delta V_{pF} | \Psi_A \Psi_i^{3B(+)} \rangle$$

$$\langle I_A^F \chi_{pF}^{(-)} | \Delta V_{pF} | \Psi_i^{3B(+)} \rangle$$

CDCC

$$\langle I_A^F \chi_{pF}^{(-)} | \Delta V_{pF} | \Psi_i^{\text{CDCC}(+)} \rangle$$

$\Psi_i^{(+)}$: exact d+A scattering function

$\Phi_f^{(-)}$ = $\Phi_F \chi_{pF}^{(-)}$ exit channel function

$\Delta V_{pF} = V_{pA} + V_{pn} - U_{pF}$

$I_A^F = \langle \Phi_A | \Phi_F \rangle$ one-body overlap

$$M^{(\text{prior})} = \langle \Psi_f^{(-)} | \Delta V_{dA} | \Phi_i^{(+)} \rangle$$

$$\Delta V_{dA} = V_{pA} + V_{nA} - U_{dA}$$

Generalized R-matrix formalism for (d,p) reactions I

Splitting the transition matrix element M:

- Interior and exterior with respect to r_{nA}

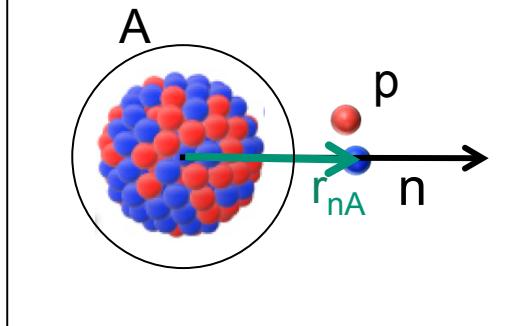
$$M^{(\text{post})} = \langle \Phi_f^{(-)} | \Delta V_{pF} | \Psi_i^{(+)} \rangle$$

DWBA

$$\langle \Phi_F \chi_{pF}^{(-)} | \Delta V_{pF} | \Phi_d \Phi_A \chi_{dA}^{(+)} \rangle$$

$$\langle I_A^F \chi_{pF}^{(-)} | \Delta V_{pF} | \Phi_d \chi_{dA}^{(+)} \rangle$$

3-body system



Interior + exterior

$$M^{(\text{post})} = M^{(\text{post})}(0, a) + M^{(\text{post})}(a, \infty)$$

$$I_A^F = \langle \Phi_A | \Phi_F \rangle = I_A^F(r_{nA})$$

Mukhamedzhanov

$$M^{(\text{post})}(a, \infty) = M_{\text{surf}}(a) + M^{(\text{prior})}(a, \infty)$$

$$M_{\text{surf}}(a) = \langle I_A^F \chi_{pF}^{(-)} | [\overleftarrow{T} - \overrightarrow{T}] | \Phi_d \chi_{dA}^{(+)} \rangle_{\text{ext}}$$

$$\int_{r \leqslant R} d\mathbf{r} f(\mathbf{r}) [\overleftarrow{T} - \overrightarrow{T}] g(\mathbf{r})$$

$$= -\frac{1}{2\mu} \oint_{r=R} dS [g(\mathbf{r}) \nabla_{\mathbf{r}} f(\mathbf{r}) - f(\mathbf{r}) \nabla_{\mathbf{r}} g(\mathbf{r})]$$

$$= -\frac{1}{2\mu} R^2 \int d\Omega_{\mathbf{r}} \left[g(\mathbf{r}) \frac{\partial f(\mathbf{r})}{\partial r} - f(\mathbf{r}) \frac{\partial g(\mathbf{r})}{\partial r} \right]_{r=R}$$

Surface term

$$M_{\text{surf}}(a) = f(a, C_A^F, B_{nA})$$

B_{nA} = log derivative of I_A^F at surface radius a

ANC: C_A^F defined through: $I_A^F(r_{nA}) \rightarrow C_A^F W(kr_{nA})$
related to reduced width amplitude $C_A^F \sim \gamma_{nA}$

Generalized R-matrix formalism for (d,p) reactions II

DWBA matrix element

$$M^{(\text{post})} = M^{(\text{post})}(0, a) + M_{(\text{surf})}(a) + M^{(\text{prior})}(a, \infty)$$

$$\begin{aligned} M_{(\text{surf})}(a) &= \sqrt{\frac{R_{nA}}{2\mu_{nA}}} \sum_{j_{nA} m_{j_{nA}} m_{l_{nA}} M_n} \langle J_A M_A j_{nA} m_{j_{nA}} | J_F M_F \rangle \langle J_n M_n l_{nA} m_{l_{nA}} | j_{nA} m_{j_{nA}} \rangle \langle J_p M_p J_n M_n | J_d M_d \rangle \gamma_{nA j_{nA} l_{nA}} \\ &\times \int d\mathbf{r}_{pF} \chi_{-\mathbf{k}_{pF}}^{(+)}(\mathbf{r}_{pF}) \int d\Omega_{\mathbf{r}_{nA}} Y_{l_{nA} m_{l_{nA}}}^*(\hat{\mathbf{r}}_{nA}) \left[\varphi_d(\mathbf{r}_{pn}) \chi_{\mathbf{k}_{dA}}^{(+)}(\mathbf{r}_{dA}) (B_{nA} - 1) - R_{nA} \frac{\partial \varphi_d(\mathbf{r}_{pn}) \chi_{\mathbf{k}_{dA}}^{(+)}(\mathbf{r}_{dA})}{\partial r_{nA}} \right] \Big|_{r_{nA}=R_{nA}} \end{aligned}$$

Assessing the approach:

- Separation into internal and external regions sensible?
- Is the surface term dominant and what is the size of the corrections?
- Study cross sections arising from different terms
- Start with DWBA and bound states
- Investigate resonances

Cases considered so far:

- $^{90}\text{Zr}(\text{d},\text{p})$ for $E_d = 11$ MeV
 - ^{91}Zr gs, 1st excited state, $2f_{7/2}$ resonance
- $^{48}\text{Ca}(\text{d},\text{p})$ for $E_d = 13$ MeV
 - ^{49}Ca gs, 1st excited state
- $^{12}\text{C}(\text{d},\text{p})$ for $E_d = 30$ MeV
- $^{40}\text{Ca}(\text{d},\text{p})$ for $E_d = 34.4$ MeV
- $^{209}\text{Pb}(\text{d},\text{p})$ for $E_d = 52$ MeV
- Planned: $^{48}\text{Ca}(\text{d},\text{p})$ for $E_d = 19.3$ and 56 MeV

Assessing the R-matrix ideas Ia

1. Interior vs exterior contributions

$$M = M(0,a) + M(a,\infty)$$

This case:

- $^{90}\text{Zr}(\text{d},\text{p})$ for $E_d = 11 \text{ MeV}$
- ^{91}Zr gs ($5/2^+$)
- 1st excited state ($1/2^+$)
- $2f_{7/2}$ resonance



Observations

- 'action is in the nuclear surface'
- Post formalism more sensitive to larger radii than prior:

$$M^{(\text{post})} = \langle \Phi_f^{(-)} | \Delta V_{pF} | \Psi_i^{(+)} \rangle$$

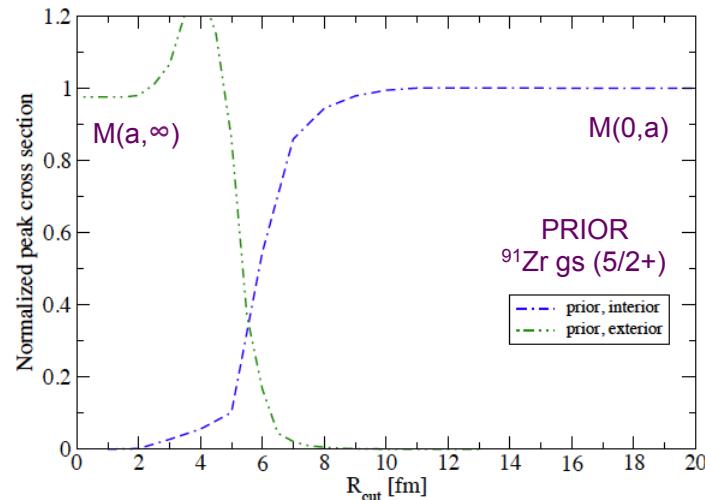
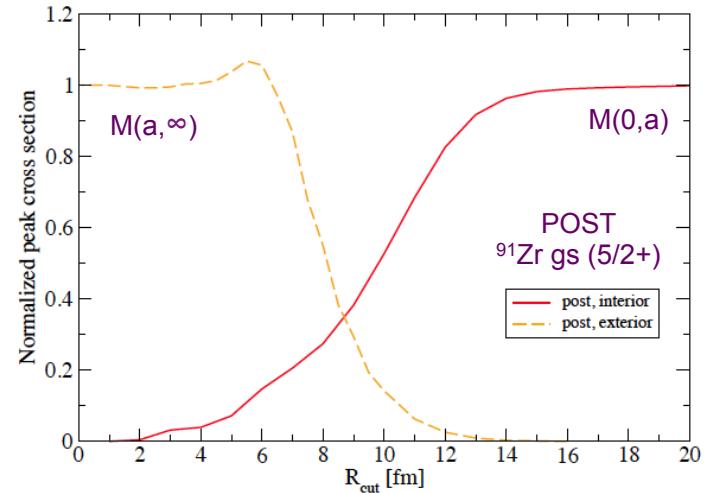
$$\Delta V_{pF} = V_{pA} + V_{pn} - U_{pF}$$

$$M^{(\text{prior})} = \langle \Psi_f^{(-)} | \Delta V_{dA} | \Phi_i^{(+)} \rangle$$

$$\Delta V_{dA} = V_{pA} + V_{nA} - U_{dA}$$

- Exterior contributions: Post requires contributions further out; also an issue for resonance calculations

Peak cross section relative to full calculation



Assessing the R-matrix ideas Ib

1. Interior vs exterior contributions

$$M = M(0,a) + M(a,\infty)$$

This case:

- $^{90}\text{Zr}(\text{d},\text{p})$ for $E_{\text{d}}=11 \text{ MeV}$
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 - 1st excited state ($1/2^+$)
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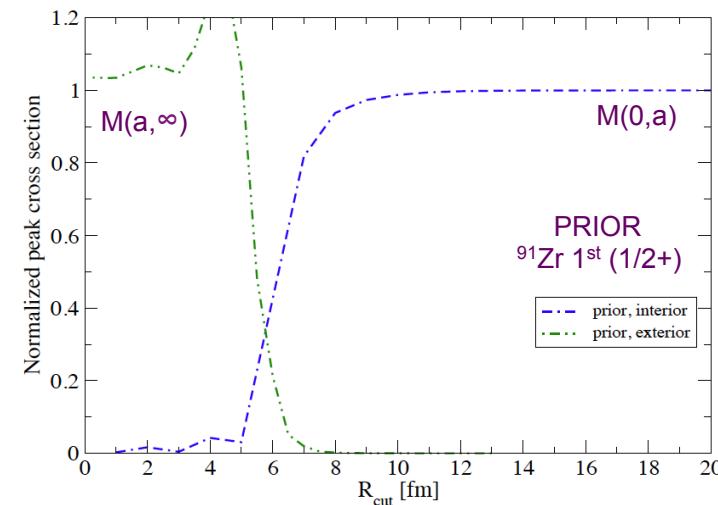
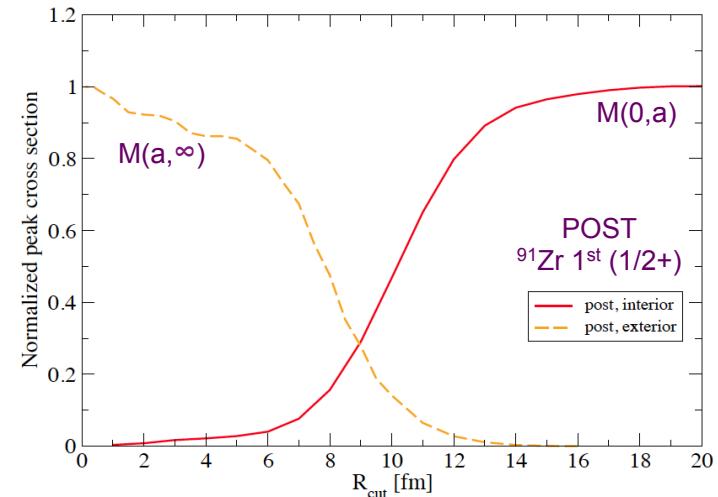
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Peak cross section relative to full calculation



Assessing the R-matrix ideas Ic

1. Interior vs exterior contributions

$$M = M(0,a) + M(a,\infty)$$

This case:

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Observations

- 'action is in the nuclear surface'
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$$M^{(\text{post})} = \langle \Phi_f^{(-)} | \Delta V_{pF} | \Psi_i^{(+)} \rangle$$

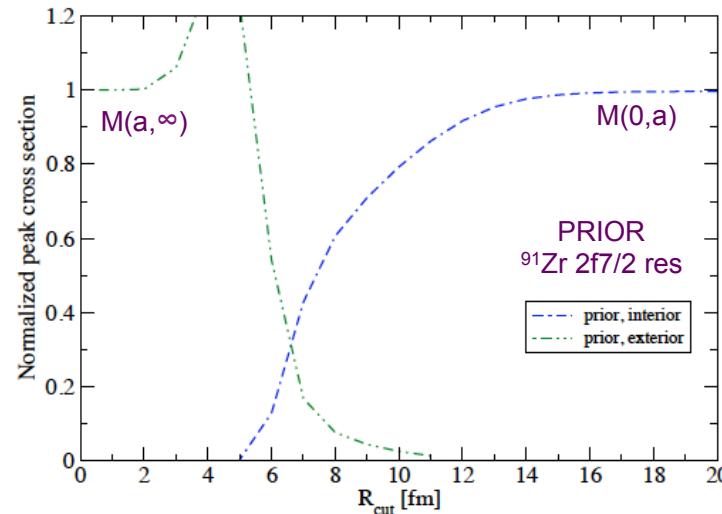
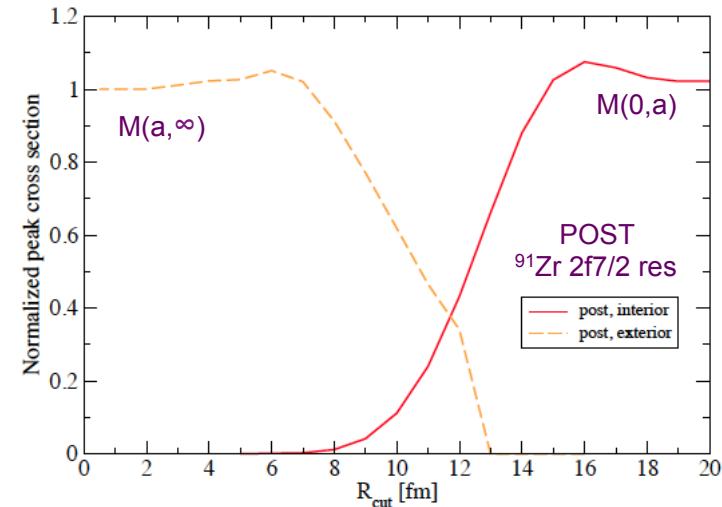
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Peak cross section relative to full calculation



Assessing the R-matrix ideas IIa

2. Surface contribution

$$M = M^{(\text{post})}(0,a) + M_{(\text{surf})}(a) + M^{(\text{prior})}(a,\infty)$$

This case:

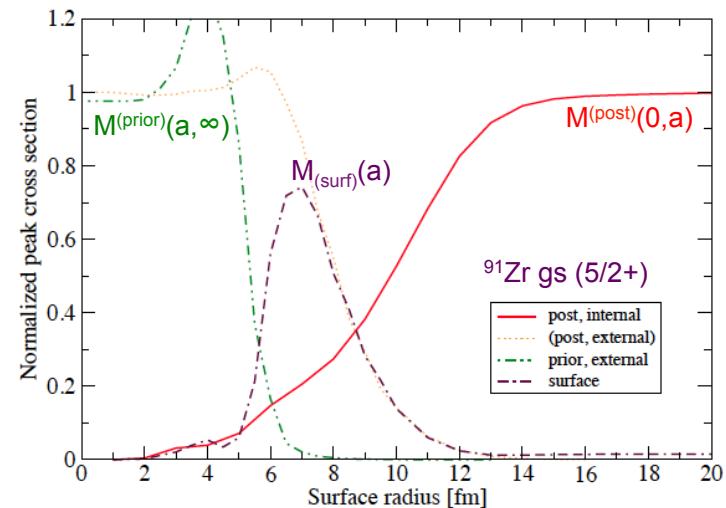
- ${}^{90}\text{Zr}(\text{d},\text{p})$ for $E_{\text{d}}=11 \text{ MeV}$
- ${}^{91}\text{Zr}$ gs ($5/2^+$)
- 1st excited state ($1/2^+$)
- $2f_{7/2}$ resonance



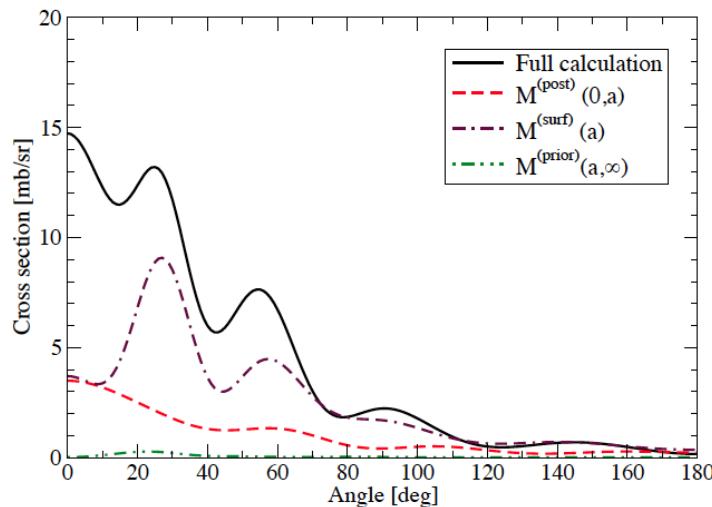
Observations

- Surface term indeed dominant 6-8 fm
- Small interior contributions → little dependence on model for interior
- Small exterior contributions → better convergence for resonance case
- Surface term does not produce the whole cross section, corrections required from internal/external

Peak cross section relative to full calculation



Angular differential cross sections



Assessing the R-matrix ideas IIb

2. Surface contribution

$$M = M^{(\text{post})}(0,a) + M_{(\text{surf})}(a) + M^{(\text{prior})}(a,\infty)$$

This case:

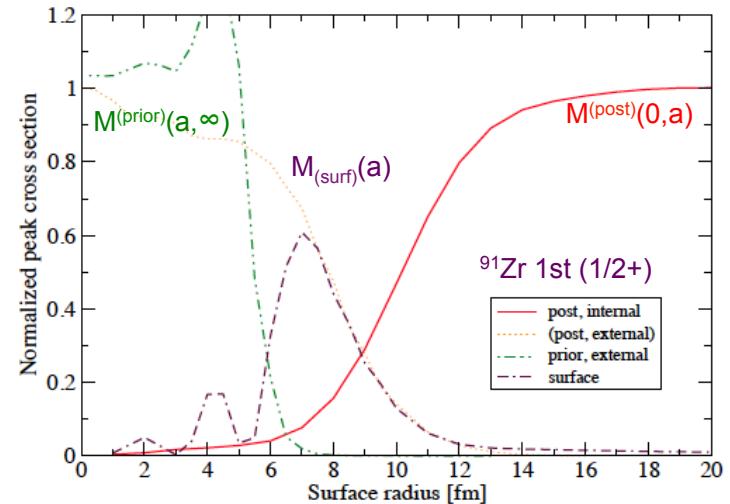
- ${}^{90}\text{Zr}(\text{d},\text{p})$ for $E_{\text{d}}=11 \text{ MeV}$
 - ${}^{91}\text{Zr}$ gs ($5/2^+$)
 - 1st excited state ($1/2^+$)
 - $2f_{7/2}$ resonance



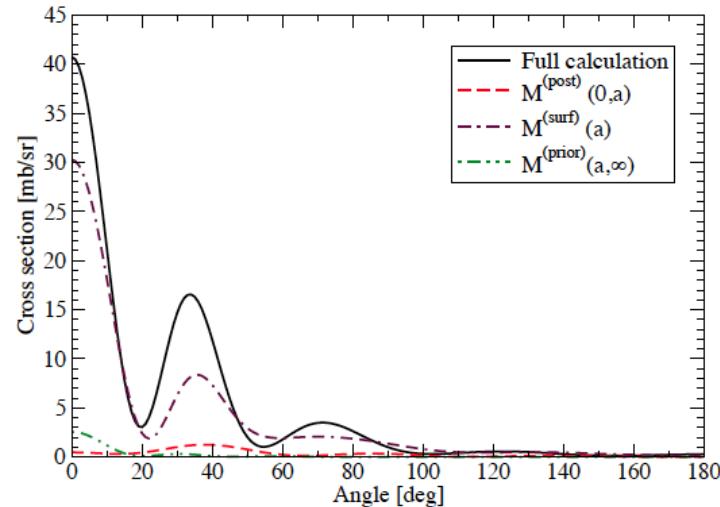
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Peak cross section relative to full calculation



Angular differential cross sections



Assessing the R-matrix ideas IIc

2. Surface contribution

$$M = M^{(\text{post})}(0,a) + M_{(\text{surf})}(a) + M^{(\text{prior})}(a,\infty)$$

This case:

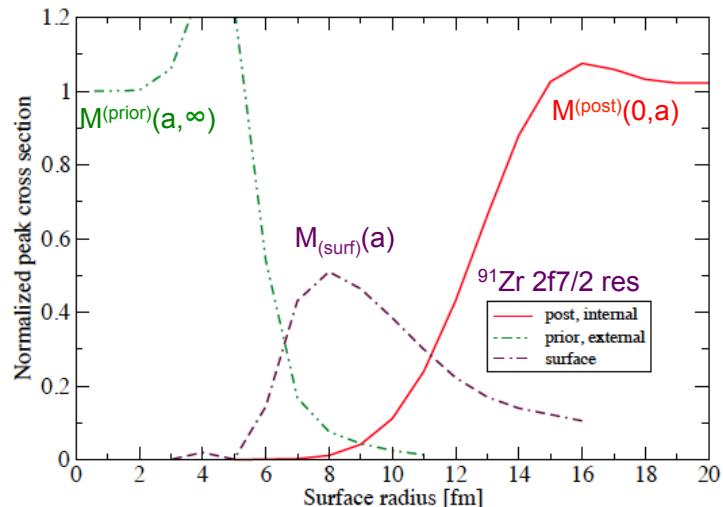
- ${}^{90}\text{Zr}(\text{d},\text{p})$ for $E_{\text{d}}=11$ MeV
 - ${}^{91}\text{Zr}$ gs ($5/2^+$)
 - 1st excited state ($1/2^+$)
 - $2f_{7/2}$ resonance



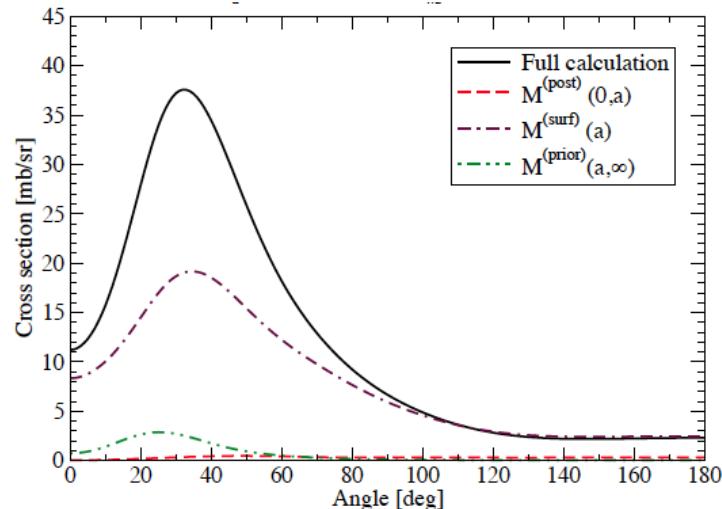
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Peak cross section relative to full calculation



Angular differential cross sections



Assessing the R-matrix ideas - ^{48}Ca

2. Surface contribution

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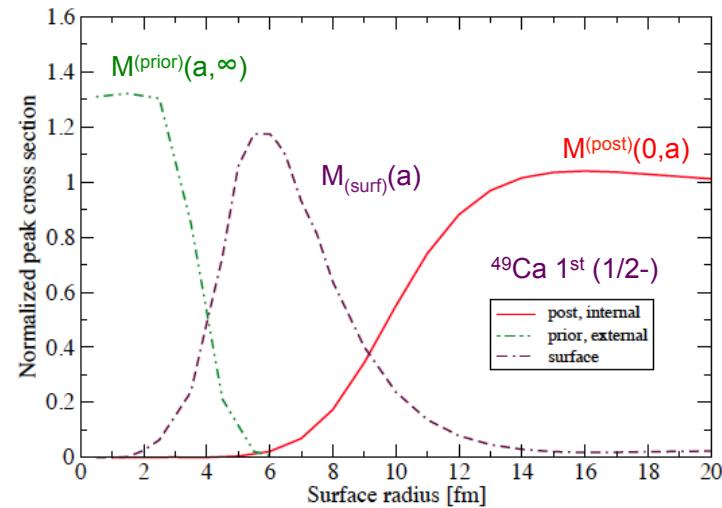
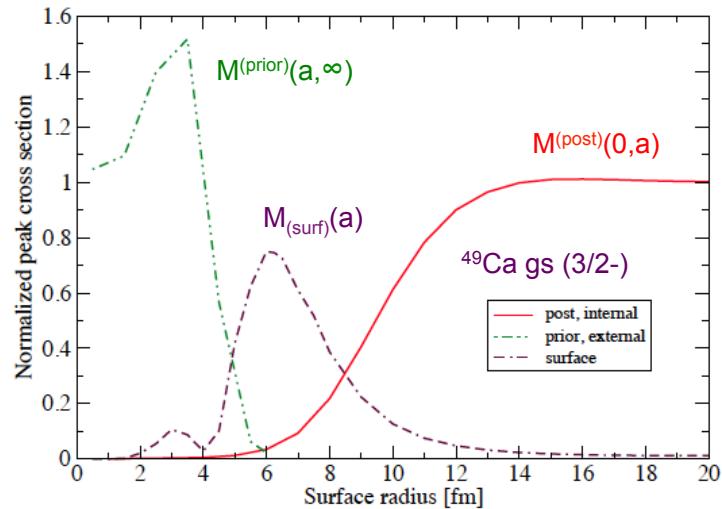
This case:

- $^{48}\text{Ca}(\text{d},\text{p})$ for $E_d=13$ MeV
 - ^{49}Ca gs ($3/2^-$)
 - 1st excited state ($1/2^-$)

Observations

- Surface term indeed dominant 5-7 fm
- Small interior contributions → little dependence on model for interior
- Small exterior contributions → better convergence for resonance case
- Surface term does not produce the whole cross section, corrections required from internal/external

Peak cross section relative to full calculation



Next: Extension of the formalism to include breakup

DWBA matrix element

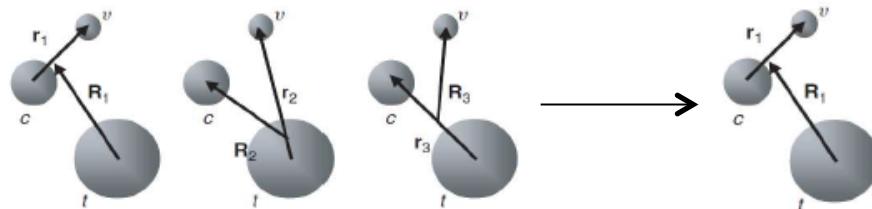
$$M^{(\text{post})} = M^{(\text{post})}(0, a) + M_{(\text{surf})}(a) + M^{(\text{prior})}(a, \infty)$$

CDCC matrix element

$$M^{(\text{post})} = M^{(\text{post})}(0, a) + M_{(\text{surf})}(a)$$
$$M^{(\text{prior})}(a, \infty) = 0 \text{ (is included in breakup)}$$

CDCC (Continuum-discretized coupled channels)

- Approximate treatment of 3-body problem
- Describes breakup of deuteron



- Successfully used for describing data
- Currently revisited via comparison with Fadeev
- To be studied in connection with R-matrix approach for (d,p) to resonances

Expectation for CDCC extension

- Simultaneous calculation of breakup and transfer cross sections
- Surface term dominant
- Good convergence

Conclusions

(d,p) reactions:

- Important for nuclear structure studies and astrophysics
- New experimental techniques for radioactive isotopes
- Improved theoretical descriptions required

Studying resonances with (d,p):

- Conceptual and practical problems have to be overcome

New formalism:

- Builds on ideas from successful R-matrix approach
- Separation into interior and exterior regions works formally well, surface term emerges as important contributor, can be expressed in terms of familiar R-matrix parameters
- Test cases show that the surface term is dominant, but other contributions may not be negligible
- Including breakup via CDCC removes exterior prior contribution, thus eliminates convergence problem for resonances

Further studies needed to clarify conditions where the surface formalism will work well.