

The Coulomb Problem in Momentum Space without Screening

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(The TORUS Collaboration)



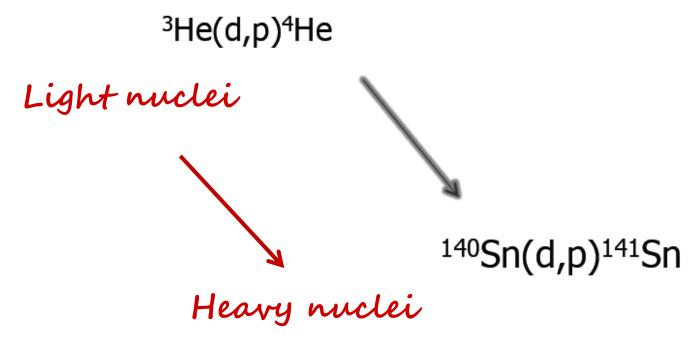




Physics Problem:

Nuclear Reactions dominated by few degrees of freedom

Reactions: Elastic Scattering, Breakup & Transfer



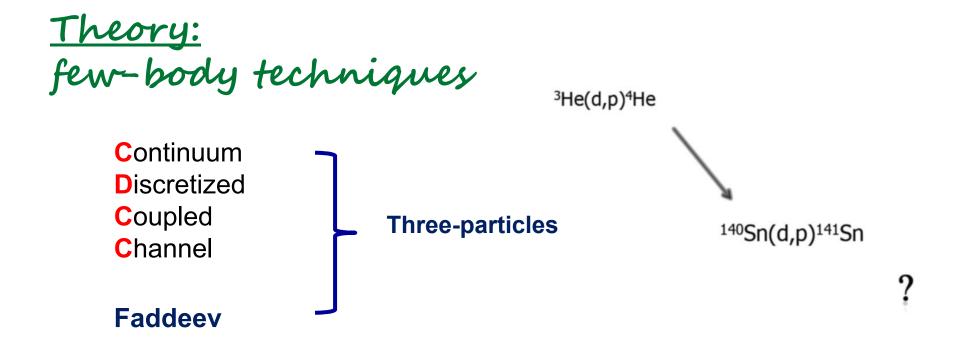




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Reduce Many-Body to Few-Body Problem



Task:

- Isolate important degrees of freedom in a reaction
- Keep track of important channels
- Connect back to the many-body problem

Hamiltonian for effective few-body poblem:

Three-Body Problem





(d,p) Reactions as three-body problem



Deltuva and Fonseca, Phys. Rev. C79, 014606 (2009)

Applied Faddeev AGS Equations to ¹²C(d,p)¹³C

Elastic, breakup, rearrangement channels are included and fully coupled

(compared to e.g. CDCC calculations)



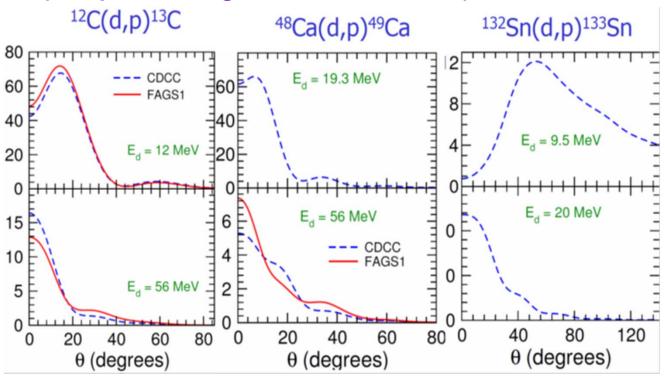


(d,p) Reactions as three-body problem



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Issue: current
momentum
space
implementation
of Coulomb
interaction
(screening) does
not converge for
Z ≥ 20

Courtesy: F.M. Nunes







A.M. Mukhamedzhanov, V. Eremenko and A.I. Sattarov, Phys.Rev. C86 (2012) 034001

Solve Faddeev equations in Coulomb basis (no screening)

Scattering: Faddeev equations best solved in momentum space



Very nasty! Oscillatory singular at p=q



Matrix elements in Coulomb basis:

Up to now

scattering

 not directly solved Example: plane wave basis: $V(p',p) \equiv \langle p' \mid V \mid p \rangle$

Coulomb basis: 2 singularities, for p'=p: "pinch" singularity

• Indirect: Chinn, CE. Thaler. PRC44, 1569 (1991) for p+A



Work with separable functions: $V(p',p) \equiv \sum g(p') \lambda g(p)$

Can we handle this?







First Test in Two-Body System



Separable t-matrix derived from p+A optical potential (generalized EST scheme)

$$t_l(E) = \sum_{i,j} u |f_{l,k_{E_i}}\rangle \tau_{ij}(E) \langle f_{l,k_{E_j}}^* | u \rangle$$

Nuclear matrix elements $\langle p|t_l(E)|p'\rangle$



$$\langle p|u|f_{l,k_E}\rangle = t_l(p, k_E; E_{k_E}) \equiv u_l(p)$$
$$\langle f_{l,k_E}^*|u|p'\rangle = t_l(p', k_E; E_{k_E}) \equiv u_l(p')$$

Coulomb distorted nuclear matrix element



$$\langle \psi_{l,p}^{C} | u | f_{l,k_{E}} \rangle = \int_{0}^{\infty} \frac{dq \ q^{2}}{2\pi^{2}} \ u_{l}(q) \psi_{l,p}^{C}(q)^{*} \equiv u_{l}^{C}(p)$$

$$\langle f_{l,k_{E}}^{*} | u | \psi_{l,p}^{C} \rangle = \int_{0}^{\infty} \frac{dq \ q^{2}}{2\pi^{2}} \ u_{l}(q) \ \psi_{l,p}^{C}(q) \equiv u_{l}^{C}(p)^{\dagger}$$



 $\psi_{n,l}^{C}(p)$ is the Coulomb scattering wave function





Challenge I: momentum space Coulomb functions

General:
$$\psi_{\mathbf{q},\eta}^{C(+)}(\mathbf{p}) = \lim_{\gamma \to +0} \int d^3\mathbf{r} \ e^{-i\mathbf{p}\mathbf{r}-\gamma r} \ \psi_{\mathbf{q},\eta}^{C(+)}(r)$$

FT: A. Chan, MS thesis **U. Waterloo (2007)**

$$= -4\pi e^{-\pi\eta/2} \Gamma(1+i\eta) \lim_{\gamma \to +0} \frac{d}{d\gamma} \left\{ \frac{[p^2 - (q+i\gamma)^2]^{i\eta}}{[|\mathbf{p} - \mathbf{q}|^2 + \gamma^2]^{1+i\eta}} \right\}$$

Partial wave decomposition (Mukhamedzanov, Dolinskii) (1966)

$$\psi_{\mathbf{q},\eta}^{C(+)}(\mathbf{p}) \equiv \sum_{l=0}^{\infty} (2l+1)\psi_{l,q,\eta}^{C}(p)P_{l}(\hat{\mathbf{p}} \cdot \hat{\mathbf{q}}), \qquad \psi_{l,q,\eta}^{C}(p) = \frac{1}{2} \int_{-1}^{1} dz \ P_{l}(z)\psi_{\mathbf{q},\eta}^{C(+)}(\mathbf{p}),$$

$$\frac{1}{2} \int_{-1}^{1} dz \ P_{l}(z)(\zeta - z)^{-1-i\eta} = \frac{e^{\pi\eta}}{\Gamma(1+i\eta)} (\zeta^{2} - 1)^{-i\eta/2} Q_{l}^{i\eta}(\zeta)$$

$$\zeta \equiv \frac{p^2 + q^2 + \gamma^2}{2pq},$$

Essential:

 $Q_{I}^{i\eta}(\zeta)$ has different representations depending on ζ





$Q_l^{i\eta}(\zeta)$ has different representations in terms of the hypergeometric function $_2{\sf F}_1$ (a;b;c;z) depending on ζ

 ζ large enough (p and q different) —— "regular" representation

$$Q_l^{i\eta}(\zeta) = \frac{e^{-\pi\eta}\Gamma(l+i\eta+1)\Gamma(1/2)}{2^{l+1}\Gamma(l+3/2)}(\zeta^2 - 1)^{i\eta/2}\zeta^{-l-i\eta-1} \times {}_2F_1\left(\frac{l+i\eta+2}{2}, \frac{l+i\eta+1}{2}; l + \frac{3}{2}; \frac{1}{\zeta^2}\right)$$

 $\zeta \approx 1 \ (p \approx q)$ "pole-proximity" representation

$$Q_l^{i\eta}(\zeta) = \frac{1}{2}e^{-\pi\eta} \left\{ \Gamma(i\eta) \left(\frac{\zeta+1}{\zeta-1} \right)^{i\eta/2} {}_2F_1 \left(-l, l+1; 1-i\eta; \frac{1-\zeta}{2} \right) + \frac{\Gamma(-i\eta)\Gamma(l+i\eta+1)}{\Gamma(l-i\eta+1)} \left(\frac{\zeta-1}{\zeta+1} \right)^{i\eta/2} {}_2F_1 \left(-l, l+1; 1+i\eta; \frac{1-\zeta}{2} \right) \right\}$$





Partial-wave momentum space Coulomb functions

"regular" representation

$$\psi_{l,q}^{C}(p) = -\frac{4\pi\eta e^{-\pi\eta/2}q(pq)^{l}}{(p^{2} + q^{2})^{1+l+i\eta}} \left[\frac{\Gamma(1+l+i\eta)}{(1/2)_{l+1}} \right] \times {}_{2}F_{1}\left(\frac{2+l+i\eta}{2}, \frac{1+l+i\eta}{2}; l+3/2; \frac{4q^{2}p^{2}}{(p^{2}+q^{2})^{2}} \right) \times \lim_{\gamma \to 0} \left[p^{2} - (q+i\gamma)^{2} \right]^{-1+i\eta}$$

"pole-proximity" representation:

$$\psi_{l,q}^C(p) = -\frac{2\pi}{p} \exp(-\pi\eta/2 + i\sigma_l) \left[\frac{(p+q)^2}{4pq} \right]^l \lim_{\gamma \to 0} 2 \operatorname{Sm} \mathcal{D}.$$

$$\mathcal{D} \equiv \frac{\Gamma(1+i\eta)e^{-i\sigma_l}(p+q)^{-1+i\eta}}{(p-q+i\gamma)^{1+i\eta}} \,_{2}F_{1}\left(-l,-l-i\eta;1-i\eta;\frac{(p-q)^{2}}{(p+q)^{2}}\right)$$

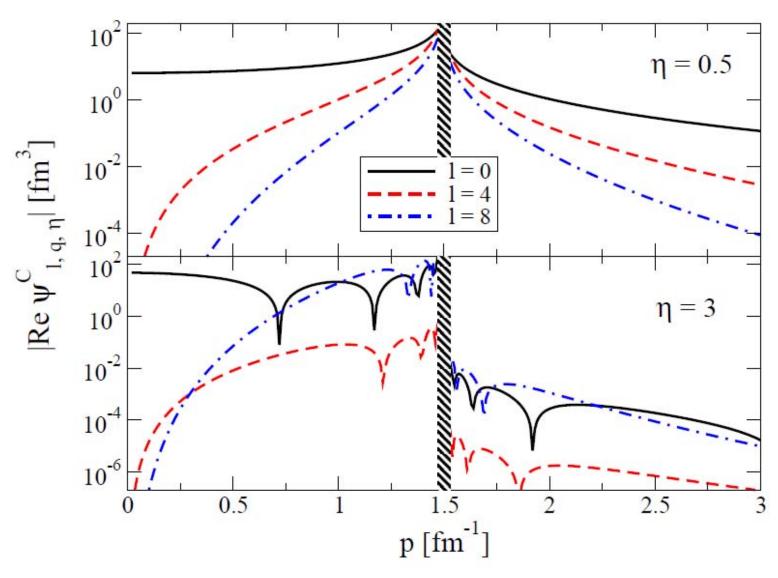


Oscillatory singularity for p-q



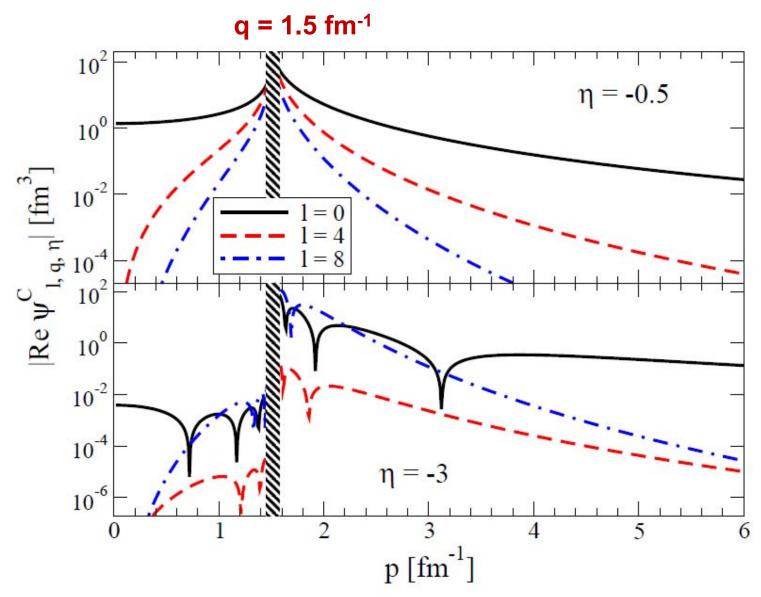












Work in progress: publish code in CPC





Challenge II:

Matrix elements with Coulomb basis functions

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"oscillatory" singularity at $q=p:\lim_{\gamma o +0} rac{1}{(q-p+i\gamma)^{1+i\eta}}$ cs + astronomy

Gel'fand-Shilov Regularization:

Generalization of Principal value regularization Idea: reduce value of integrand near singularity

$$\int_{-\Delta}^{\Delta} dx \frac{\varphi(x)}{x^{1+i\eta}} = \int_{-\Delta}^{\Delta} dx \frac{\varphi(x) - \varphi(0) - \varphi'(0)x}{x^{1+i\eta}}$$
 simplified

Reduce integrand around pole by subtracting 2 terms of the Laurent series

$$-\frac{i\varphi(0)}{\eta}[\Delta^{-i\eta}-(\Delta)^{-i\eta}]+\dots$$





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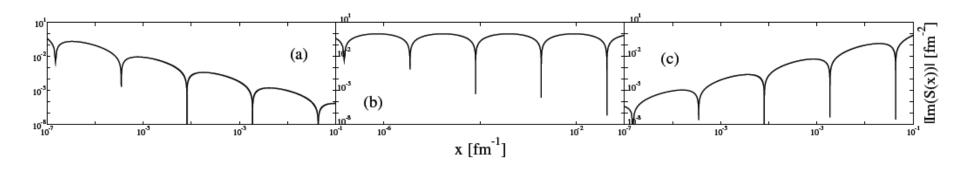


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➤ Reduce integrand around pole by subtracting 2 terms of the Laurent series

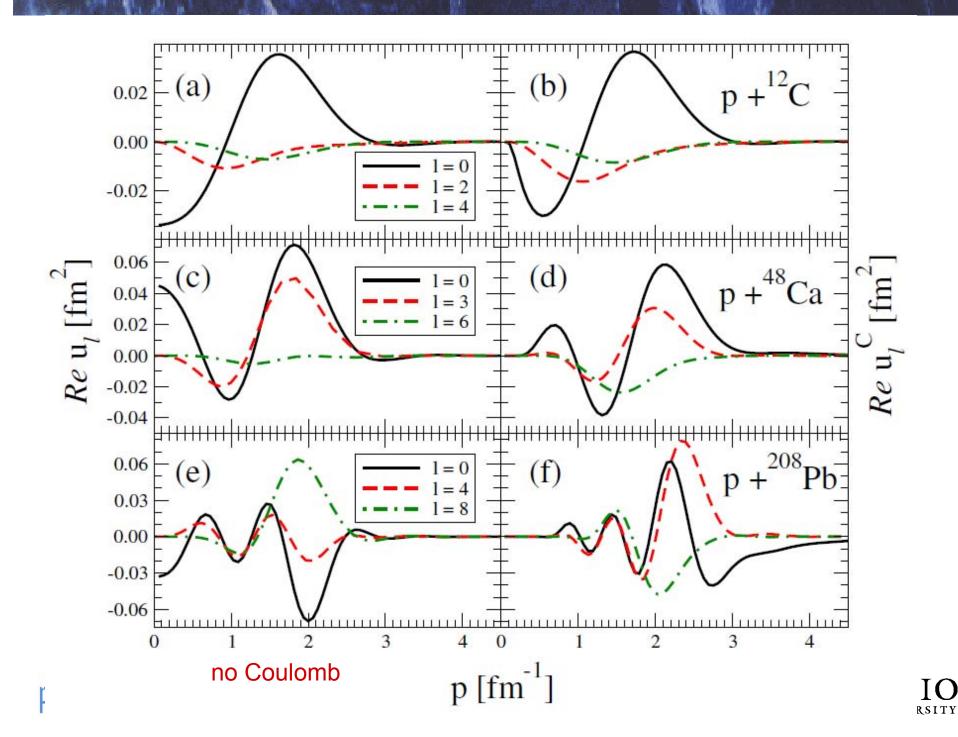
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I. M. Gel'fand and G. E. Shilov. "Generalized Functions". Vol. 1.
Academic Press, New York and London. 1964.



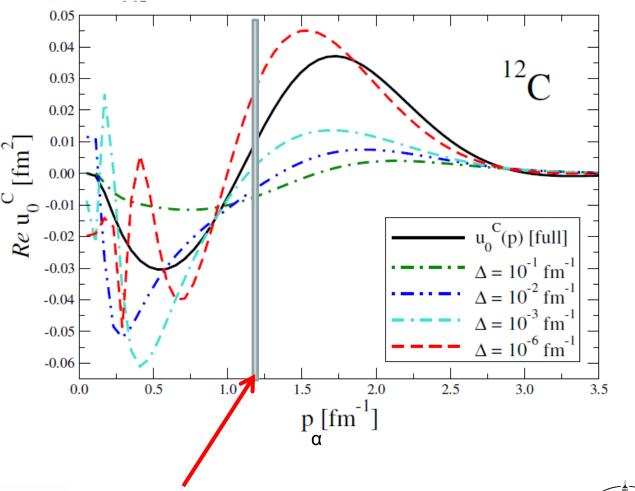




$p + {}^{12}C$



$$u_l^C(p_\alpha) = \underbrace{\int_0^{p_\alpha - \Delta} \frac{dp}{2\pi^2} \, p^2 u_l(p) \psi_{p_\alpha l}^C(p)}_{l} + \underbrace{\int_{p_\alpha - \Delta}^{p_\alpha + \Delta} \dots + \int_{p_\alpha + \Delta}^{\infty} \dots}_{l}$$



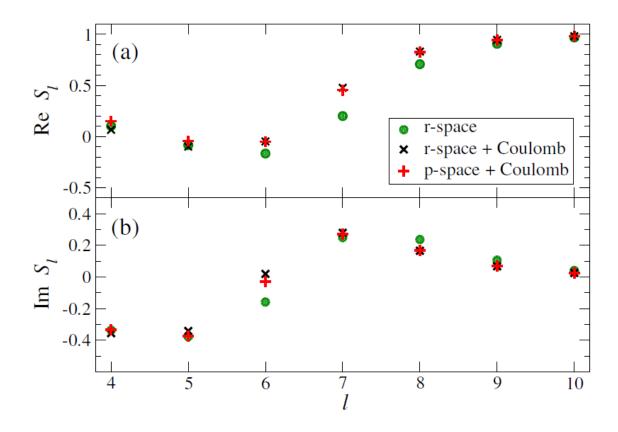
physics + astronomy

Fixed p_{α}



First Physics Check:

p + ⁴⁸Ca



Selected partial wave S-matrix elements S_{l+1} for p+48Ca (CH89 optical potential) with Coulomb distorted n+48Ca formfactors

Method not designed for two-body scattering!





Summary & Outlook

Faddeev-AGS framework in Coulomb basis passed first test!

- Momentum space nuclear form factors obtained in a Coulomb distorted basis for high charges for the first time.
- \succ "Oscillatory singularity" of $\psi_{q,l}{}^c(p)$ at p \rightarrow q successfully regularized.
- ightharpoonup Algorithms to compute $\psi_{q,l}{}^c(p)$ and overlap integrals successfully implemented



Near Future:

Implementation of Faddeev-AGS equations in Coulomb basis







TORUS: Theory of Reactions for Ustable iSotopes

A Topical Collaboration for Nuclear Theory

http://www.reactiontheory.org/



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