



INPP

INSTITUTE OF NUCLEAR & PARTICLE PHYSICS

@OHIO UNIVERSITY

Towards (d,p) Reactions with Heavy Nuclei in a Faddeev Description

Ch. Elster

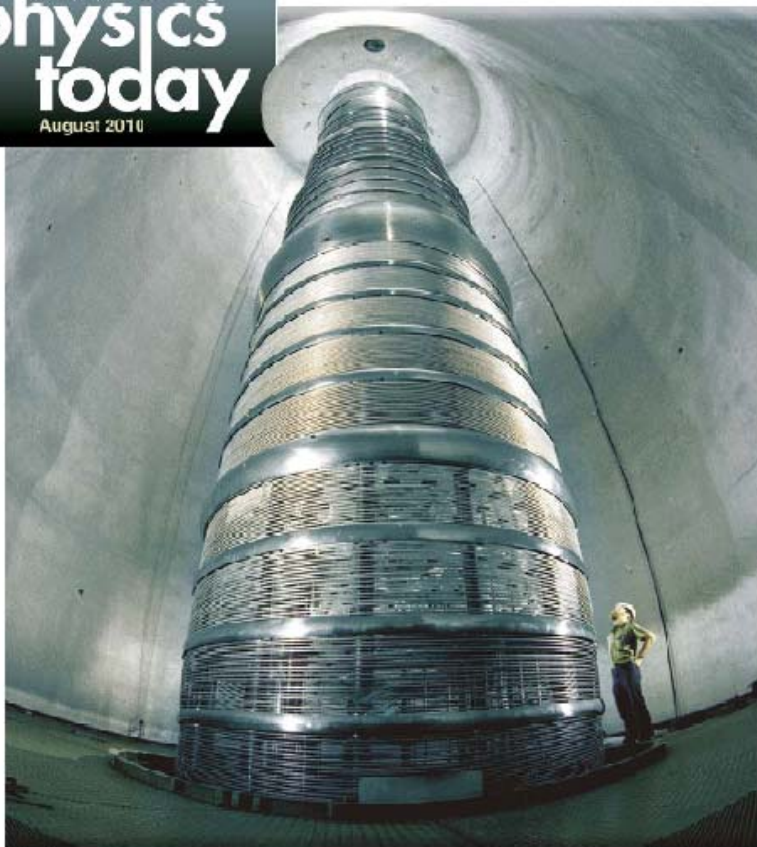
TORUS Collaboration

International Workshop on
Nuclear Dynamics with Effective Field Theories
Ruhr-Universität Bochum, July 1 - 3, 2013

Dedicated to the memory of late Prof. Walter Glöckle

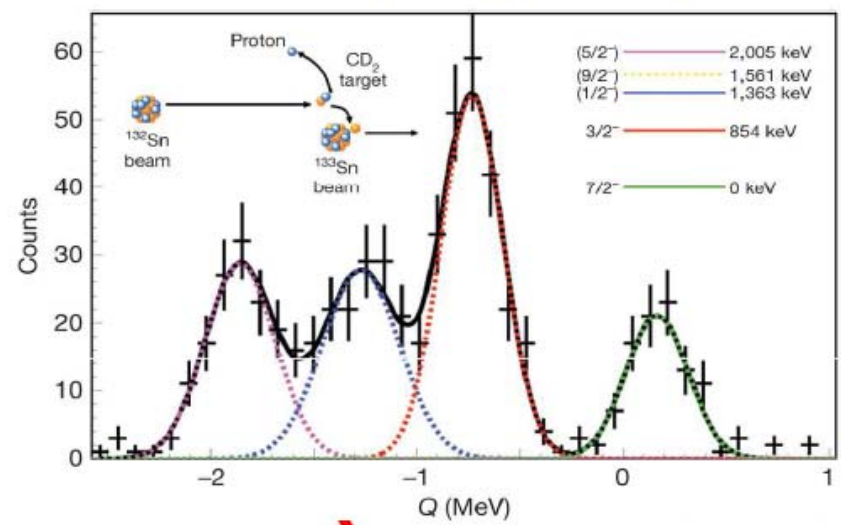
(d,p) Reactions: Tool to study structure

www.physics.ohio-state.edu
physics today
 August 2010

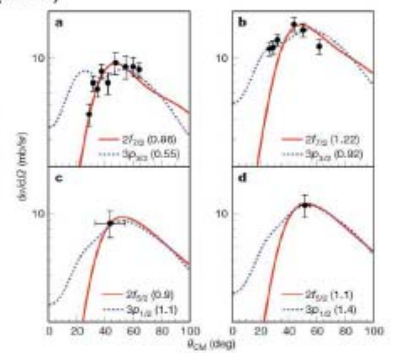


Doubly magic shell game

$d(^{132}\text{Sn}, ^{133}\text{Sn})p @ 5 \text{ MeV/u}$

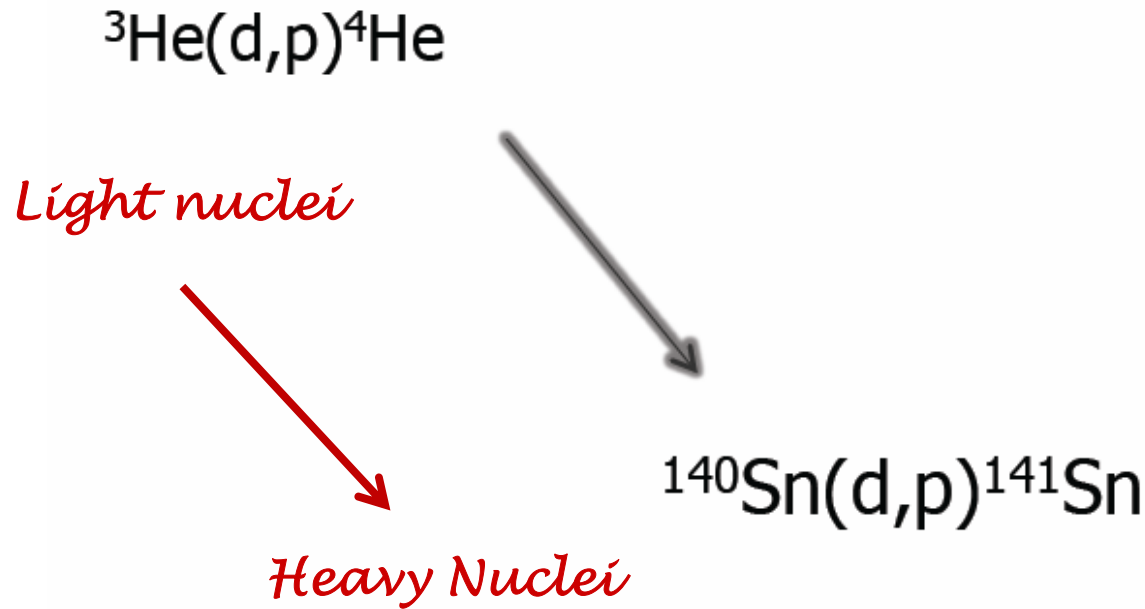


connection
 to r-process
 $^{132}\text{Sn}(n,\gamma)^{133}\text{Sn}$



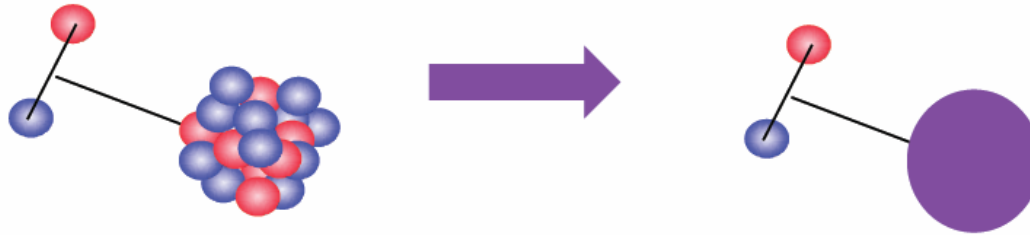
[K. Jones et al, Nature 465 (2010) 454]

What Reactions are we interested in?



Reactions: Elastic Scattering, Breakup, Transfer

Reduce Many-Body to Few-Body Problem



Task:

- Isolate important degrees of freedom in a reaction
- Keep track of important channels
- Connect back to the many-body problem

Hamiltonian for effective few-body problem:

$$H = H_0 + V_{np} + V_{nA} + V_{pA}$$

np interaction

Optical potentials p+A and n+A

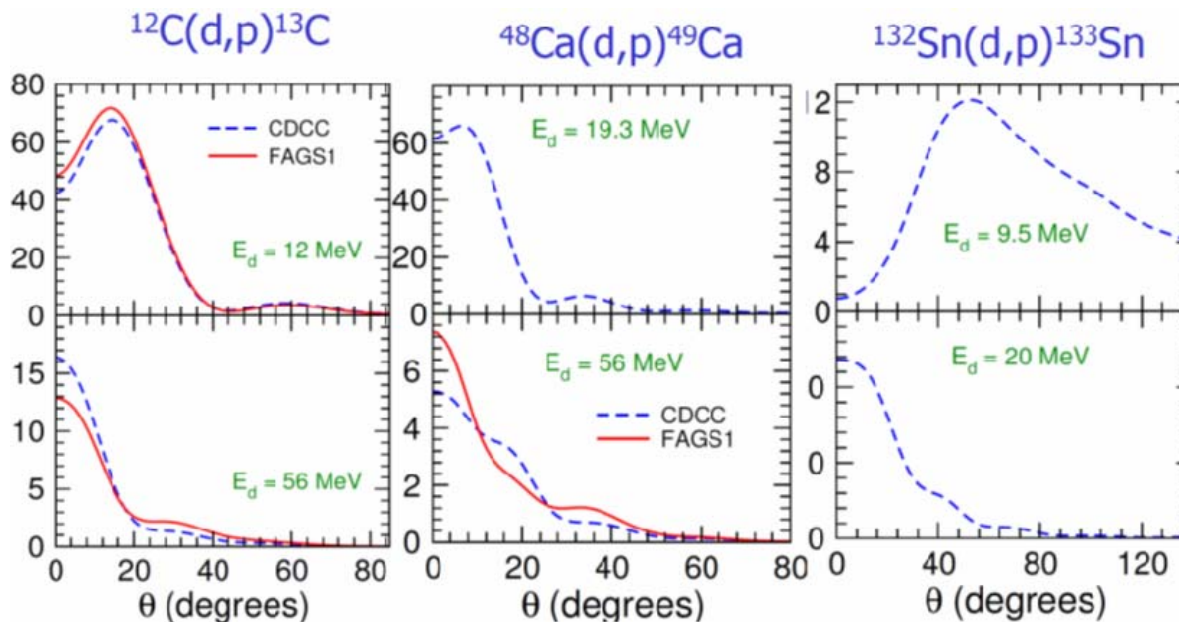
Three-Body Problem

(d,p) Reactions as three-body problem



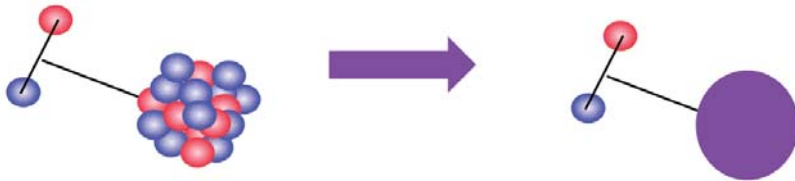
Deltuva and Fonseca, Phys. Rev. C **79**, 014606 (2009).

Elastic, breakup, rearrangement channels are included and fully coupled (compared to e.g. CDCC calculations)



Issue: current momentum space implementation of Coulomb interaction (shielding) does **not** converge for $Z \geq 20$

(d,p) Reactions as three-body problem



Deltuva and Fonseca, Phys. Rev. C **79**, 014606 (2009).

Elastic, breakup, rearrangement channels are included and fully coupled (compared to e.g. CDCC calculations)

Issue: traditional Faddeev formulation does **not** contain target excitations

Probably especially important for reactions with exotic nuclei

Proposed Plan to address issues:

A.M. Mukhamedzhanov, V.Eremenko and A.I. Sattarov, Phys.Rev. C **86** (2012) 034001

TORUS Collaboration

www.reactiontheory.org



A.M. Mukhamedzhanov, V.Eremenko and A.I. Sattarov,
Phys.Rev. C86 (2012) 034001

Generalized Faddeev formulation of (d,p) reactions with

- (a) explicit inclusion of target excitations
- (b) explicit inclusion of the Coulomb interaction

Faddeev formulation \rightarrow momentum space

Suggestions:



Target excitations:

Including specific excited states \rightarrow Formulation with separable interactions



Explicit inclusion of Coulomb interaction:

Formulation of Faddeev equations in Coulomb basis instead of plane wave basis

Both need prep work!



Hamiltonian: $H = H_0 + V_{np} + V_{nA} + V_{pA}$

V_{np} : NN interaction – separable representations exist



V_{nA} : Optical potential

Phenomenological optical potentials fitted to data from ^{12}C to ^{208}Pb
 given in coordinate-space
 parameterized in terms of Woods-Saxon functions

$$U_{nucl}(r) = V(r) + i[W(r) + W_s(r)] + V_{ls}(r) \mathbf{l} \cdot \boldsymbol{\sigma}$$

$$V(r) = -V_r f_{ws}(r, R_0, a_0)$$

$$W(r) = -W_v f_{ws}(r, R_w, a_w)$$

$$W_s(r) = -W_s(-4a_w) f'_{ws}(r, R_w, a_w)$$

$$V_{ls}(r) = -(V_{so} + iW_{so})(-2)g_{ws}(r, R_{so}, a_{so}),$$

$$f_{ws}(r, R, a) = \frac{1}{1 + \exp(\frac{r-R}{a})}$$

$$f'_{ws}(r, R, a) = \frac{d}{dr} f_{ws}(r, R, a)$$

$$g_{ws}(r, R, a) = f'_{ws}(r, R, a)/r$$

Not useful in this form

Woods-Saxon functions have a semi-analytic Fourier transform:

(fast converging series expansion)

Central term:

$$\bar{V}(\mathbf{q}) = \frac{V_r}{\pi^2} \left\{ \frac{\pi a_0 e^{-\pi a_0 q}}{q (1 - e^{-2\pi a_0 q})^2} [R_0 (1 - e^{-2\pi a_0 q}) \cos(qR_0) - \pi a_0 (1 + e^{-2\pi a_0 q}) \sin(qR_0)] - a_0^3 e^{-\frac{R_0}{a_0}} \left[\frac{1}{(1 + a_0^2 q^2)^2} - \frac{2e^{-\frac{R_0}{a_0}}}{(4 + a_0^2 q^2)^2} \right] \right\}$$

Surface term:

$$\bar{W}_s(\mathbf{q}) = -4a_w \frac{W_s}{\pi^2} \left\{ \frac{\pi a_w e^{-\pi a_w q}}{(1 - e^{-2\pi a_w q})^2} \left[\left(\pi a_w (1 + e^{-2\pi a_w q}) - \frac{1}{q} (1 - e^{-2\pi a_w q}) \right) \cos(qR_w) + R_w (1 - e^{-2\pi a_w q}) \sin(qR_w) \right] + a^2 e^{-R_w/a_w} \left[\frac{1}{(1 + a_w^2 q^2)^2} - \frac{4e^{-R_w/a_w}}{(4 + a_w^2 q^2)^2} \right] \right\}.$$

Separabilization: Method of Ernst-Shakin-Thaler

BUT: Needs to be generalized for complex potentials

Example: Rank-1 separable potential:

$$V = \frac{v |\Psi_{k_E}^{(-)}\rangle \langle \Psi_{k_E}^{(+)}| v}{\langle \Psi_{k_E}^{(-)}| v | \Psi_{k_E}^{(+)}\rangle},$$

Definition with In-state necessary to fulfill reciprocity theorem

$$\mathcal{K}V\mathcal{K}^{-1} = V^\dagger \quad \mathcal{K} \text{ is the time-reversal operator.}$$

$$t(p', p, E) = \frac{t(p', k_E, E) t(p, k_E, E)}{\langle \Psi_{k_E}^{(-)}| v (1 - g_0(E)v) | \Psi_{k_E}^{(+)}\rangle} \equiv t(p', k_E, E) \tau(E) t(p, k_E, E)$$

$$\begin{aligned} \tau(E)^{-1} &= t(k_E, k_E, E_{k_E}) \\ &+ 2\mu \left[\mathcal{P} \int dp p^2 \frac{t(p, k_E, E_{k_E}) t(p, k_E, E_{k_E})}{k_E^2 - p^2} - \mathcal{P} \int dp p^2 \frac{t(p, k_E, E_{k_E}) t(p, k_E, E_{k_E})}{k_0^2 - p^2} \right] \\ &+ i\pi\mu \left[k_0 t(k_0, k_E, E_{k_E}) t(k_0, k_E, E_{k_E}) - k_E t(k_E, k_E, E_{k_E}) t(k_E, k_E, E_{k_E}) \right] \end{aligned}$$

Generalization to arbitrary rank

$$V = \sum_{i,j} v |\Psi_i^{(-)}\rangle \langle \Psi_i^{(-)} | M | \Psi_j^{(+)}\rangle \langle \Psi_j^{(+)} | v.$$

with

$$\delta_{ik} = \sum_j \langle \Psi_i^{(-)} | M | \Psi_j^{(+)}\rangle \langle \Psi_j^{(+)} | v | \Psi_k^{(-)}\rangle = \sum_j \langle \Psi_i^{(+)} | v | \Psi_j^{(-)}\rangle \langle \Psi_j^{(-)} | M | \Psi_k^{(+)}\rangle$$

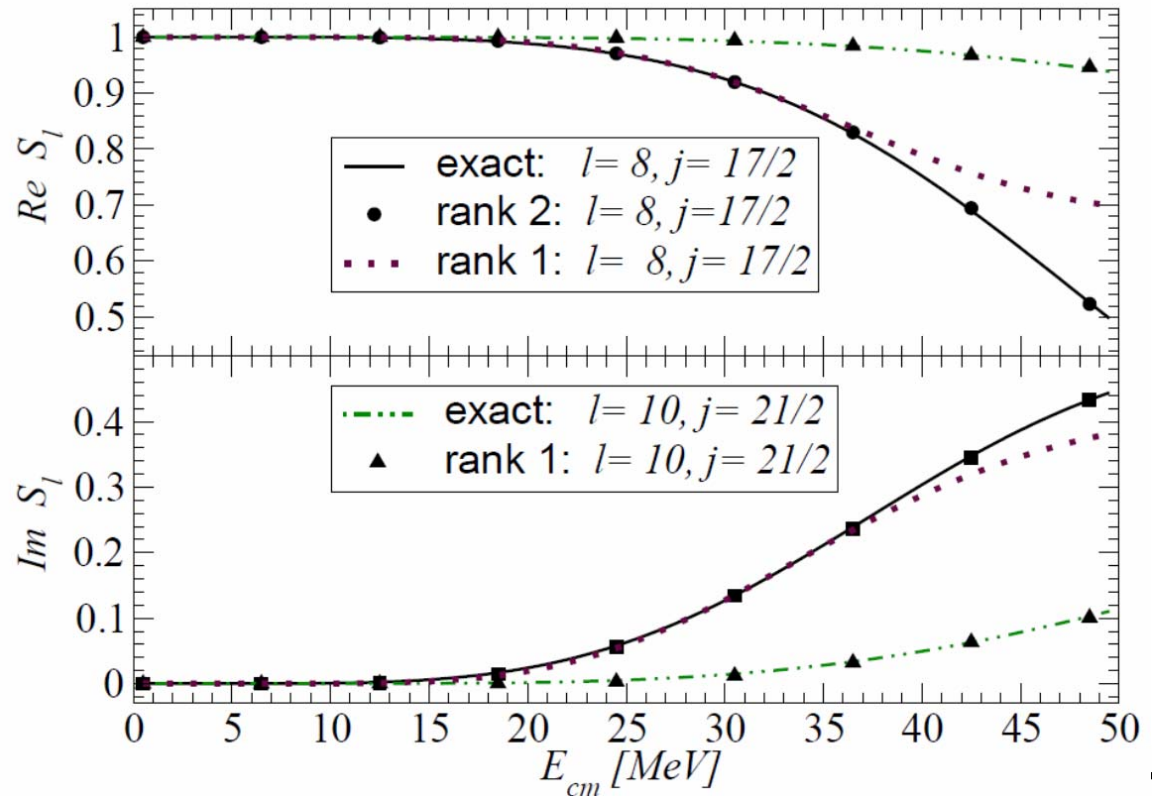
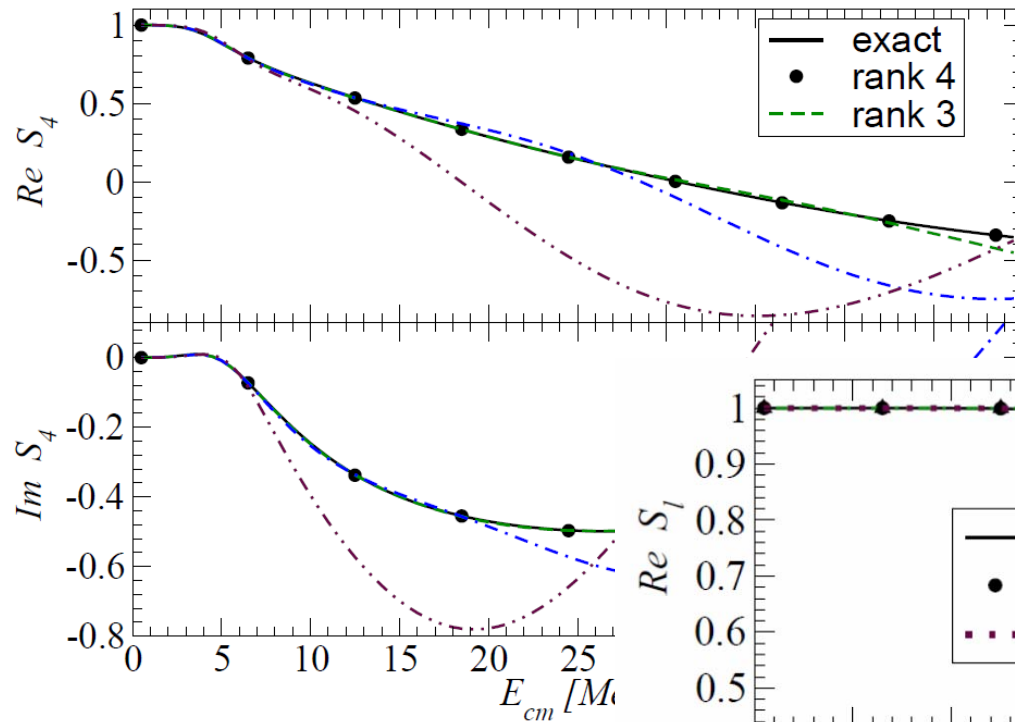
t-matrix

$$\hat{t}(E) = \sum_{i,j} v |\Psi_i^{(-)}\rangle \tau_{ij}(E) \langle \Psi_j^{(+)} | v$$

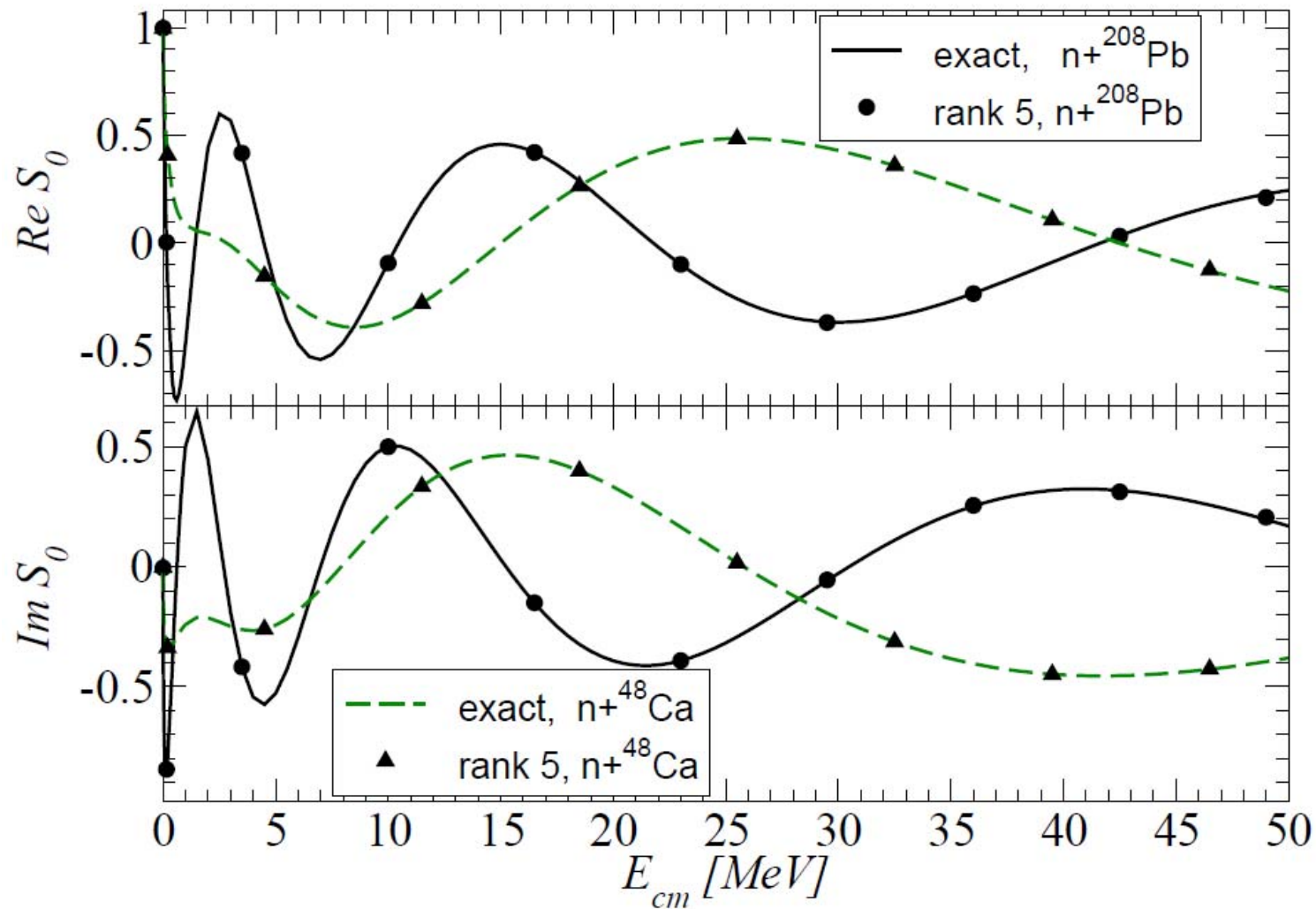
$$\sum_j \tau_{ij}(E) \underbrace{\langle \Psi_j^{(-)} | v - v g_0(E) v | \Psi_k^{(+)}\rangle}_{\text{Compute and solve system of linear equations}} = \delta_{ik}$$

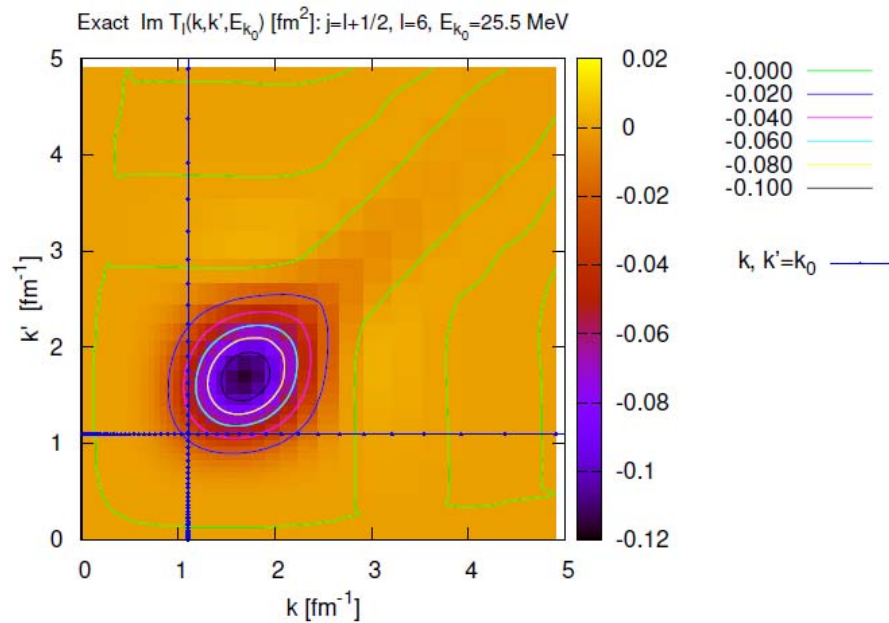
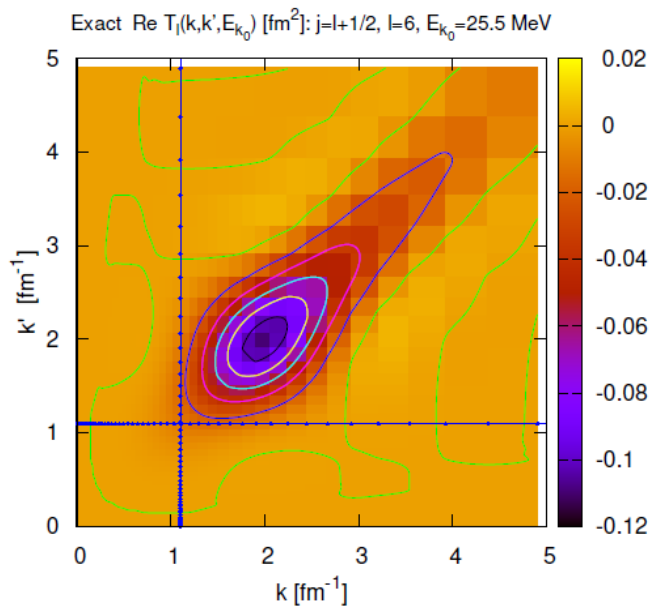
Compute and solve system of linear equations

$n + {}^{48}\text{Ca} : l=4, j=9/2$

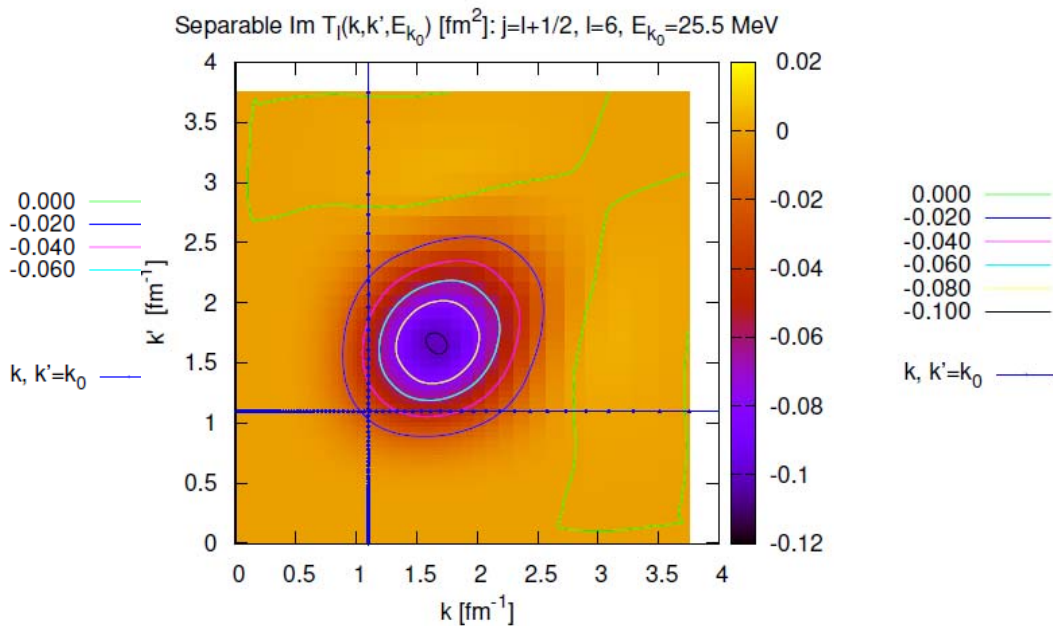
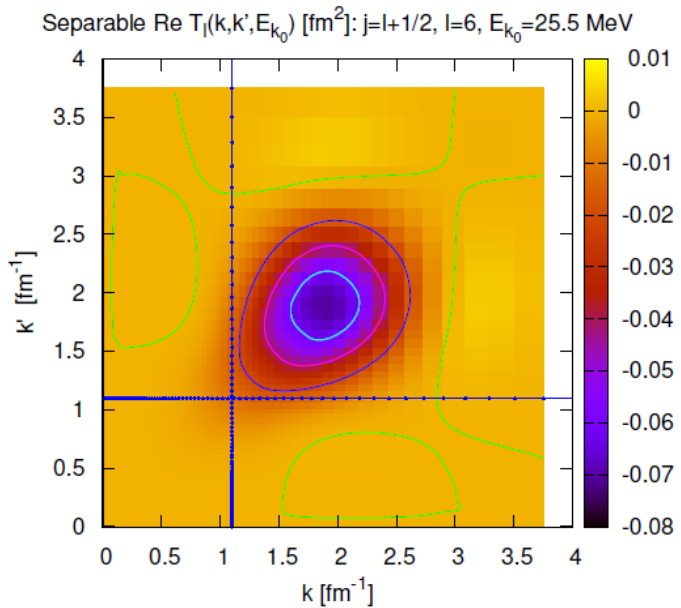


$n + {}^{48}\text{Ca}$ and $n + {}^{208}\text{Pb}$: $l=0$





$n + {}^{48}\text{Ca}$



Separabilization of Optical Potentials

Method of Ernst-Shakin-Thaler

- In momentum space
- **Generalized to non-Hermitian potentials**
- Universal Rank-5 representation: ^{12}C to ^{208}Pb
- Allows to use all phenomenological Woods-Saxon based optical potentials in our q-space Faddeev equations
- EST projects out high-momentum off-shell components.
- (spawned project at MSU to compare with SRG methods)
- Next: coupling to excited states

Suggestion: A.M. Mukhamedzhanov, V.Eremenko and A.I. Sattarov,
Phys.Rev. C86 (2012) 034001

Solve Faddeev equations in Coulomb basis

This implies calculation of integrals like

$$Z_l^{SC}(p, p_\alpha) = \int \frac{dp' p'^2}{2\pi^2} V_l(p, p') \psi_{p_\alpha}^C(p')$$

If $V_l(p, p')$ separable, then integral contains a smooth function and $\psi_{p_\alpha}^C(p')$

Very nasty!

Coulomb wave function in momentum space and pw decomposition

Ref.: AMM *et al.*, Soviet Journal of Physics 3, 180 (1966).

Pole at $p_\alpha = p'$

Some insights on momentum space Coulomb wave functions:
There are two representations: pole and non-pole regions

$$\psi_{p_\alpha l}^C(p') = \frac{-4\pi e^{-\eta_\alpha \pi/2}}{p'} \left(\frac{(p' + p_\alpha)^2 + \gamma^2}{4p' p_\alpha} \right)^l \times \Gamma(1 + i\eta_\alpha) e^{i\alpha l}$$

$$\times \lim_{\gamma \rightarrow +0} \text{Im} \left\{ \left[e^{-i\alpha l} \frac{(p' + p_\alpha + i\gamma)^{i\eta_\alpha - 1}}{(p' - p_\alpha + i\gamma)^{i\eta_\alpha + 1}} \right. \right.$$

$$\left. \left. \times {}_2F_1 \left(-l, -l - i\eta_\alpha; 1 - i\eta_\alpha; \frac{(p' - p_\alpha)^2 + \gamma^2}{(p' + p_\alpha)^2 + \gamma^2} \right) \right] + \gamma \left[\dots \right] \right\}$$

$\psi_{p_\alpha l}^C(p')$ at low & high mc **Switch:** $\frac{4p'^2 p_\alpha^2}{(p'^2 + p_\alpha^2 + \gamma^2)^2} = \frac{(p' - p_\alpha)^2 + \gamma^2}{(p' + p_\alpha)^2 + \gamma^2}$

$$\psi_{p_\alpha l}^C(p') = -2\pi e^{-\eta_\alpha \pi/2} (p' p_\alpha)^l \left[\frac{\Gamma(l + 1 + i\eta_\alpha) \Gamma(\frac{1}{2})}{\Gamma(l + \frac{3}{2})} \right]$$

$$\times \lim_{\gamma \rightarrow +0} \left\{ \left[\left(\frac{2(p'^2 - (p_\alpha + i\gamma)^2)^{i\eta_\alpha}}{(p'^2 + p_\alpha^2 + \gamma^2)^{l+i\eta_\alpha+1}} \right) \left(\frac{\eta_\alpha(p_\alpha + i\gamma)}{p'^2 - (p_\alpha + i\gamma)^2} - \frac{\gamma(l + i\eta_\alpha + 1)}{p'^2 + p_\alpha^2 + \gamma^2} \right) \right. \right.$$

$$\left. \left. \times {}_2F_1 \left(\frac{l + i\eta_\alpha + 2}{2}, \frac{l + i\eta_\alpha + 1}{2}; l + \frac{3}{2}; \frac{4p'^2 p_\alpha^2}{(p'^2 + p_\alpha^2 + \gamma^2)^2} \right) \right] + \gamma \left[\dots \right] \right\}$$

Code will eventually be published

Integral is oscillatory singular:

Integral:

$$\int_0^\infty \frac{dp' p'^2}{2\pi^2} V_l(p_\alpha, p') \psi_{p_\alpha l}^C(p')$$

$V_l(p_\alpha, p')$ \rightarrow well-behaved function

$\psi_{p_\alpha l}^C(p')$ \rightarrow contains singularity!

Irregular (containing singularity)

$$\int_0^\infty \frac{dp' p'^2}{2\pi^2} V_l(p_\alpha, p') \psi_{p_\alpha l}^C(p') = \int_0^{p_\alpha - \Delta} \dots + \int_{p_\alpha - \Delta}^{p_\alpha + \Delta} \dots + \int_{p_\alpha + \Delta}^\infty \dots$$

Regular

Singularity regularized with Gel'fand-Shilov method

$$\int_{p_\alpha - \Delta}^{p_\alpha + \Delta} \dots \rightarrow \int_{p_\alpha - \Delta}^{p_\alpha + \Delta} \frac{\phi(p' - p_\alpha)}{(p' - p_\alpha + i\gamma)^{i\eta_\alpha + 1}} dp'$$

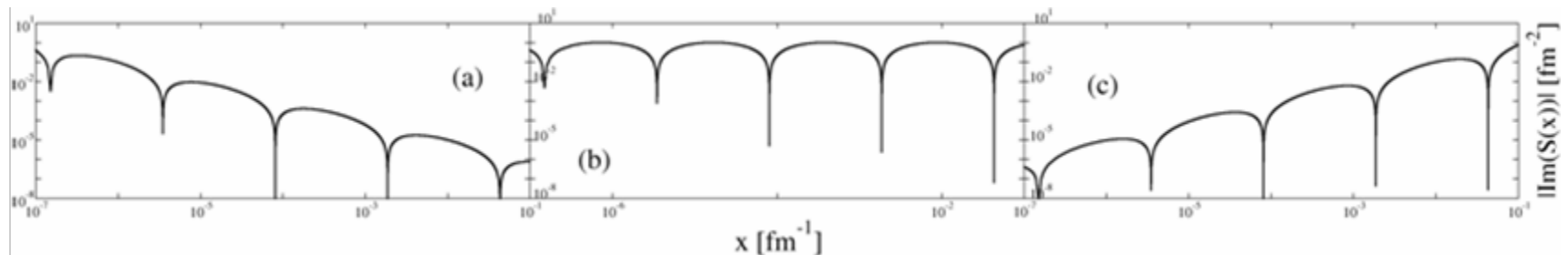
$$= \int_{p_\alpha - \Delta}^{p_\alpha + \Delta} dp' \frac{\left[\phi(p' - p_\alpha) - \phi(p_\alpha) - (p' - p_\alpha + i\gamma) \phi'(p_\alpha) \right]}{(p' - p_\alpha + i\gamma)^{i\eta_\alpha + 1}}$$

$$+ \frac{i\phi(p_\alpha)}{\eta_\alpha} \left[(\Delta + i\gamma)^{-i\eta_\alpha} - (-\Delta + i\gamma)^{-i\eta_\alpha} \right]$$

$$\Delta \sim 10^{-6} - 10^{-4}$$

$$+ \frac{\phi'(p_\alpha)}{(1 - i\eta_\alpha)} \left[(\Delta + i\gamma)^{1 - i\eta_\alpha} - (-\Delta + i\gamma)^{1 - i\eta_\alpha} \right]$$

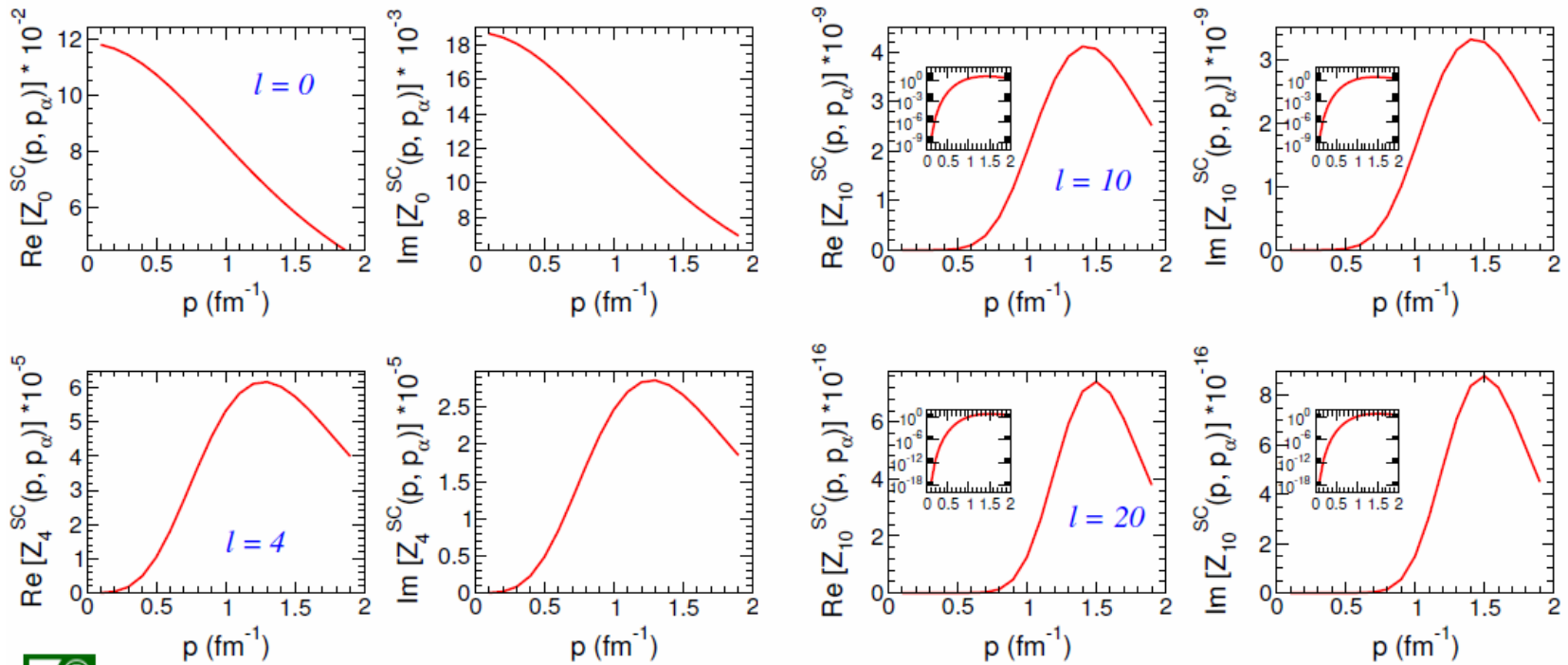
Then contribution of pole region very small



$V_l(p, p') \longrightarrow$ Yamaguchi form

$$Z_l^{SC}(p, p_\alpha) = \int_0^\infty \frac{dp' p'^2}{2\pi^2} V_l(p, p') \psi_{p_\alpha l}^C(p')$$

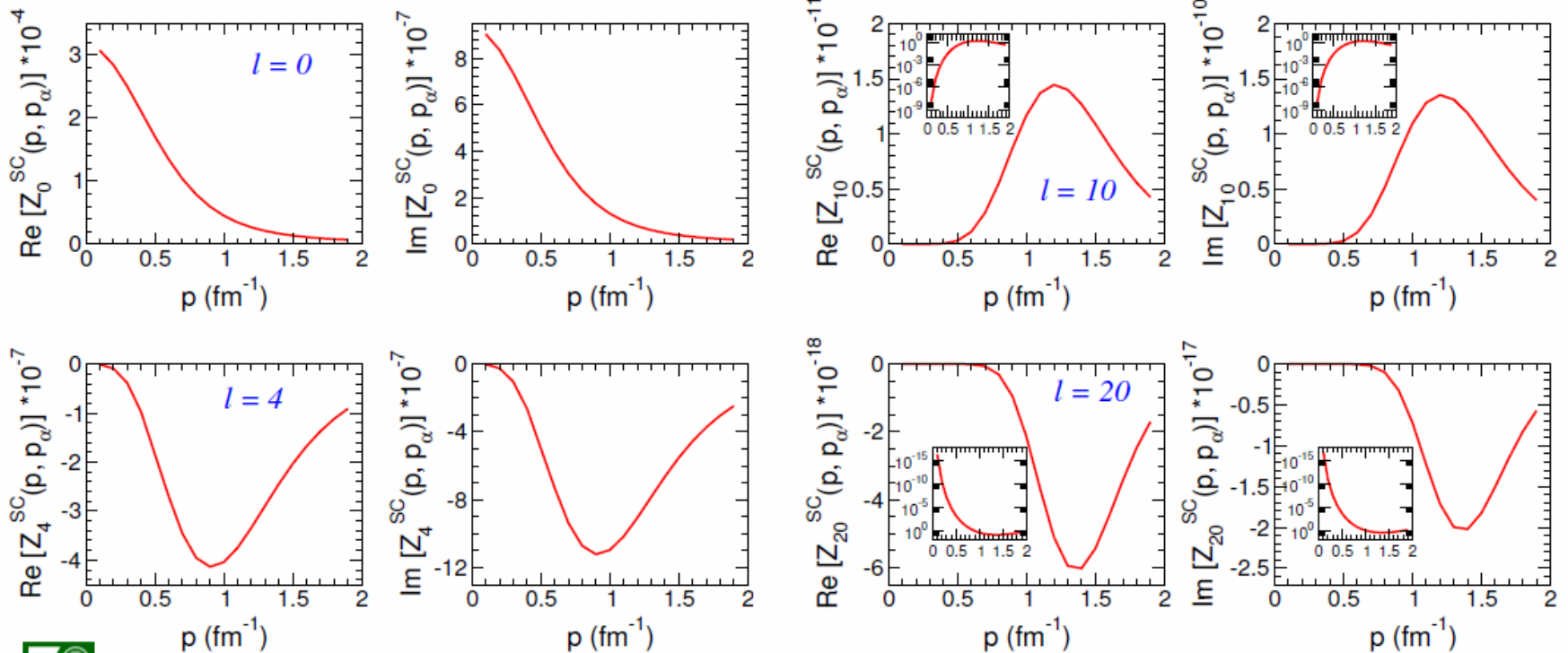
$p + {}^{12}\text{C}$ @ $E_{\text{cm}} = 10$ MeV:



$V_l(p, p') \longrightarrow$ Yamaguchi form

$$Z_l^{SC}(p, p_\alpha) = \int_0^\infty \frac{dp' p'^2}{2\pi^2} V_l(p, p') \psi_{p_\alpha l}^C(p')$$

$p + {}^{208}\text{Pb}$ @ $E_{\text{cm}} = 8$ MeV:



Looks promising



Faddeev Equations in Coulomb basis

Developing pw Coulomb wave functions in momentum space
essential to distinguish pole and non-pole regions
by use of different expansions of hypergeometric functions

Regularization of integrals over Coulomb wave function with
Gel'fand-Shilov method
tested so far with Yamaguchi type functions
for numerical and Mathematica cross checks

Next calculations with EST form factors



Persistence of Neelam Upadhyay and Vasily Eremenko

TORUS postdocs

TORUS: Theory of Reactions for Unstable Isotopes

A Topical Collaboration for Nuclear Theory



Ian Thompson, Jutta Escher (LLNL)

Filomena Nunes (NSCL MSU)

Akram Mukhamedzhanov (TAMU)

Ch. E. (OU)

**Few-Body
Collaboration**

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Post docs: Neelam Upadhyay (NSCL)

Vasily Eremenko (OU)

Grad Student: Linda Hlophe (OU)