



... for a brighter future



U.S. Department  
of Energy

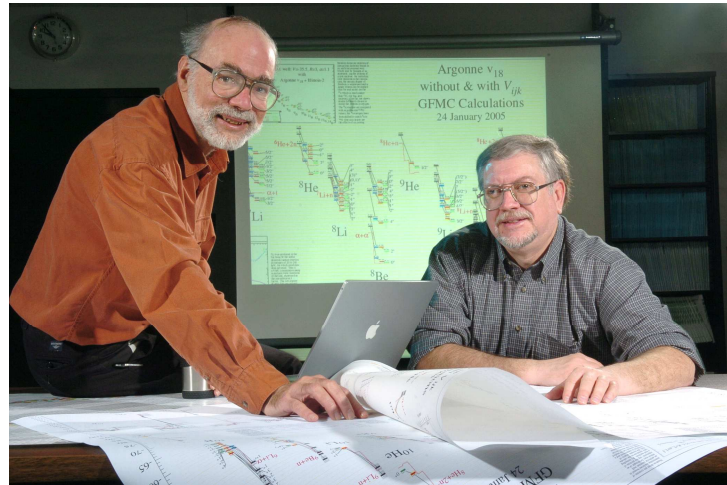
UChicago ►  
Argonne<sub>LLC</sub>



A U.S. Department of Energy laboratory  
managed by UChicago Argonne, LLC

# Tests of ab-initio shell models of light nuclei:

## The special case of $^{10}\text{Be}$



Steve Pieper

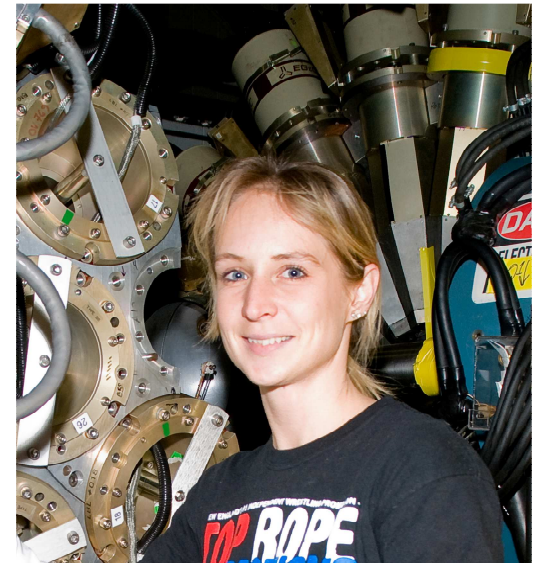
Bob Wiringa

C.J.(Kim) Lister

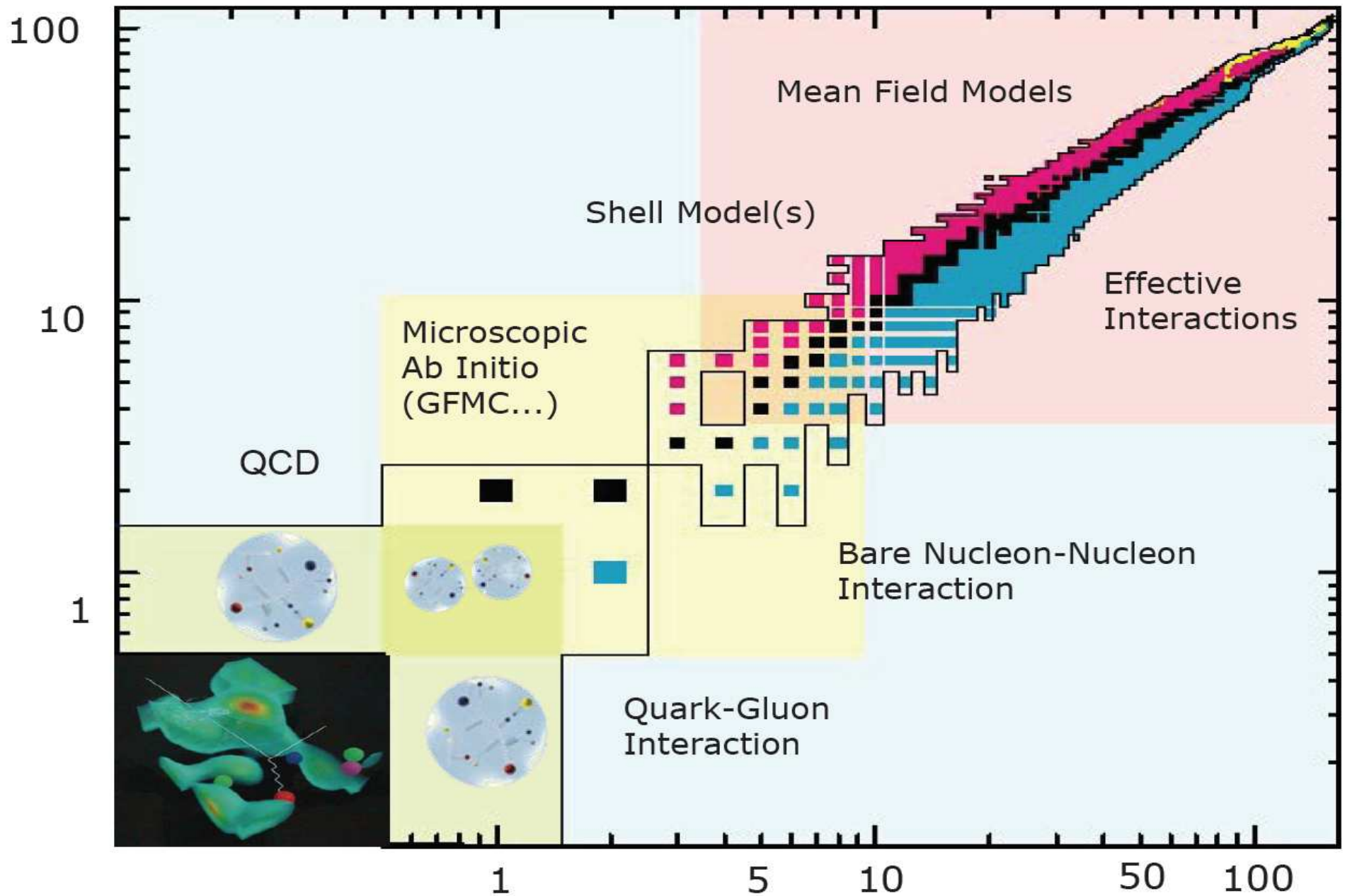
[Lister@anl.gov](mailto:Lister@anl.gov)

E.A. (Libby) McCutchan

[McCutchan@phy.anl.gov](mailto:McCutchan@phy.anl.gov)



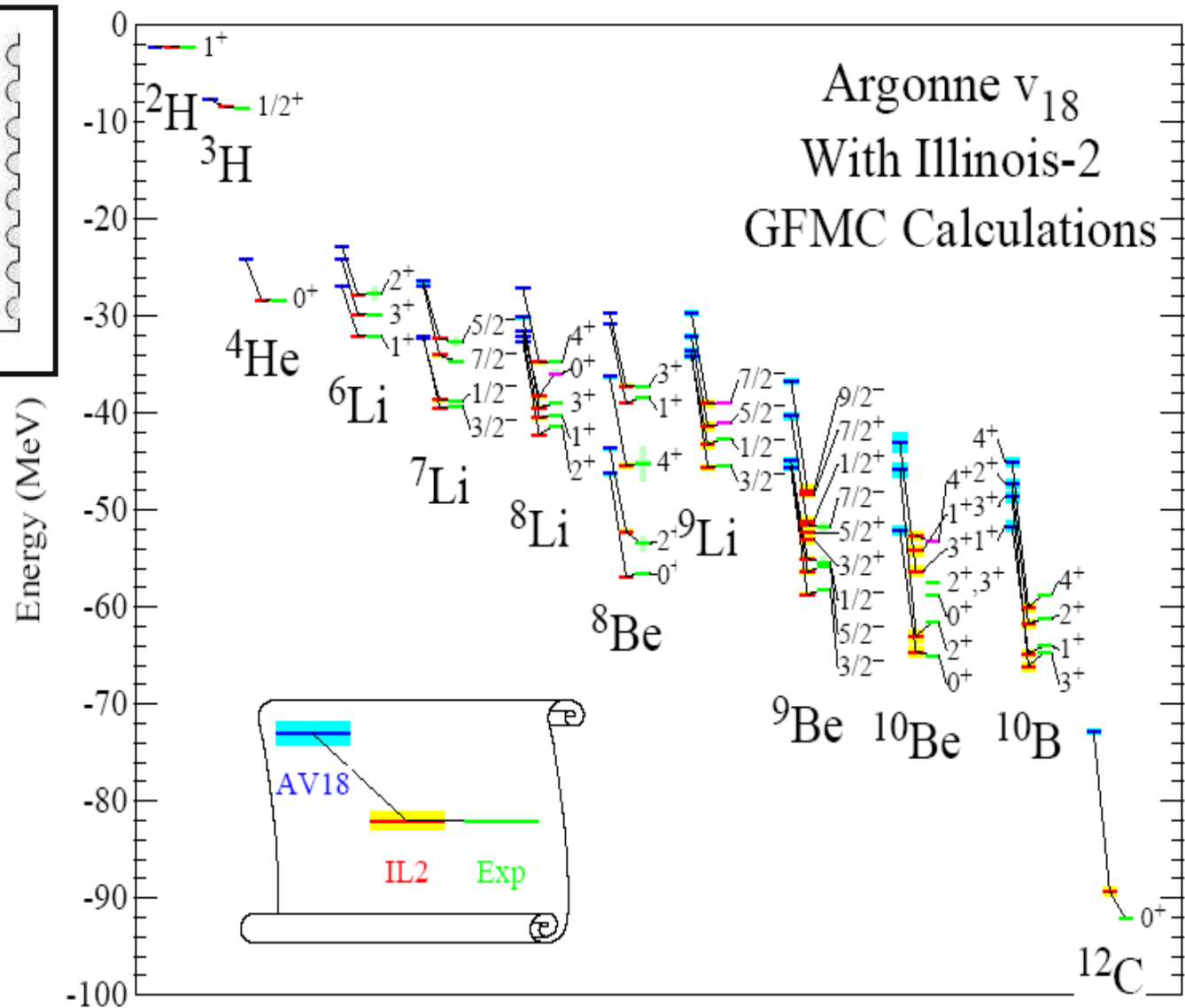
# Our nuclear world



# Ab-initio Greens Functional Monte Carlo Calculations

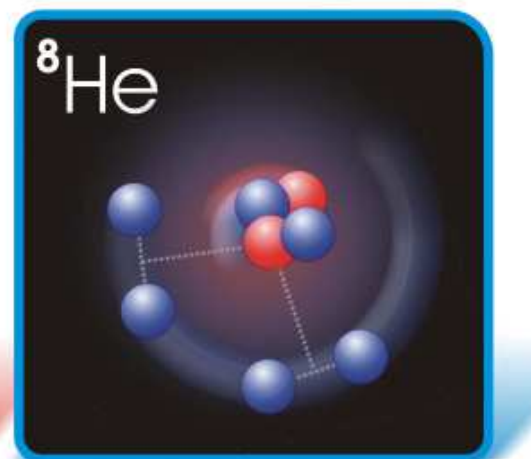
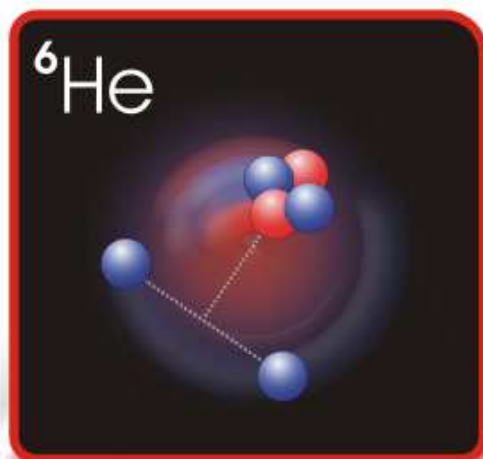
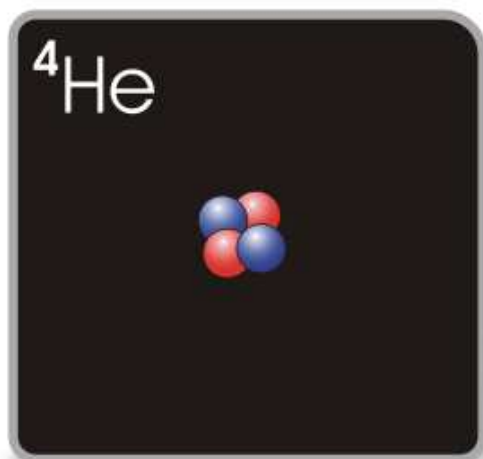
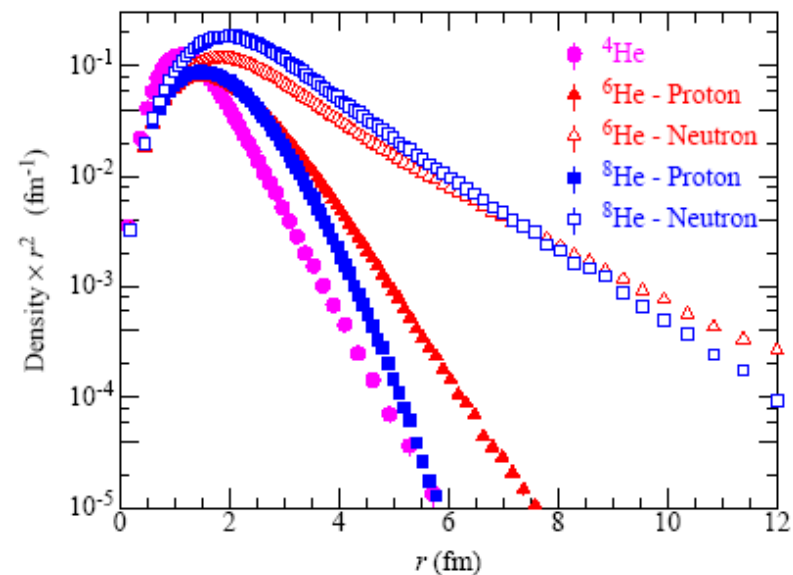
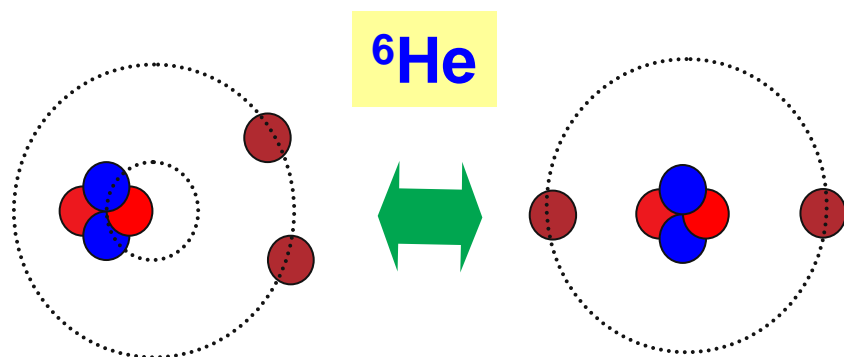


Pieper and  
Wiringa (ANL)  
Described in  
NP A751 516c  
(2005)



# RMS Radii of ${}^6, {}^8\text{He}$

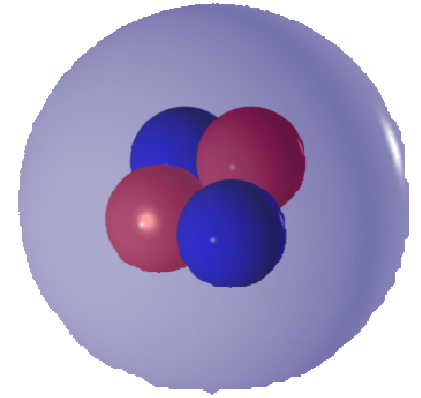
P. Mueller et. al (ANL)



# Transfer Reactions

ANL-UWM Collaboration

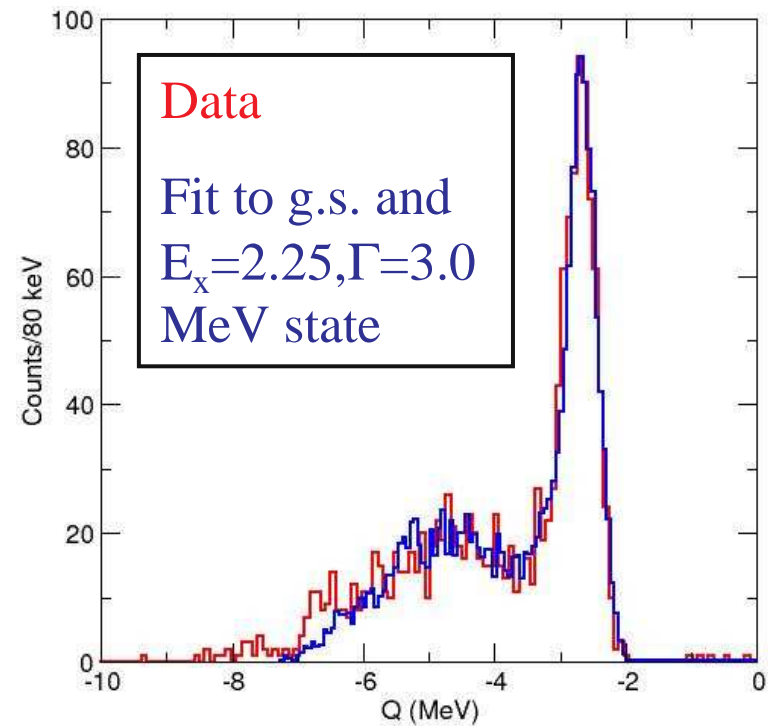
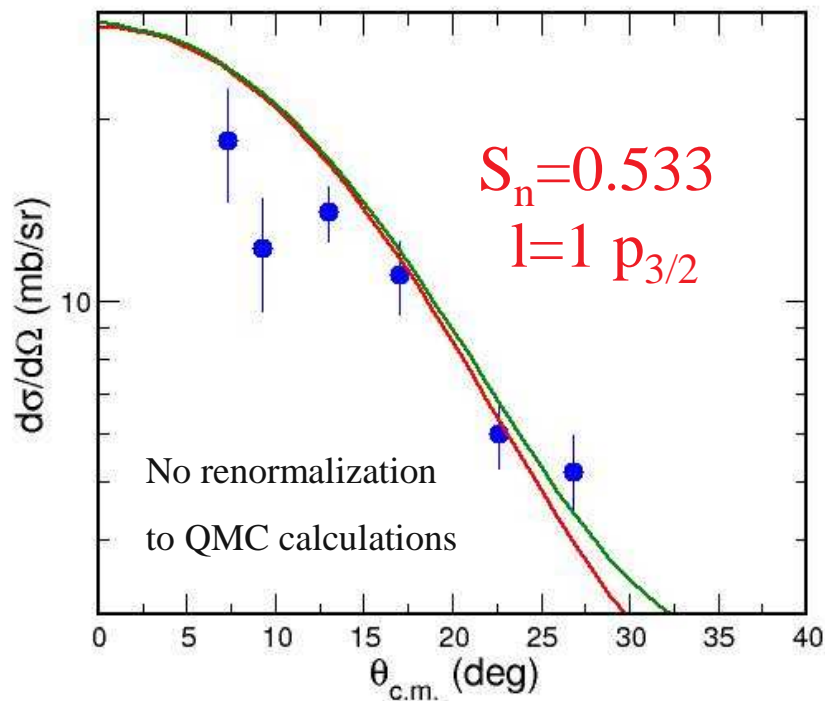
A. H. Wuosmaa et. al. PRL94 082592 and PRC72 061301(R) 2005



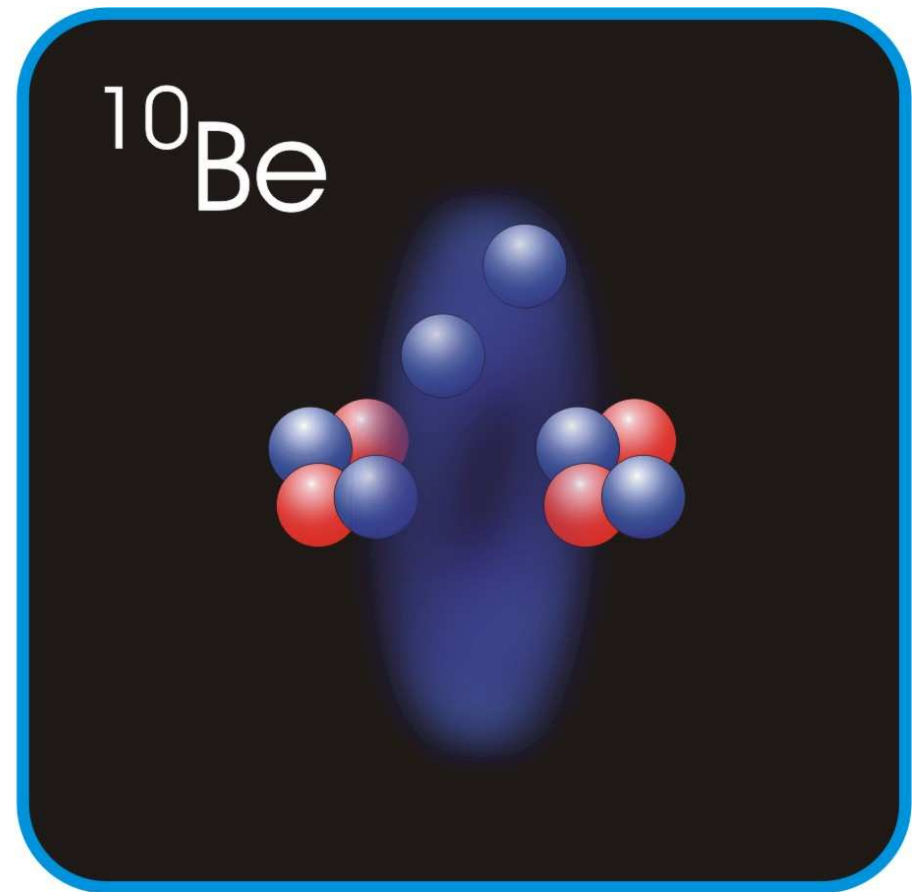
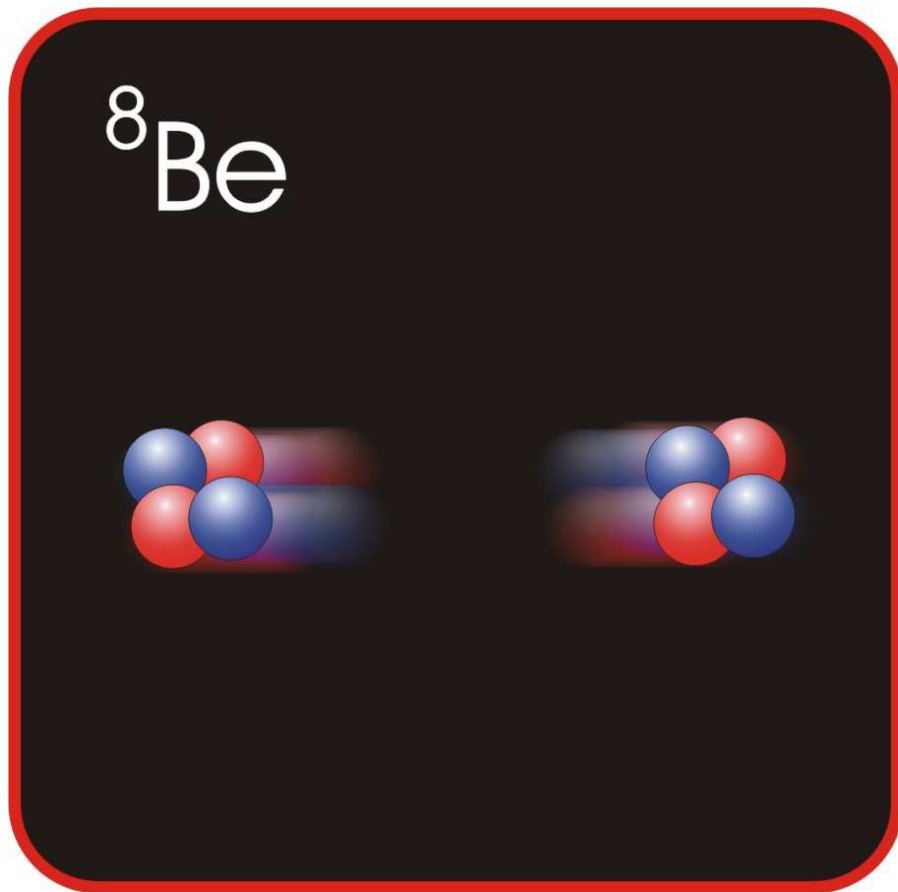
Example:  ${}^6\text{He}(d,p){}^7\text{He}$

Issues: How reliable are calculations of unbound states?

Is there a low lying state below 1MeV?

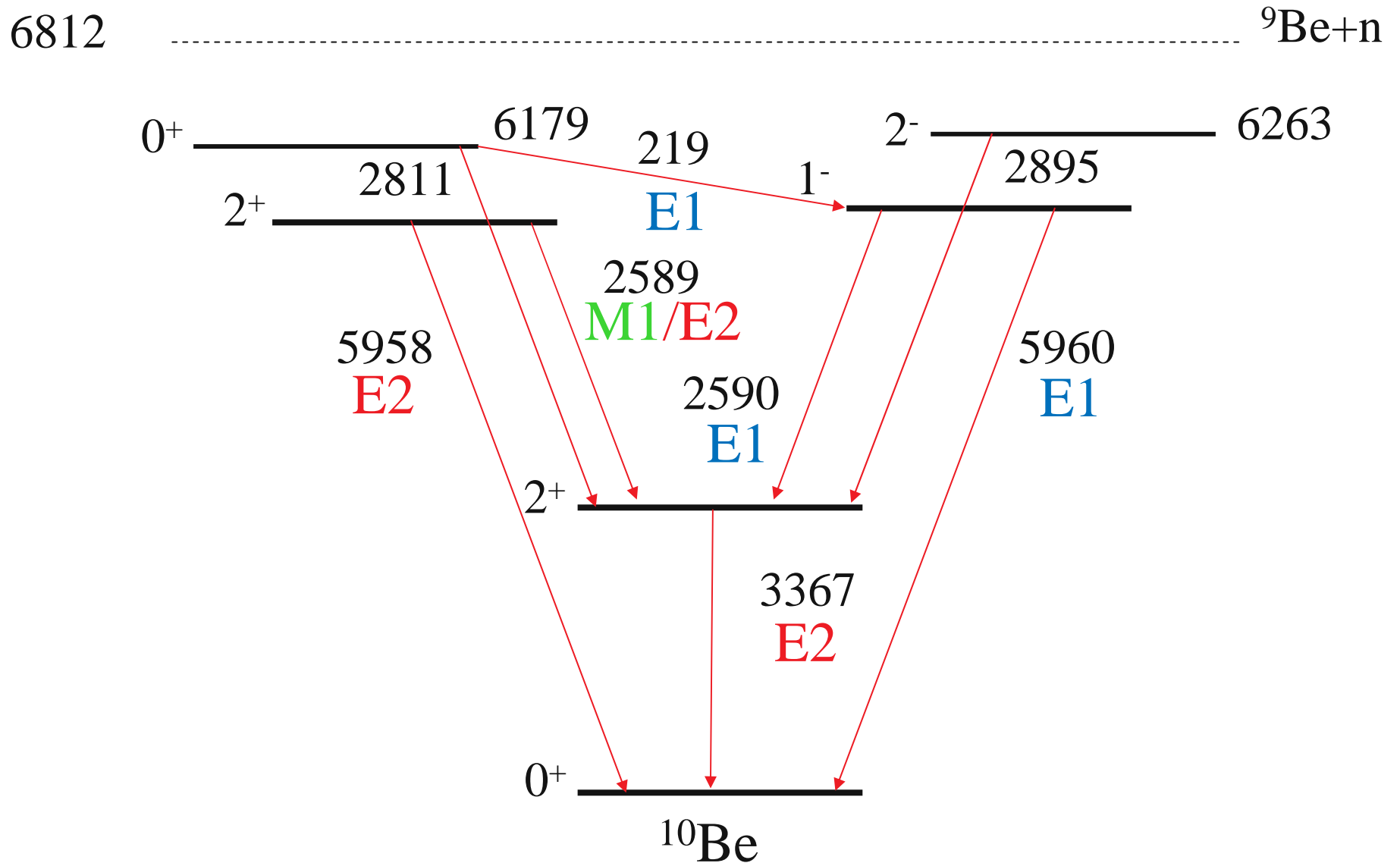


# The Special case of $^{10}\text{Be}$

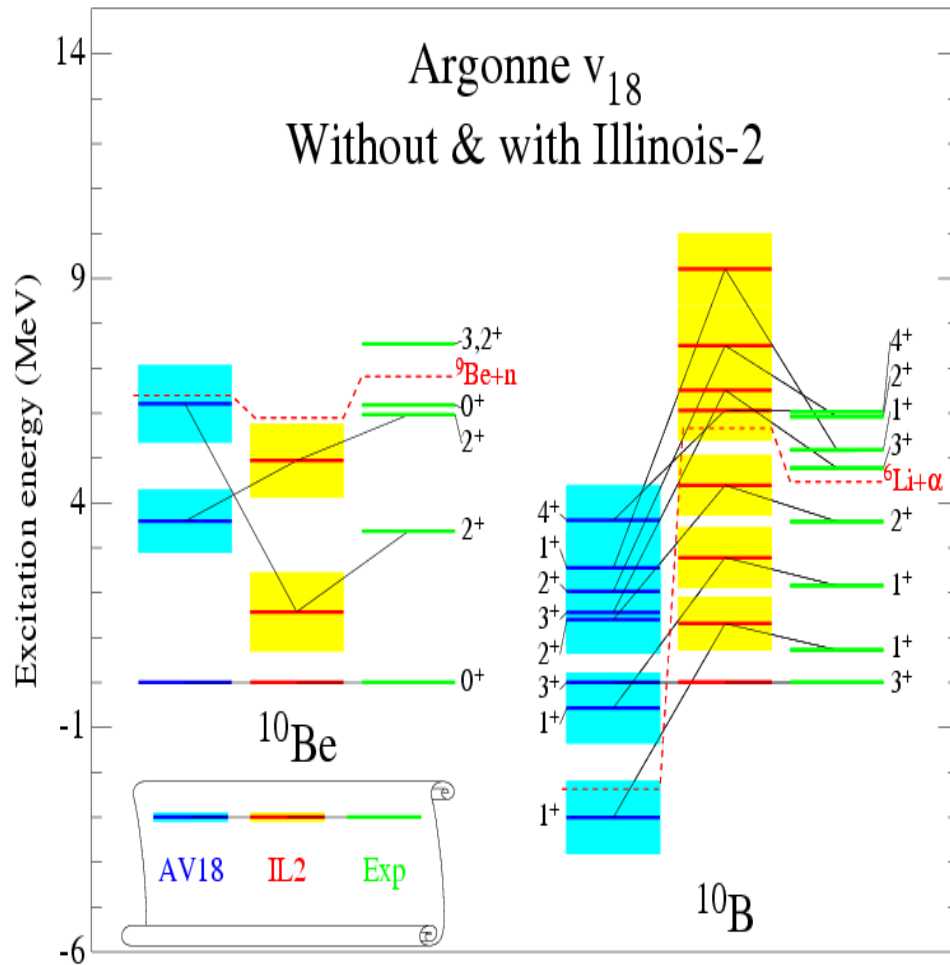


Key Issues: Microscopically, how do the neutrons really form a potential to bind the nucleus? Is it the same for all states?

# The Bound States of $^{10}\text{Be}$



# A=10 nuclei provide a sensitive test of GFMC



Very sensitive to the effects  
of 3-body correlations.

## $^{10}\text{Be}$ Data Compilation

$$B(E2; 2_1^+ \rightarrow 0_1^+) = 10.5(10)e^2 \text{fm}^4$$

## GFMC Theory

$$B(E2; 2_1^+ \rightarrow 0_1^+) = 7.89(30)e^2 \text{fm}^4 \quad Q \leq 0$$

$$B(E2; 2_2^+ \rightarrow 0_1^+) = 2.08(26)e^2 \text{fm}^4$$

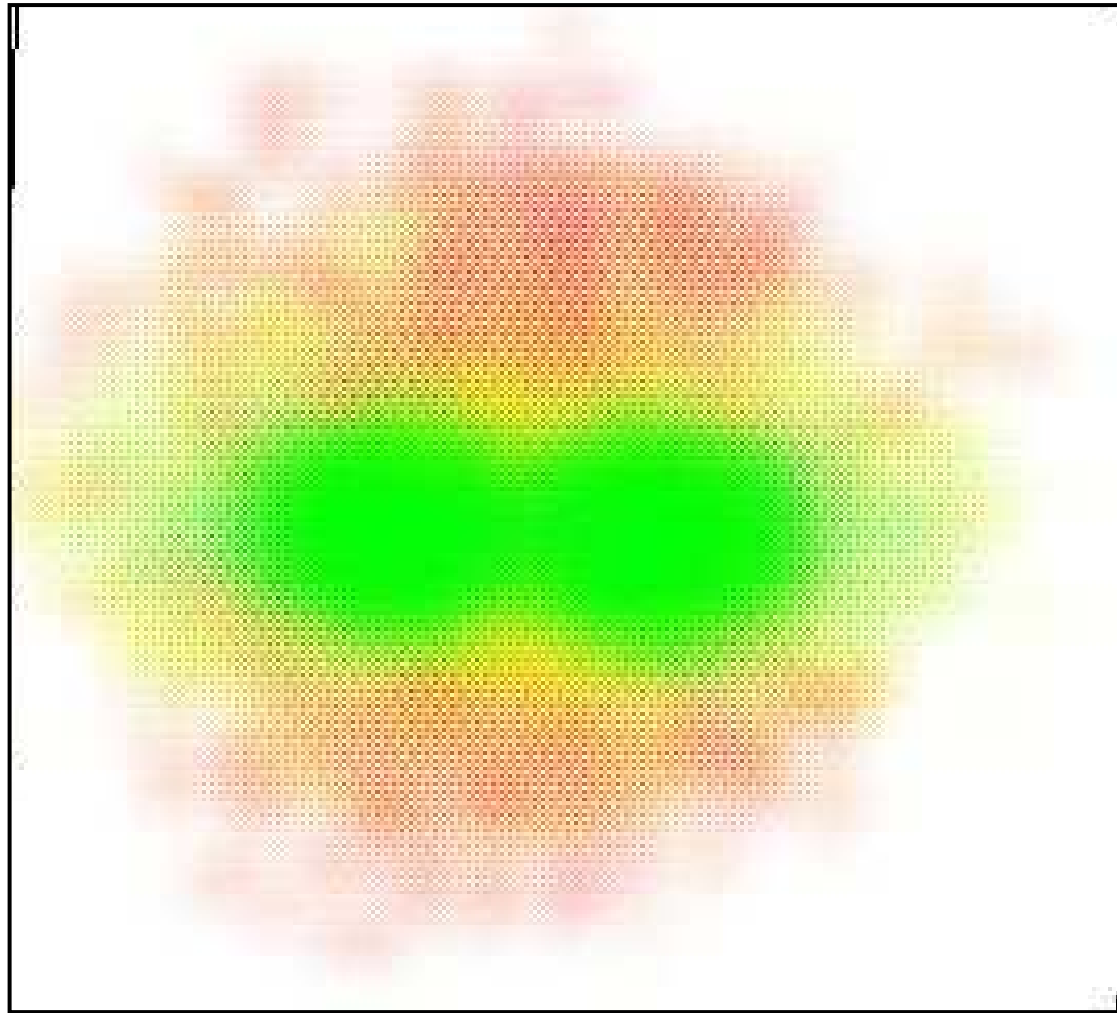
$$B(E2; 2_2^+ \rightarrow 2_1^+) = 6.9(10)e^2 \text{fm}^4 \quad Q \geq 0$$

$$R = \frac{B(E2: 2_1 \rightarrow 0_1)}{B(E2: 2_2 \rightarrow 0_1)} = 3.8$$

S.C. Pieper, K. Varga and R.B. Wiringa,  
Phys. Rev. C **66**, 044310 (2002).



# The Ab-Initio projection of T=0 and T=1 density



Ken Nollett (ANL) and Brent Graner (UC)  
visualization of  $^{10}\text{Be}$   $J=2_1$  state (2008)

# The Traditional Shell Model (1956+++)

PHYSICAL REVIEW

VOLUME 106, NUMBER 5

JUNE 1, 1957

## Radiative Transition Widths in the $1p$ Shell\*

DIETER KURATH

Argonne National Laboratory, Lemont, Illinois

(Received February 25, 1957)

The nuclear wave functions, obtained from a shell model with variable strength of spin-orbit coupling, are used to compute  $M1$  and  $E2$  transition widths. Comparison with experiment is made for  $\text{Be}^8$ ,  $\text{B}^{10}$ ,  $\text{B}^{11}$ , and  $\text{C}^{12}$ . The agreement is not nearly so good as was that obtained for energy level schemes. The pure  $M1$  transitions are in good agreement with experiment. The values computed for  $E2$  transition strengths are found to be generally low, though about the right order of magnitude. This suggests the need for adding some collective behavior to the model.

Kurath, Cohen and Kurath, Millener and Kurath, Millener and Warburton

185

DECAY OF 6.18-MeV  $J^\pi = 0^+$  LEVEL OF  $\text{Be}^{10}$

1249

TABLE I. The (8-16)POT results of Cohen and Kurath<sup>a</sup> for the decay of the lowest two  $(J^\pi, T) = (2^+, 1)$  states of  $(1p)^6$ .

Transition	Multi-polarity	Transition strength <sup>b</sup> (W.u.)	$\Gamma_\gamma$ (meV) for $\text{Be}^{10}$
$(2^+, 1)_2 \rightarrow (2^+, 1)_1$	$M1$	$[0.0044 + 0.3824T_z]^2$	54.61
$(2^+, 1)_2 \rightarrow (2^+, 1)_1$	$E2$	$[1.8518\epsilon(0) - 0.1174T_z\epsilon(1)]^2$	1.58
$(2^+, 1)_2 \rightarrow (0^+, 1)_1$	$E2$	$[-0.0948\epsilon(0) + 1.1566T_z\epsilon(1)]^2$	7.41
$(2^+, 1)_1 \rightarrow (0^+, 1)_1$	$E2$	$[1.5925\epsilon(0) + 0.0584T_z\epsilon(1)]^2$	4.81

R=5.8

<sup>a</sup> Calculated from the results of Ref. 30 by Dr. I. S. Towner.

<sup>b</sup> The phase convention is that of Rose and Brink [H. J. Rose and D. M. Brink, Rev. Mod. Phys. 39, 306 (1967)].

# Previously measured DSAM $J=2_1$ lifetime in $^{10}\text{Be}$

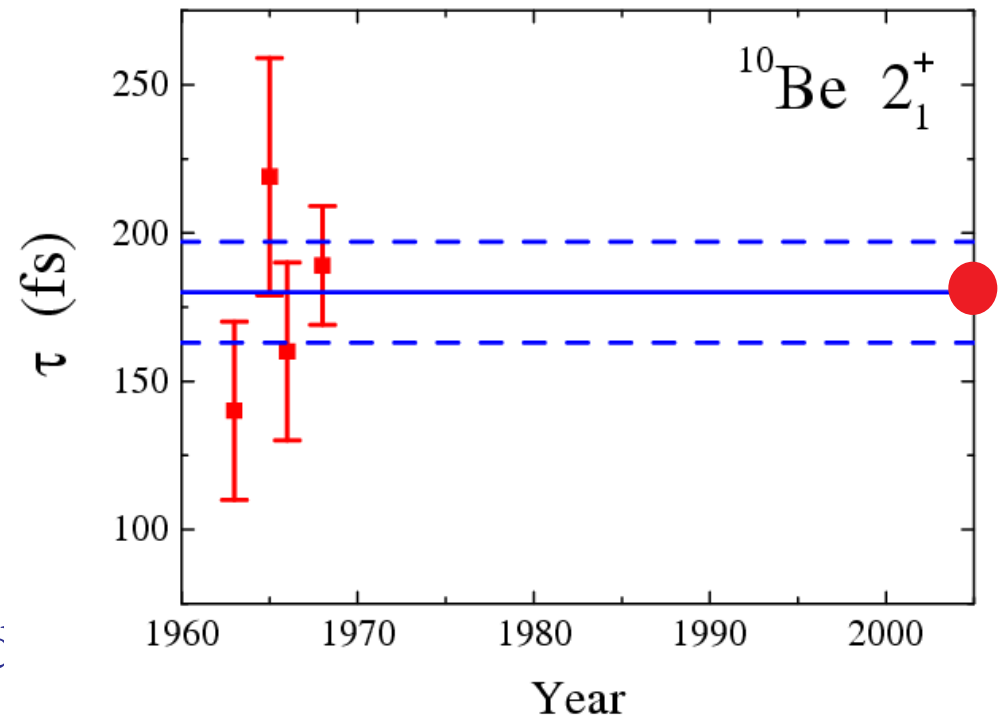
Systematic problems of:-

- Initial recoil velocity ill-defined
- Stopping powers poorly known
- Very slow recoils
- State feeding poorly known
- Detector angle poorly defined

At the time these measurements WERE cutting edge and provided discriminating tests of the (new) intermediate-coupling shell model (Kurath et al 1958) .

**But now we need better!**

Evaluated  $\tau = 180(18)$  fs  
All DSAM measurements



E.K. Warburton *et al.*, Phys. Rev. **129**, 2180 (1963).

G.C. Morrison *et al.*, - unpublished (1965).

E.K. Warburton *et al.*, Phys. Rev. **148**, 1072 (1966).

T.R. Fisher *et al.*, Phys. Rev. **176**, 1130 (1968).

# $^{10}\text{Be}$ : A popular nucleus

Author	Calc.	$B(E2; 2_1^+ \rightarrow 0_1^+)$ ( $\text{e}^2\text{fm}^4$ )	$B(E2; 2_2^+ \rightarrow 0_1^+)$ ( $\text{e}^2\text{fm}^4$ )	Ratio
Okabe	MO	11.26	0.44	25.6
Caurier	NCSM	6.58	0.13	49.5
Itagaki	TMO	11.8	0.70	16.9
Navratil	NCSM	5.4	0.17	32.7
Arai	MCM	4.8	0.10	48.0
Wiringa	GFMC	7.89(30)	2.08(26)	3.8
<b>EXPT</b>		<b>10.5(1.0)</b>	<b>??</b>	<b>??</b>

MO : Molecular orbit model, N. Itagaki and S. Okabe, Phys. Rev. C **61**, 044306 (2000).

NCSM : No-core shell model, E. Caurier *et al.*, Phys. Rev. C **66**, 024314 (2002).

TMO : Triaxial Molecular Orbital, N. Itagaki *et al.*, Phys. Rev. C **65**, 044302 (2002).

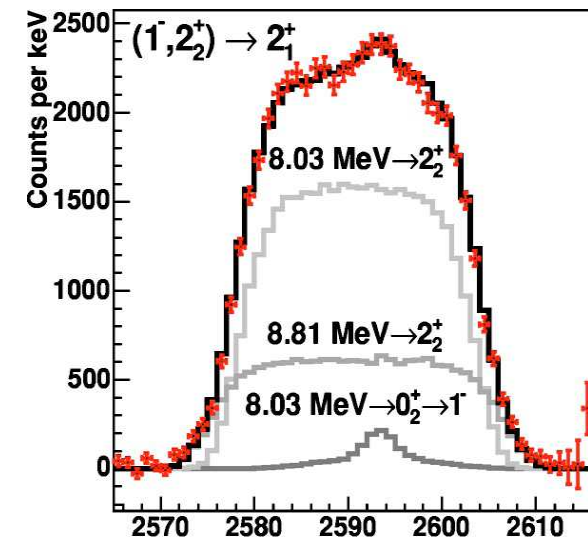
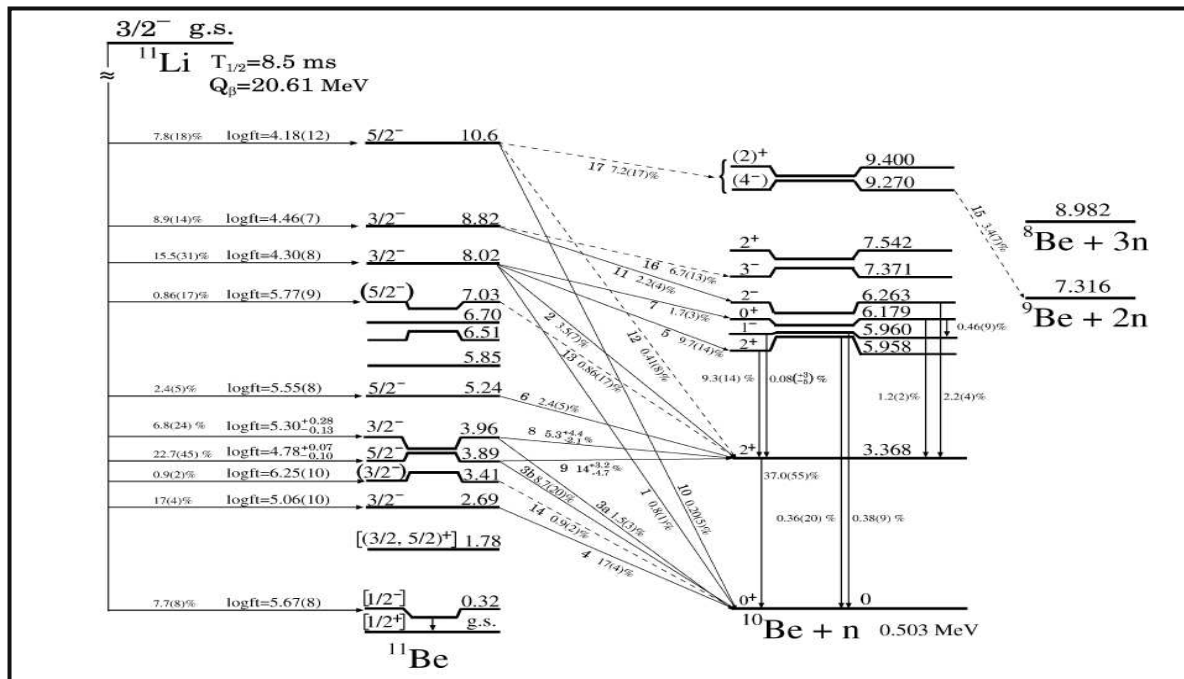
MCM : Microscopic cluster model, K. Arai, Phys. Rev. C **69**, 014309 (2004).

GFMC : Greens Functional Monte Carlo, R. Wiringa and S. Pierper *et al.*, Private Communication.

# $^{11}\text{Li}$ $\beta$ -delayed 1-n emission

Line shape depends on

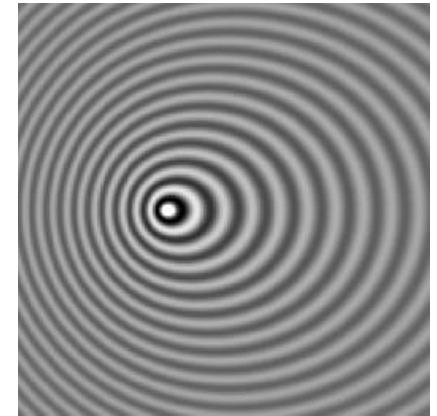
- Energies and intensities of all neutron branches feeding the level
- Lifetime of the level
- Angular correlation between gamma ray and neutron



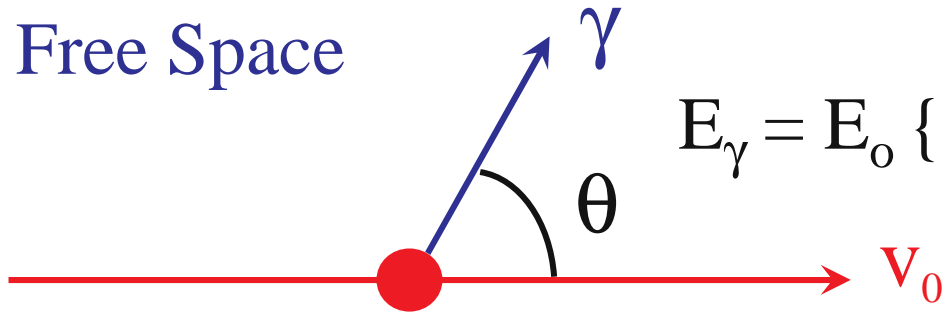
Y. Hirayama *et al.* Phys Lett **B611** 239 (2005).....polarized  $\beta$ -decay at ISAC/TRIUMF (high stats)

F. Sarazin *et al.*, Phys. Rev. C **70**, 031302(R) (2004).... $\beta$ -decay at ISAC / TRIUMF

# The Doppler Shift Attenuation Method

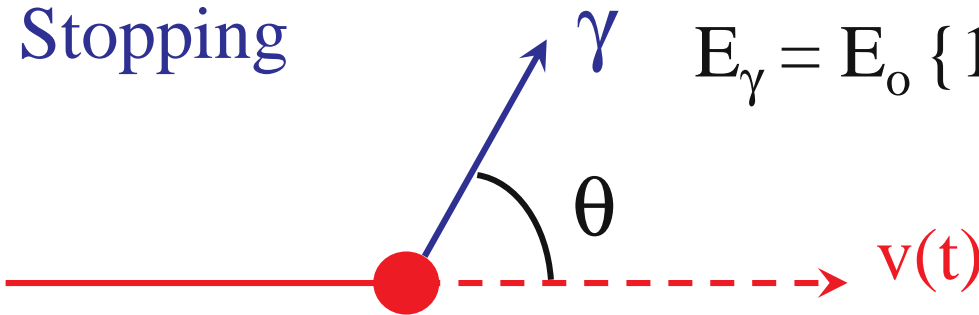


Free Space

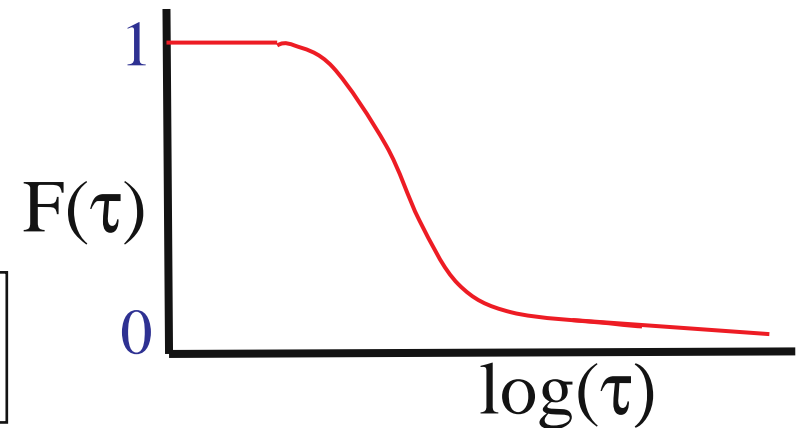


$$E_{\gamma} = E_0 \{ 1 + (v_0/c) \cos \theta \}$$

Stopping

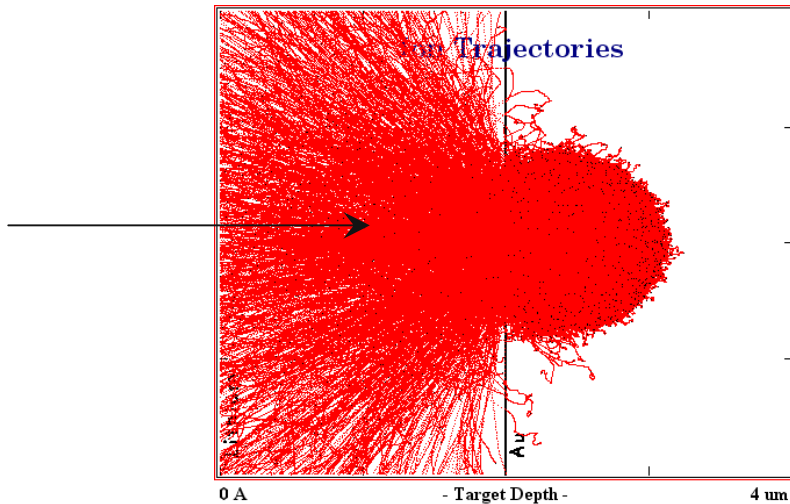
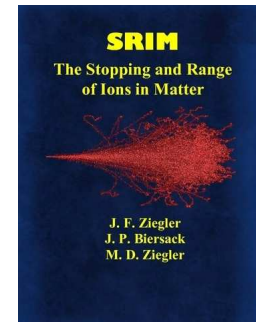


$$E_{\gamma} = E_0 \{ 1 + F(t) \cos \theta \}$$

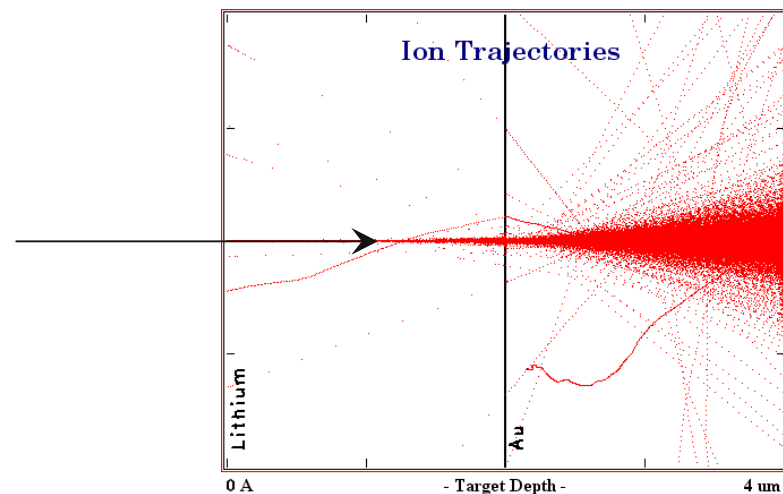


Actually, recoils are sufficiently swift ( $v/c \sim 6\%$ ) that 2<sup>nd</sup> order relativistic is correction needed

# What SRIM shows us



**SLOW** Ions.... 1 MeV <sup>10</sup>Be in  
100μg/cm<sup>2</sup> Li on 6mg/cm<sup>2</sup> Au



**FAST** Ions.... 15 MeV <sup>10</sup>Be in  
100μg/cm<sup>2</sup> Li on 6mg/cm<sup>2</sup> Au  
(99.99% electronic stopping)

J.F. Ziegler, J.P. Biersack and M.D. Ziegler <http://www.srim.org>

# DSAM\*

E.A. McCutchan et. al. (ANL)

DSAM done properly:-

2-body kinematics

High recoil velocity

Well defined recoil direction

Well defined gamma direction

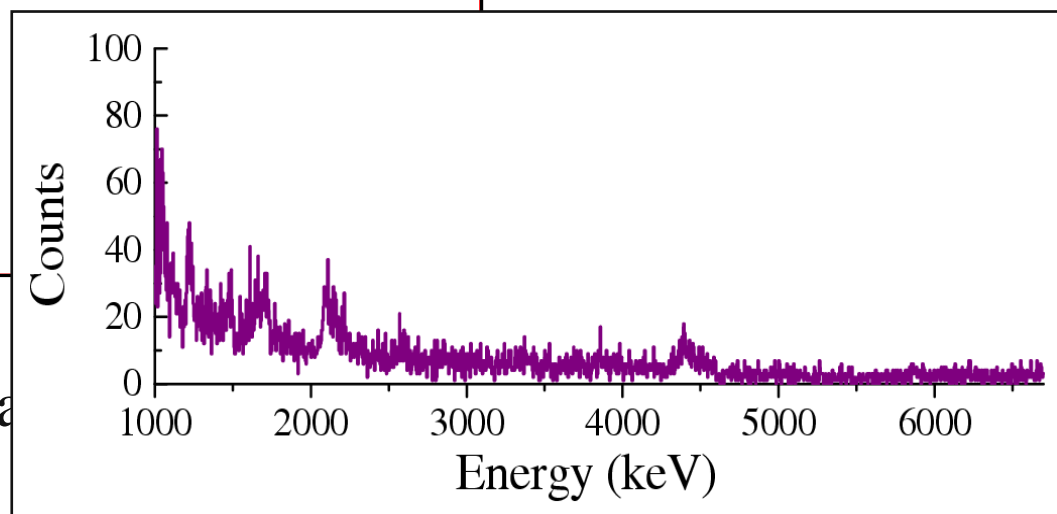
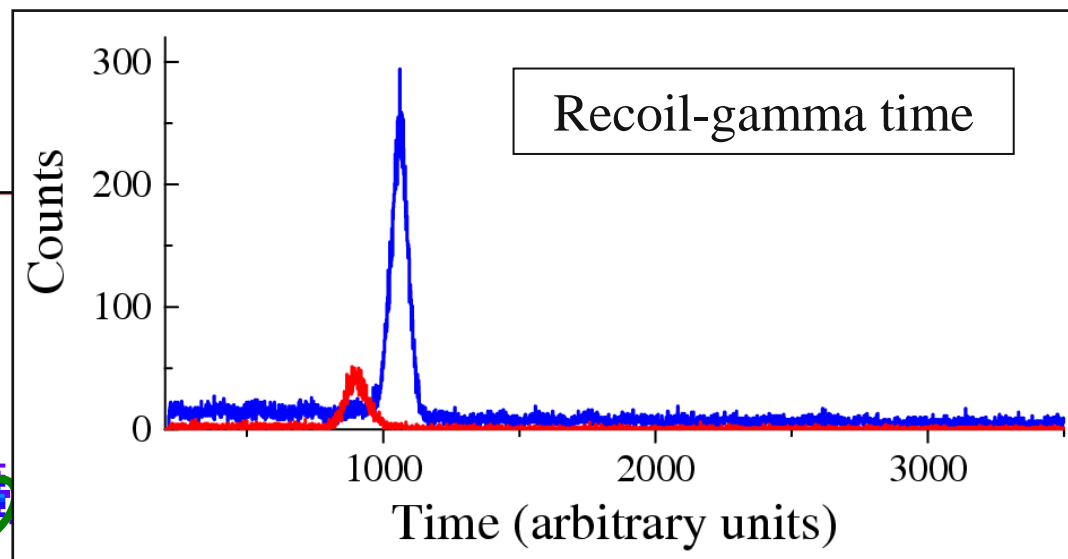
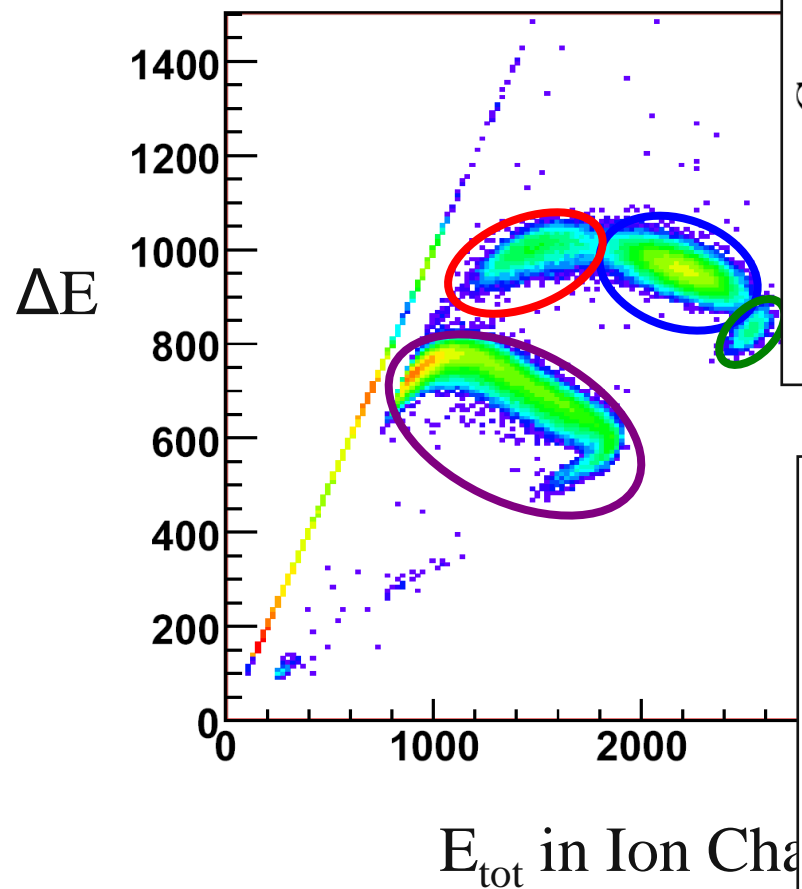
Control of state population





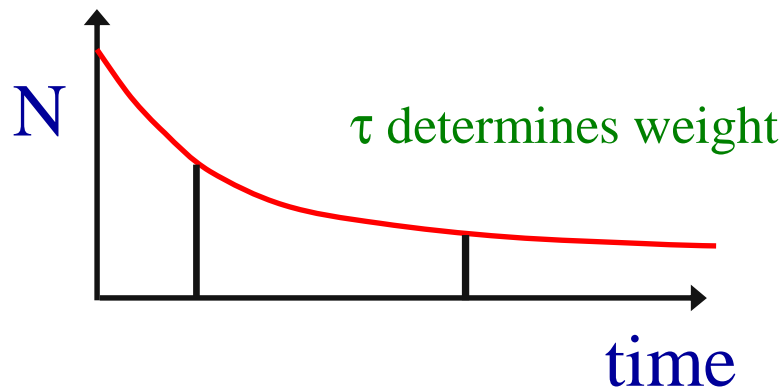
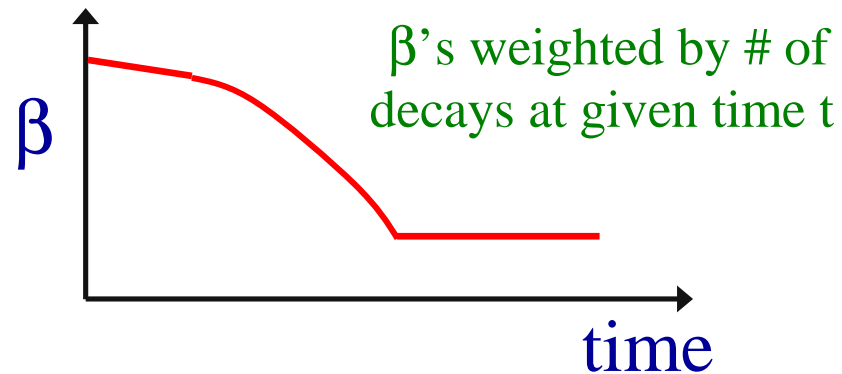
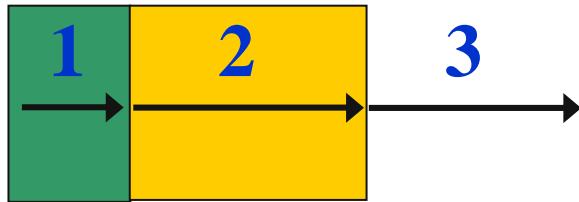
# Recoil - $\gamma$ Gated Data

2 mg/cm<sup>2</sup> Cu backing

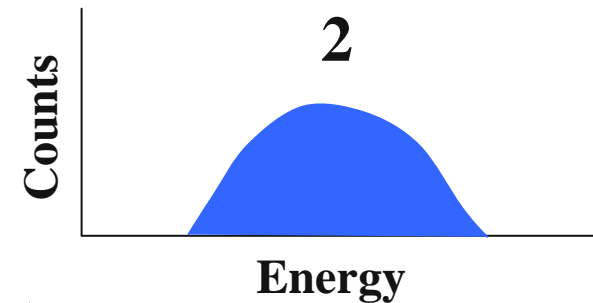


A=10 q=3 selected with FMA

# Extracting the lifetime from mean $\beta$

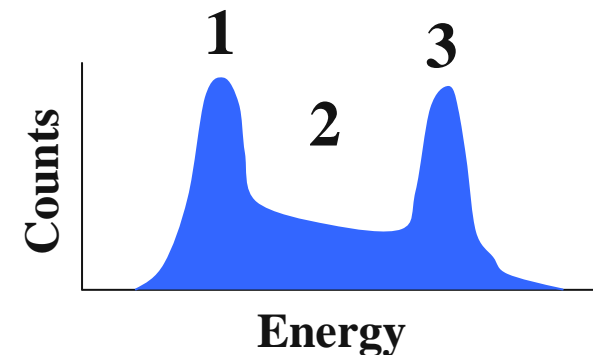


Ideally you want all decays in the stopping material

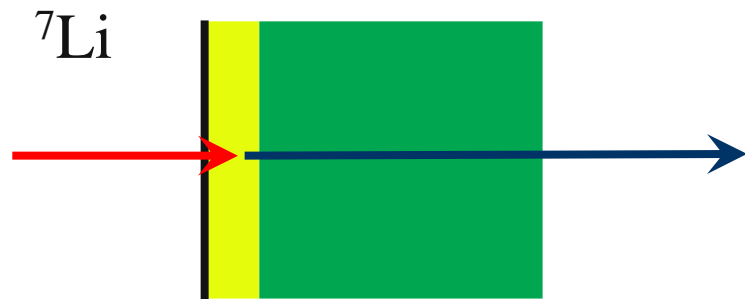


But...

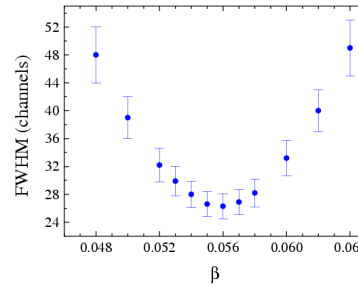
- Need some finite thickness of production material
- Need nuclei to recoil out of target



# Results (1): First excited State with $J^\pi=2^+$ at 3367 keV

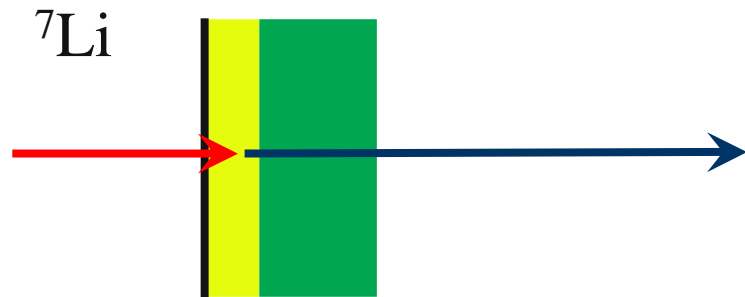


$66\mu\text{g}/\text{cm}^2$   ${}^7\text{Li}$  on  $6.95\text{mg}/\text{cm}^2$  Au

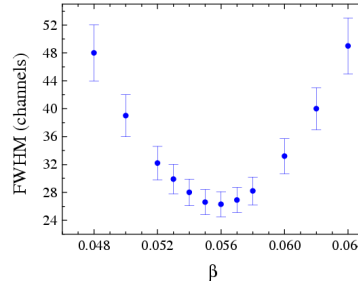


$$\beta = 0.0525(5)$$

$$\tau = 245(35)\text{fs}$$

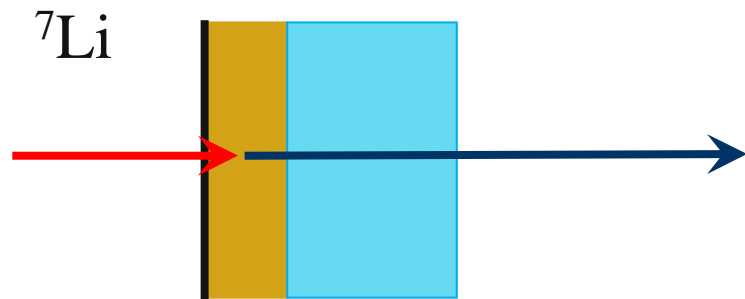


$115\mu\text{g}/\text{cm}^2$   ${}^7\text{Li}$  on  $2.69\text{mg}/\text{cm}^2$  Au

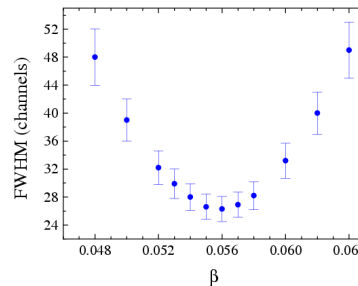


$$\beta = 0.0560(5)$$

$$\tau = 290(+200)_{-100}\text{fs}$$



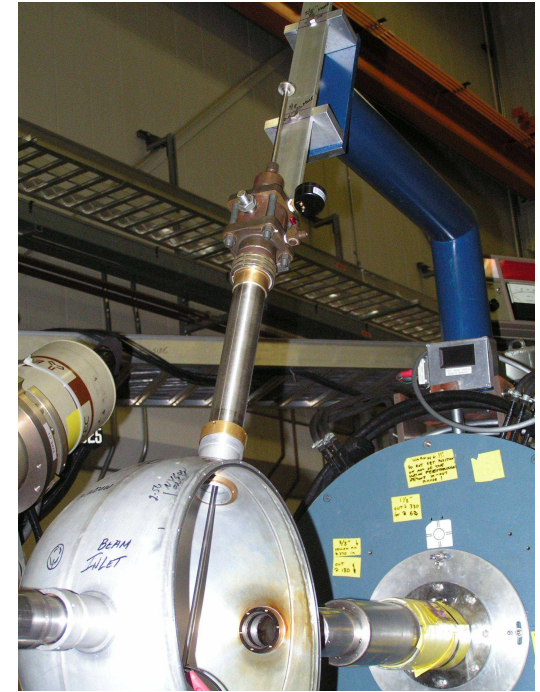
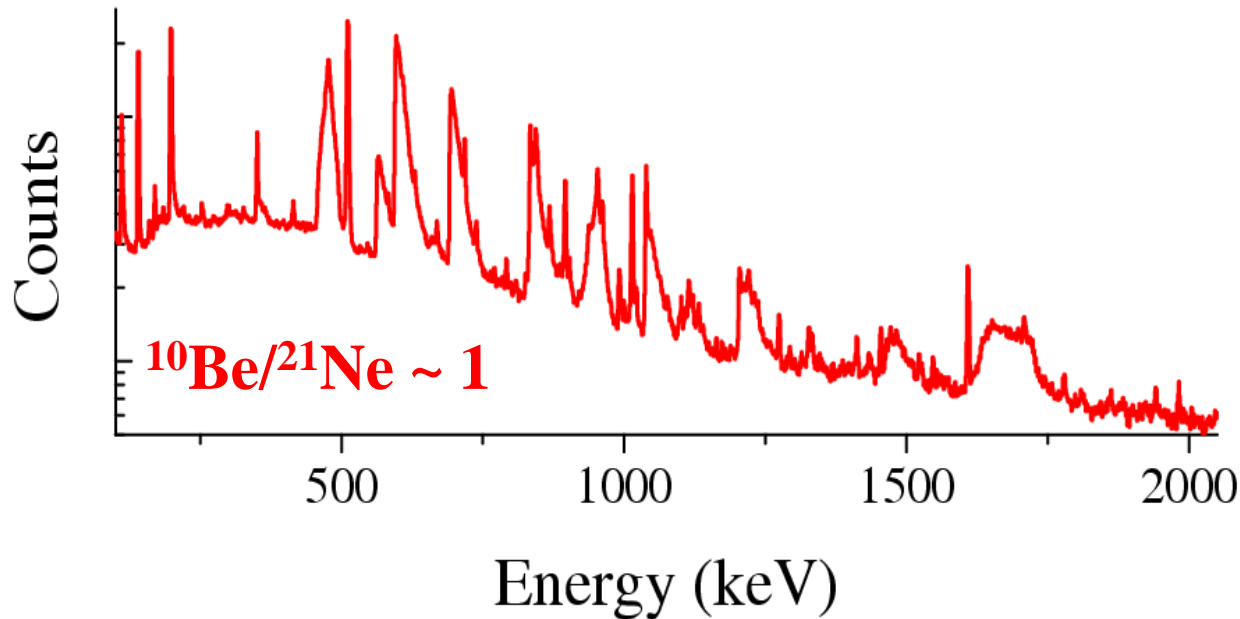
$288\mu\text{g}/\text{cm}^2$   ${}^7\text{LiF}_2$  on  $2.48\text{mg}/\text{cm}^2$  Cu



$$\beta = 0.0535(5)$$

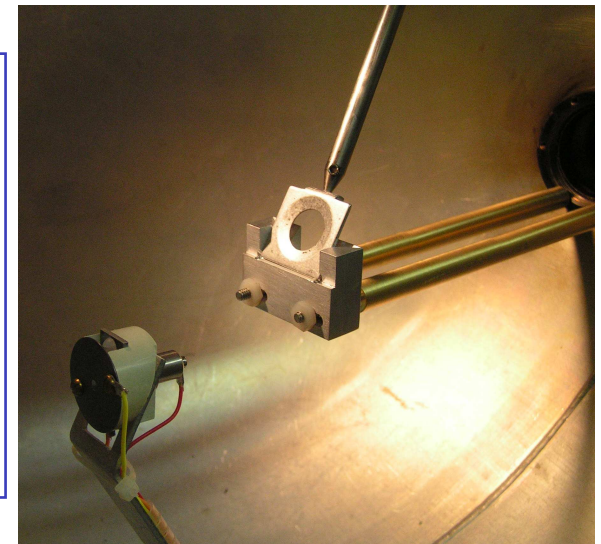
$$\tau = 270(+100)_{-60}\text{fs}$$

# Target Oxidation Issues



Lithium is volatile and oxidizes easily. A new vacuum interlock was developed to move from the evaporator to Gammasphere at  $\sim 10^{-5}$  torr.

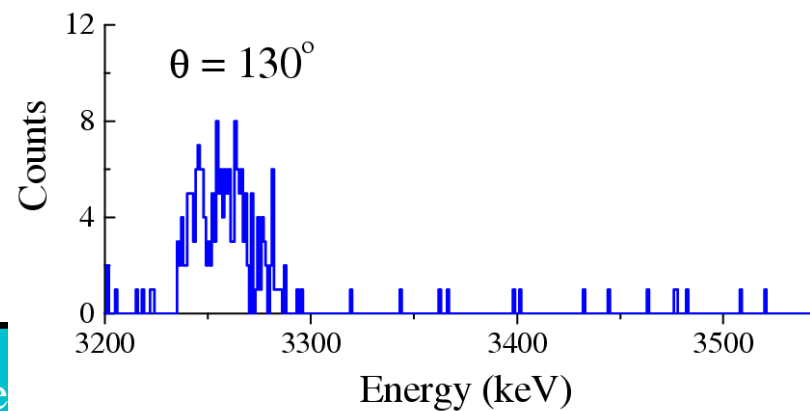
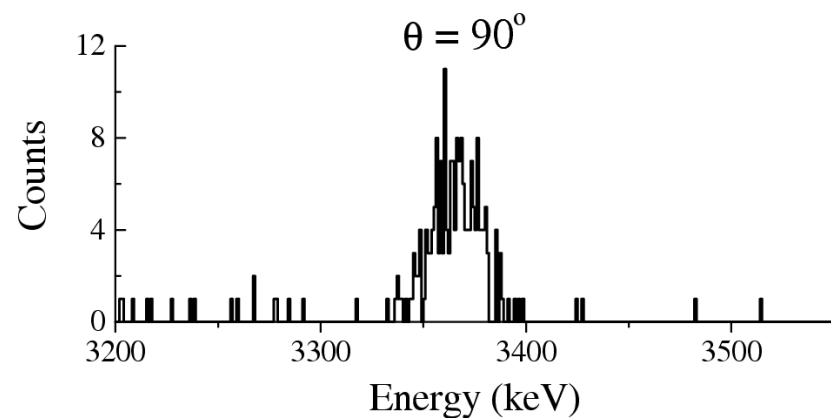
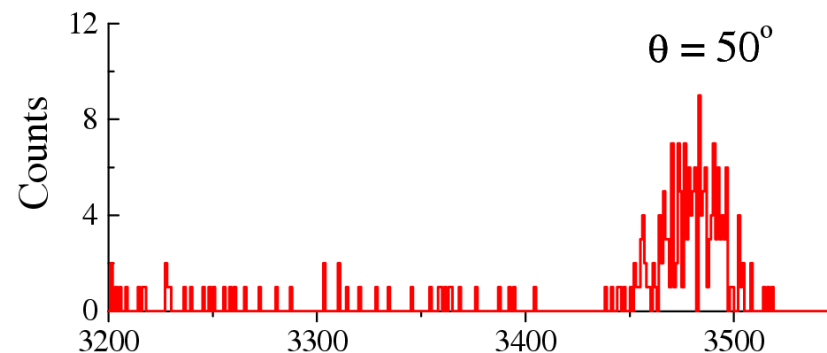
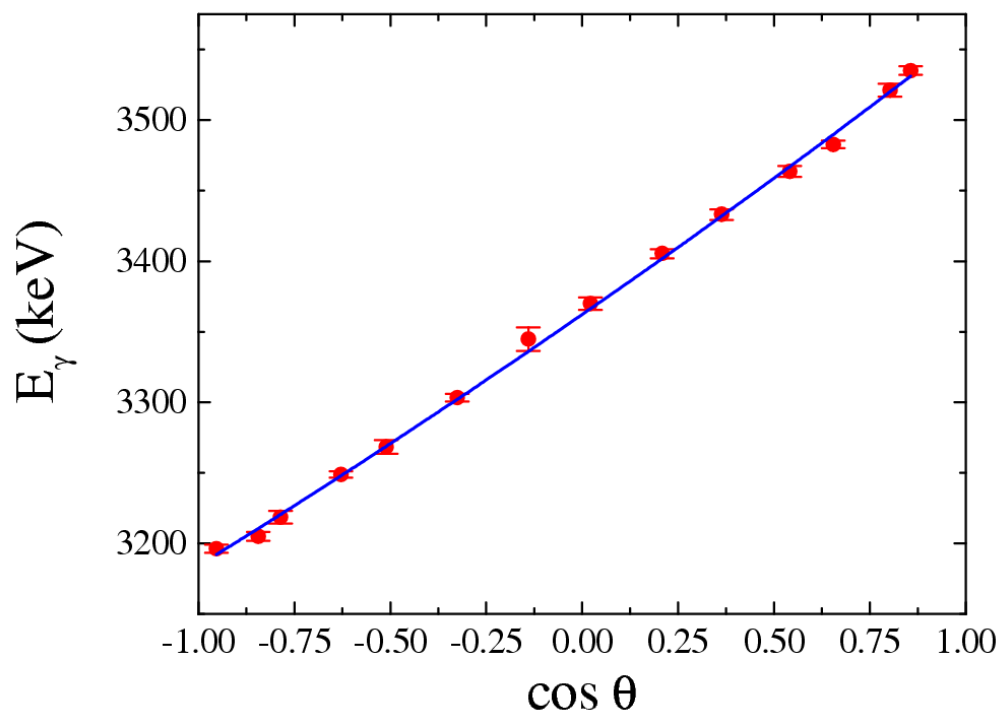
This reduced the contamination by a factor 10. With better vacuum, we expect another factor 10



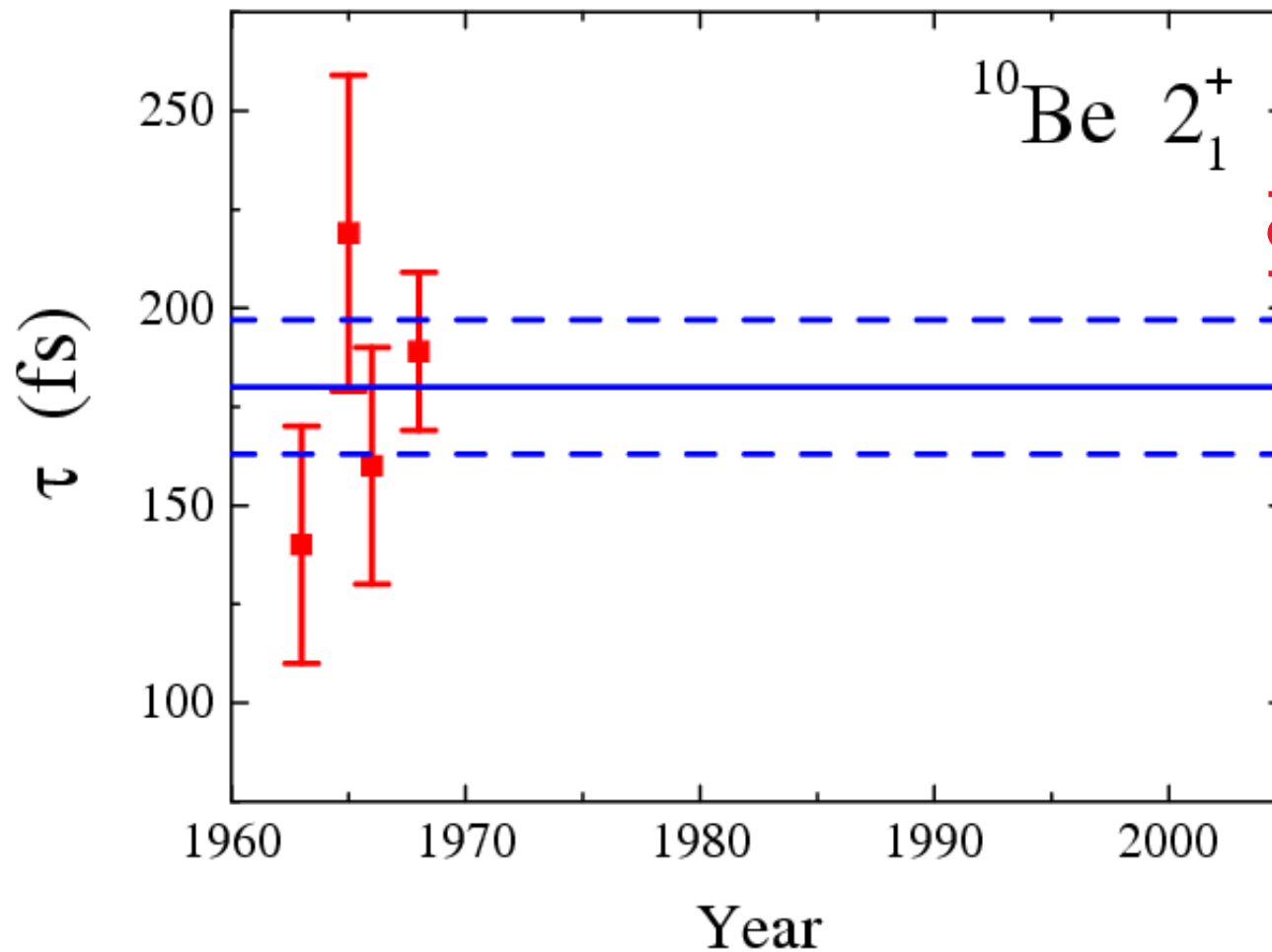
# Traditional DSAM analysis.....and GS alignment

$$E_{\gamma} = E_o \left( \frac{\sqrt{1 - \beta^2}}{1 - \beta \cos \theta} \right)$$

100  $\mu\text{g}/\text{cm}^2$   $^7\text{Li}$  on 2.33  $\text{mg}/\text{cm}^2$  Au



## Results so far for 3369keV J=2 state



# First $J = 2^+$ state in $^{10}\text{Be}$

NNDC transition strength

$$B(E2; 2_1^+ \rightarrow 0_1^+) = 10.5(10) e^2 \text{fm}^4$$

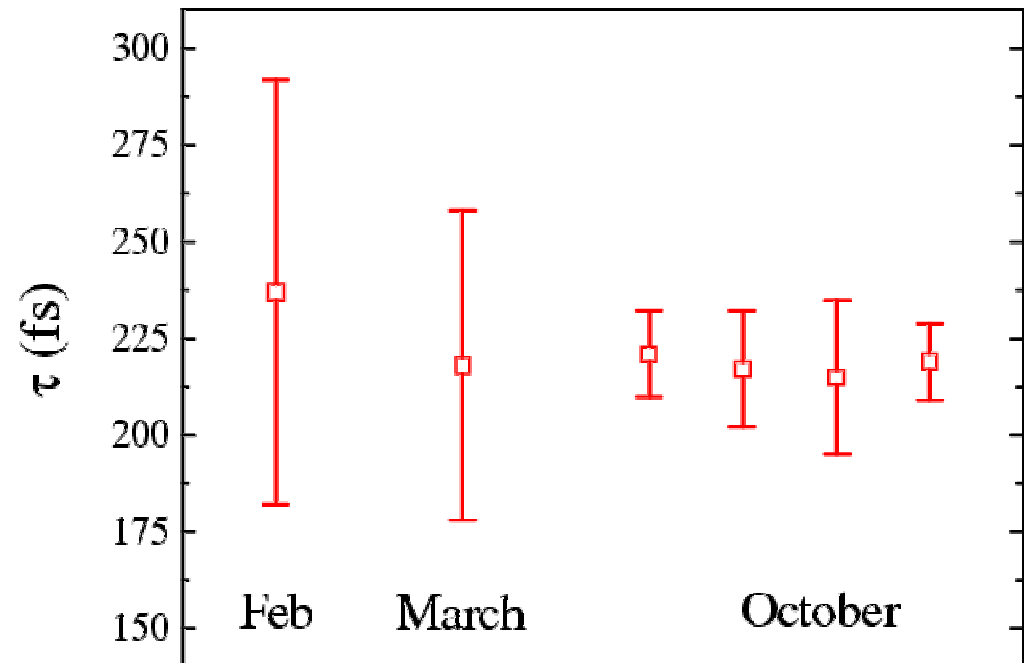
Revised transition strength

$$B(E2; 2_1^+ \rightarrow 0_1^+) = 8.7(3) e^2 \text{fm}^4$$

Ab-initio transition strength

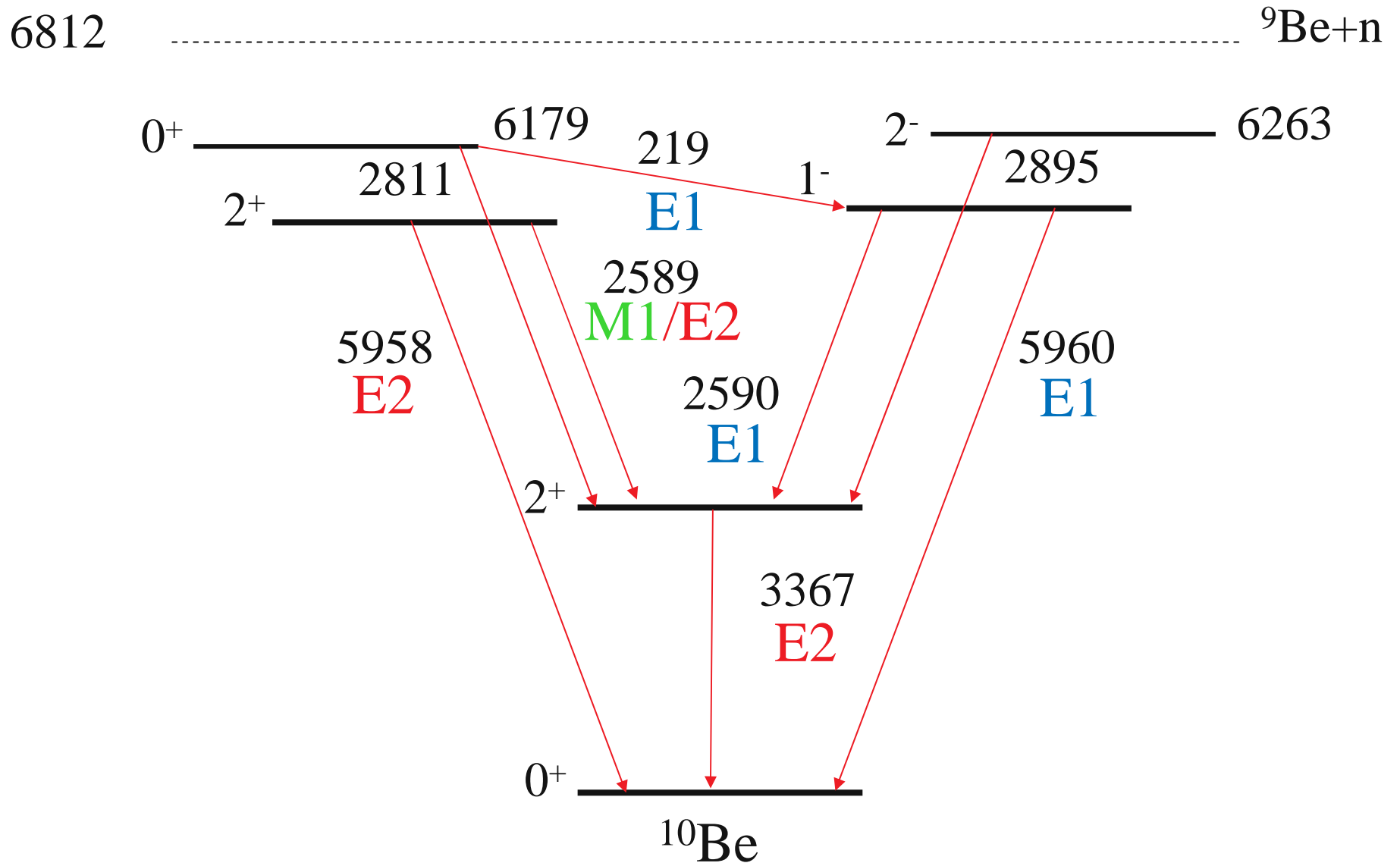
$$B(E2; 2_1^+ \rightarrow 0_1^+) = 7.89(30) e^2 \text{fm}^4$$

$$\tau = 217(\pm 7)_{stat} (\pm ??)_{sys}$$



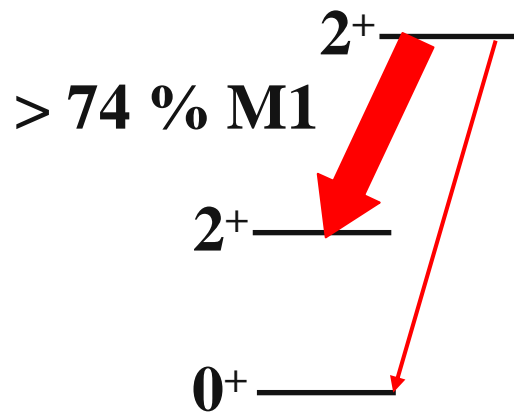
Theory seems pretty darn good at calculating E2 strengths

# The Bound States of $^{10}\text{Be}$





# Decay of the second $2^+$ state



$$|\delta| < 0.60$$

M.L. Roush *et al.*, Nucl. Phys. **A128**, 401 (1969).

$2_2^+ \rightarrow 2_1^+$	> 90	90(3)%
$2_2^+ \rightarrow 0_1^+$	< 10	10(3)%
	<b>NDS</b>	<b><math>^{11}\text{Li}</math>, <math>\beta\text{n}</math></b>

$$\Gamma_{tot} = \frac{1}{1 + \delta^2} (\delta^2 \Gamma_{E2} + \Gamma_{M1})_2 + (\Gamma_{E2})_0$$

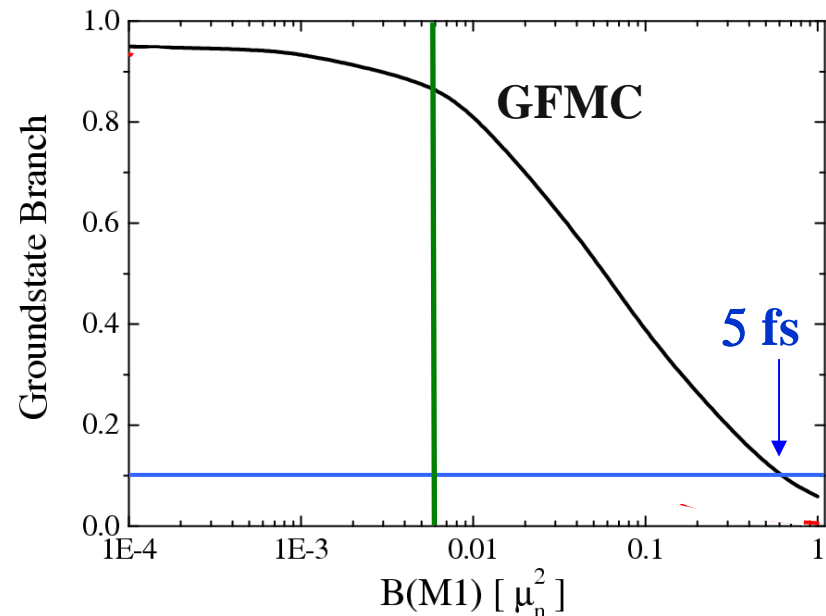
**GFMC**

$$B(E2; 2_2^+ \rightarrow 0_1^+) = 2.08(26) e^2 fm^4$$

$$B(E2; 2_2^+ \rightarrow 2_1^+) = 6.9(10) e^2 fm^4$$

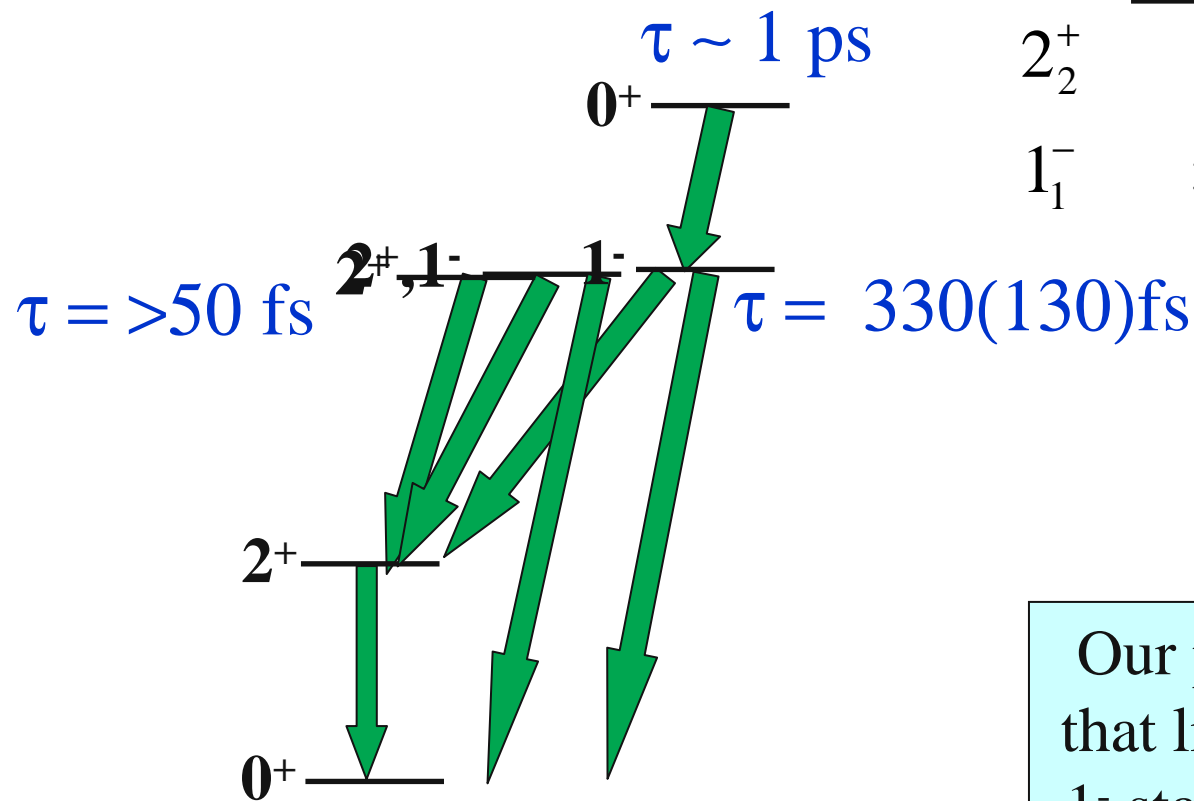
**VMC**

$$B(M1; 2_2^+ \rightarrow 2_1^+) \sim 0.006 \mu_n^2$$



# Lifetime of $2^+/1^-$ doublet

## $^{11}\text{Li}$ $\beta$ -delayed 1-n emission



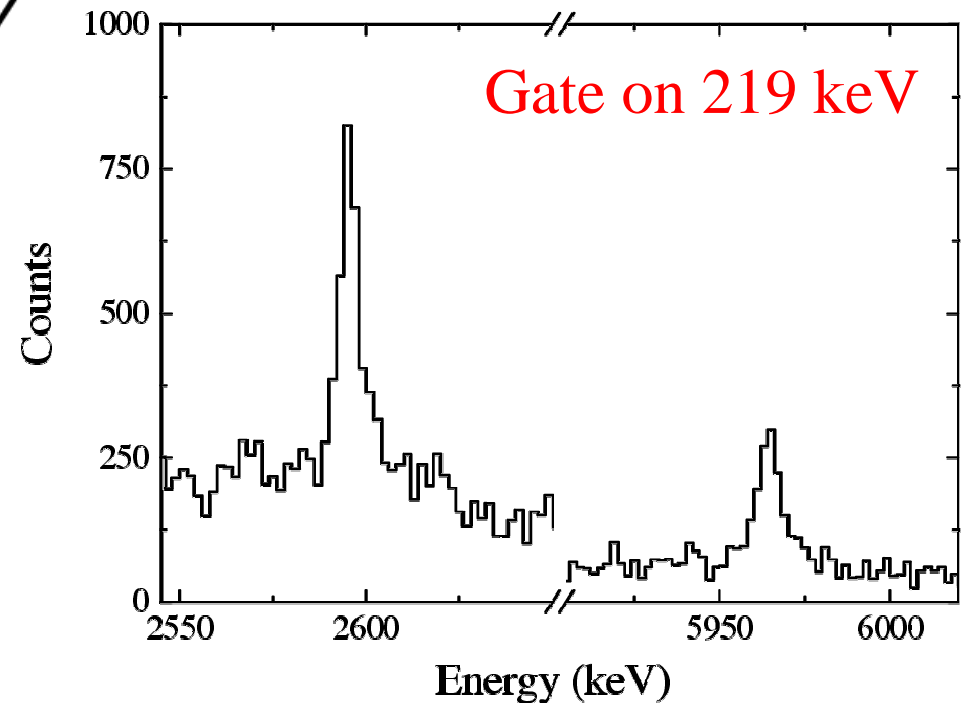
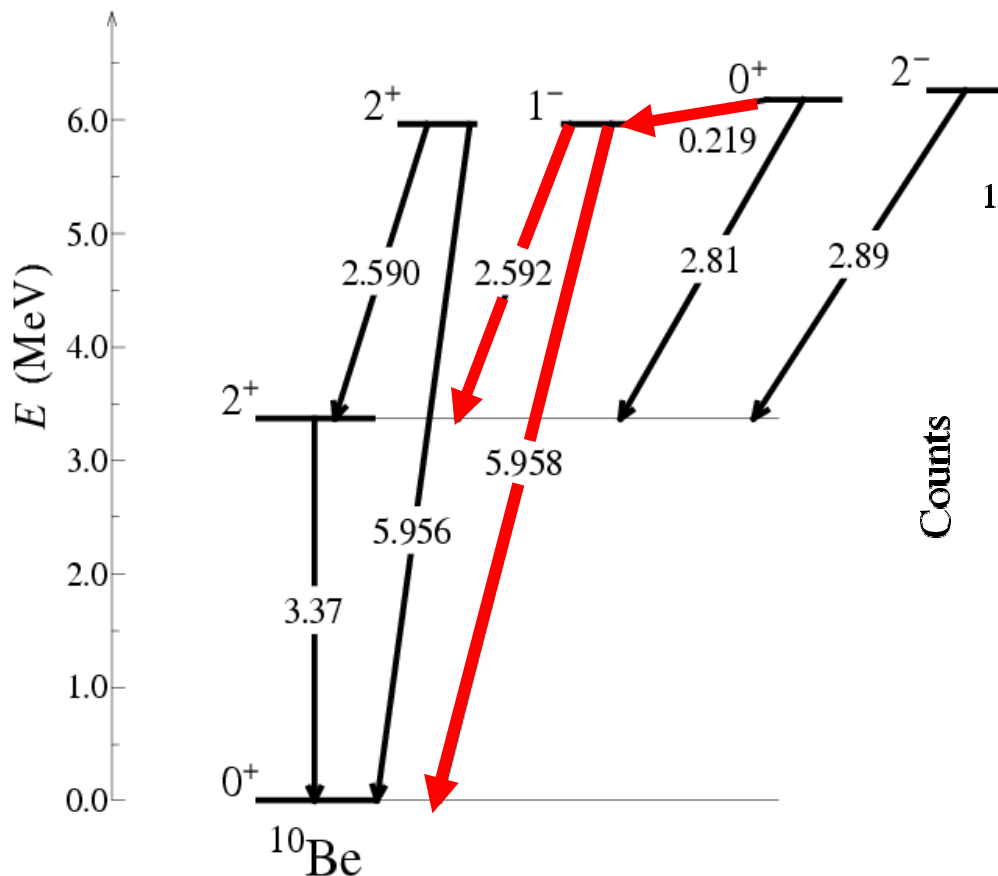
	ISOLDE	TRIUMPH
$2_2^+$	> 50 fs	86(11) fs
$1_1^-$	330(130) fs	< few 100 fs

Our preliminary indication is that lifetimes of both  $J=2^+$  and  $1^-$  state are significantly faster than the ISOLDE results

# Gamma-gamma analysis for branching ratio

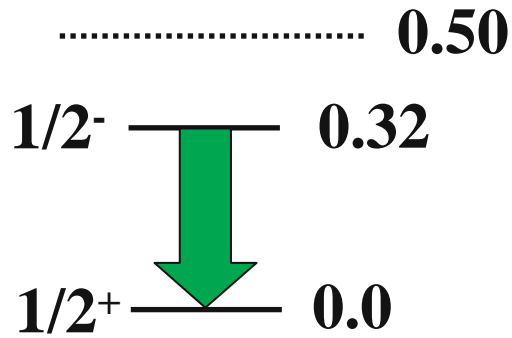
- Previous measurements fit doublet
- $\gamma$ - $\gamma$  isolates  $1^-$  branch

	$1^- \rightarrow 2_1^+$	$1^- \rightarrow 0_1^+$
NDS	$17_{-10}^{+6}$	$83_{-6}^{+10}$
$^{11}\text{Li } \beta\text{n}$	35(15)	65(15)
Present	41(4)	59(4)



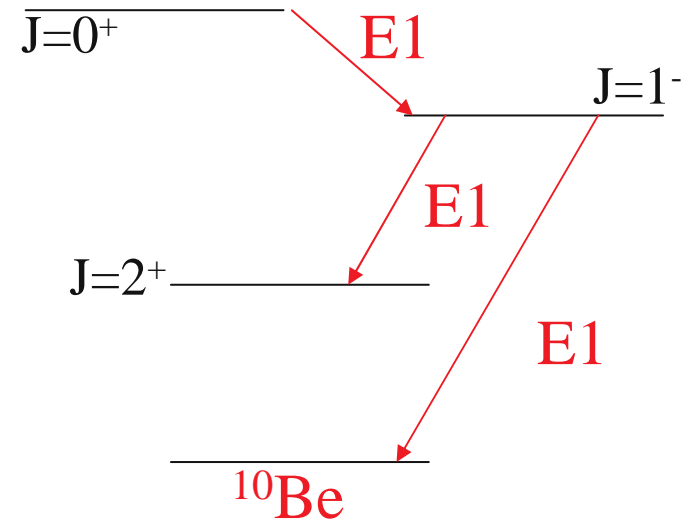
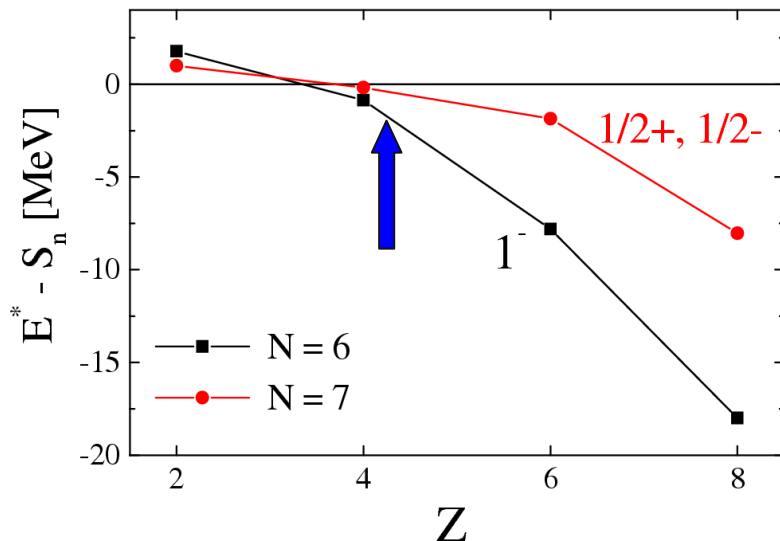
# Electric Dipole Physics.....hints of Halos

“Classic” halo nucleus  $^{11}\text{Be}$



$$B(E1) = 0.12(1) \text{ e}^2\text{fm}^2$$

D.J. Millener et al., Phys. Rev. C **28**, 497 (1983).



$^{10}\text{Be}$  B(E1) in  $\text{e}^2\text{fm}^2$

	$1^- \rightarrow 2_1^+$	$1^- \rightarrow 0_1^+$
GPMC	$2.4 \times 10^{-2}$	$7.4 \times 10^{-3}$
MO		$2.4 \times 10^{-3}$
Exp	$2(1) \times 10^{-5}$	$7(3) \times 10^{-6}$

MO : N. Itagaki and S. Okabe, Phys. Rev. C **61**, 044306 (2000).

Exp : H.O.U. Fynbo, Nucl. Phys. **A736**, 39 (2004).

# Conclusions

## In Experiment

Precision spectroscopy near stability and “discovery” studies far from stability are **BOTH** important and complimentary approaches.

We must make precise measurements (improve technique).  
Reliable <5% DSAM measurement **IS** possible.

## In Theory

We need to bridge what we learn from ab-initio calculations to more global models using NCSM, Hartree-Fock, Relativistic Mean Field and Density Functional approaches.

# Spares

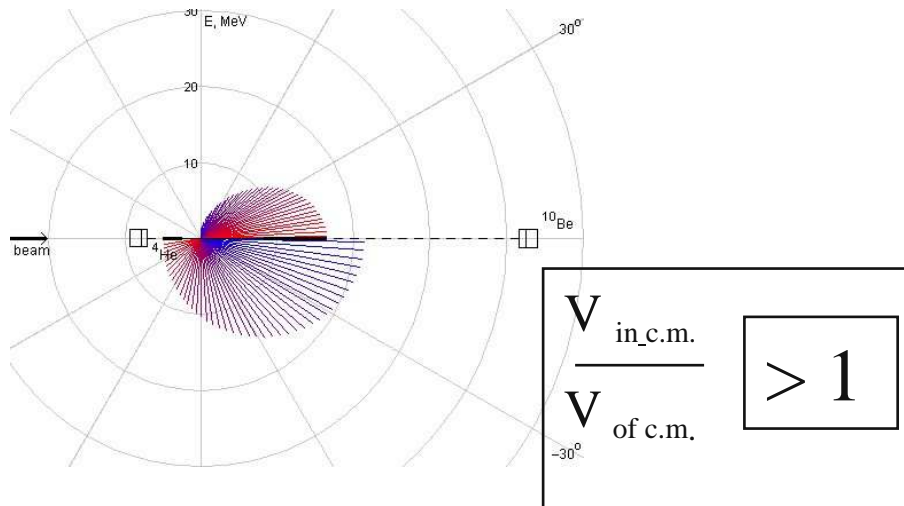
## Choice of reaction



$$Q = +14.784 \text{ MeV}$$

$$E_b = 10 \text{ MeV}$$

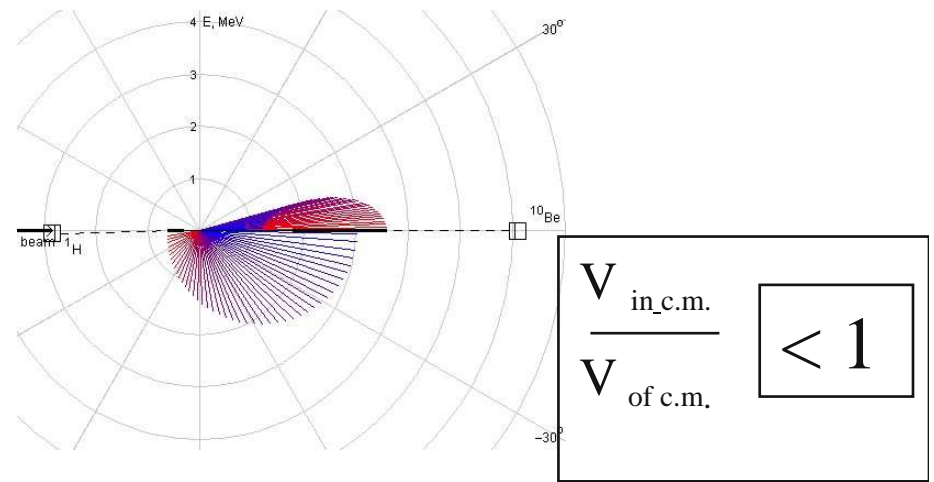
Very high recoil velocity  
 $(\beta_i = 6\%)$   
 evaporated target  
 good FMA separation  
 Low efficiency



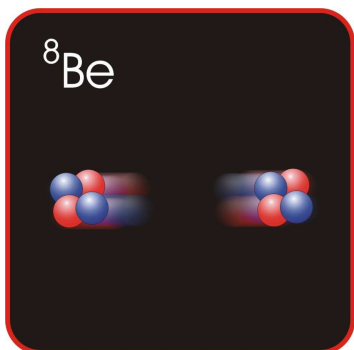
$$Q = +4.588 \text{ MeV}$$

$$E_b = 13.5 \text{ MeV}$$

Very high recoil velocity  
 $(\beta_i = 6\%)$   
 implanted target  
 difficult FMA separation  
 High efficiency

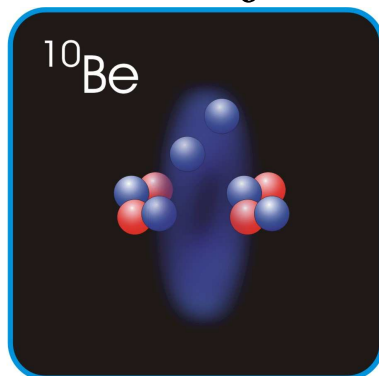


# Simple toy model



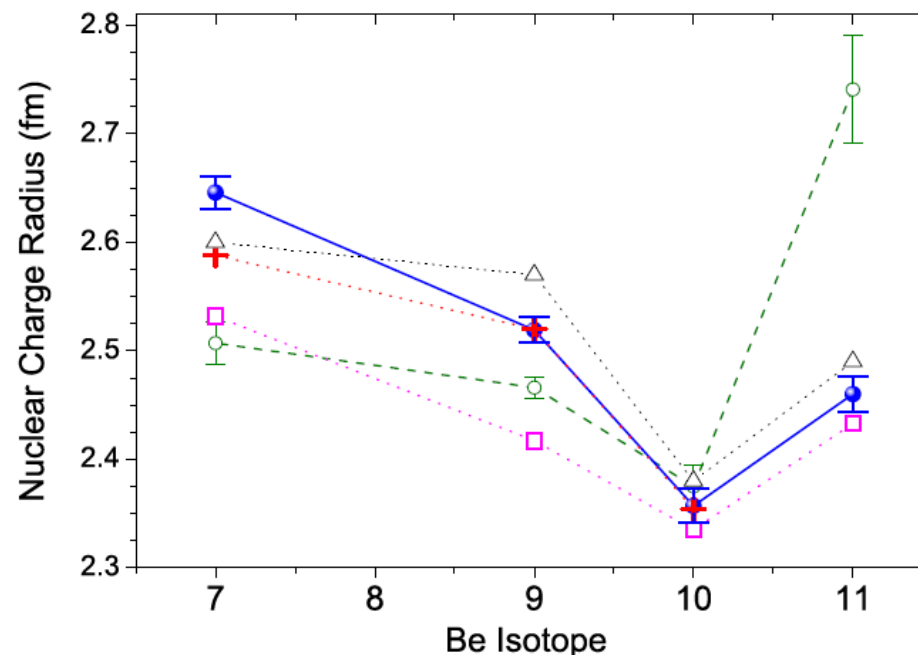
$$B(E2; 2_1^+ \rightarrow 0_1^+)_{cal} = 14.8 e^2 \text{ fm}^4$$

$$r = 2.6 \text{ fm}$$



$$B(E2; 2_1^+ \rightarrow 0_1^+) = 7.8(10) e^2 \text{ fm}^4$$

$$r = 2.2 \text{ fm}$$



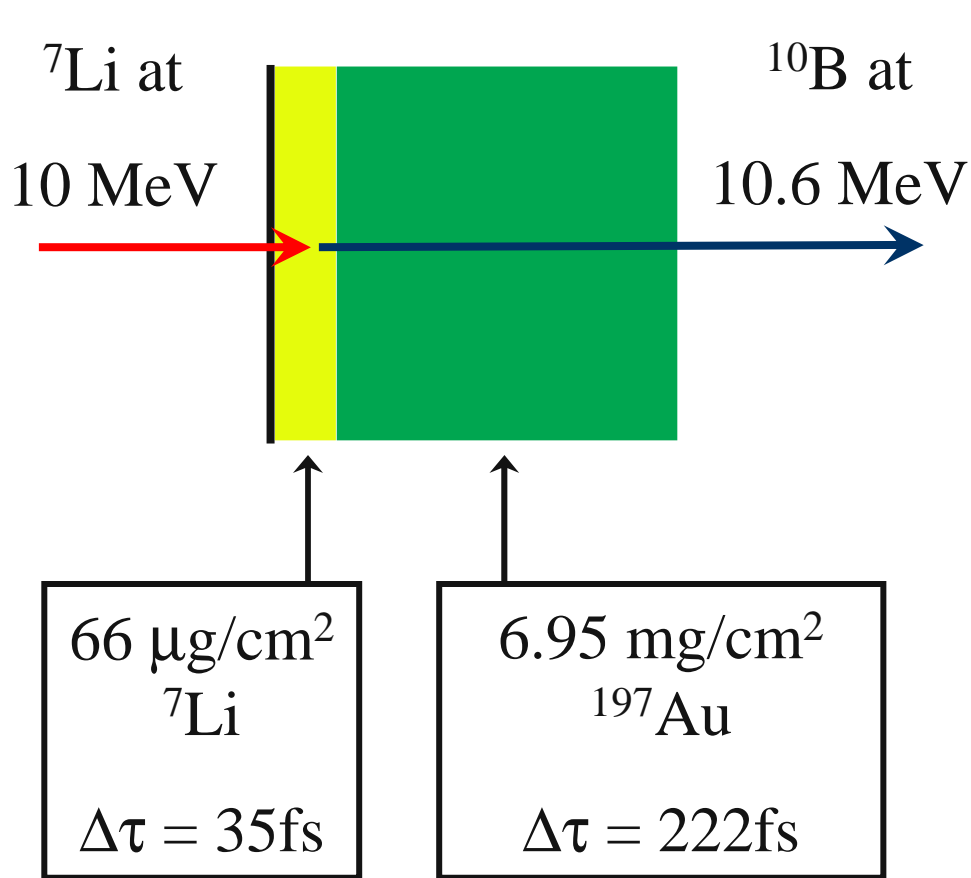
Nuclear Charge Radii of  $^{7,9,10}\text{Be}$  and the one-neutron halo nucleus  $^{11}\text{Be}$

W. Nörtershäuser,<sup>1,2</sup> D. Tiedemann,<sup>2</sup> M. Žáková,<sup>2</sup> Z. Andjelkovic,<sup>2</sup> K. Blaum,<sup>3</sup> M. L. Bissell,<sup>4</sup>  
 R. Cazan,<sup>2</sup> G.W.F. Drake,<sup>5</sup> Ch. Geppert,<sup>6,7</sup> M. Kowalska,<sup>8</sup> J. Krämer,<sup>2</sup> A. Krieger,<sup>2</sup> R. Neugart,<sup>2</sup>  
 R. Sánchez,<sup>6</sup> F. Schmidt-Kaler,<sup>9</sup> Z.-C. Yan,<sup>10</sup> D. T. Yordanov,<sup>3</sup> and C. Zimmermann<sup>7</sup>

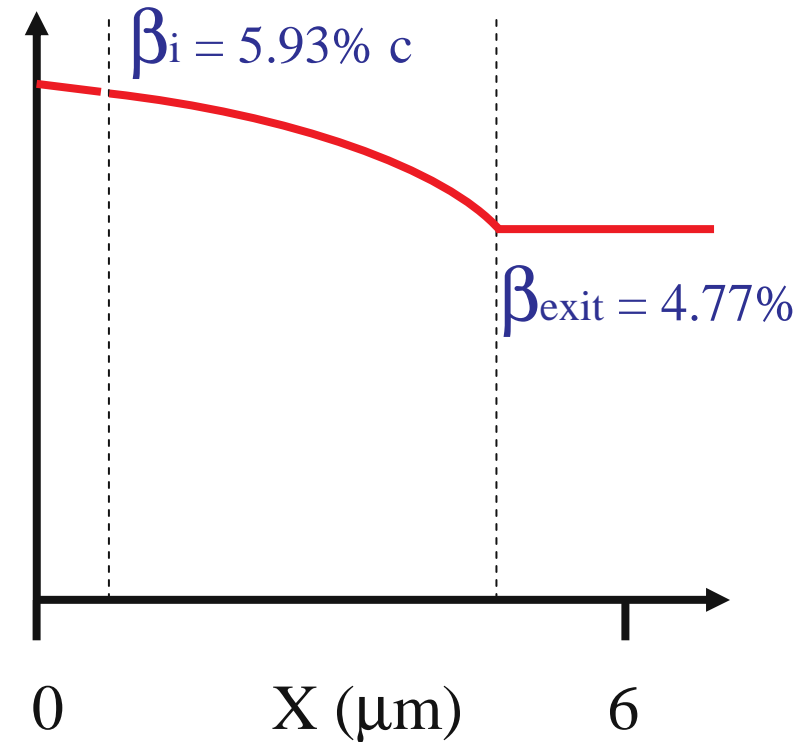
arXiv:0809.2607v1 [nucl-ex]



# Results (1): First excited State with $J^\pi=2^+$ at 3367 keV



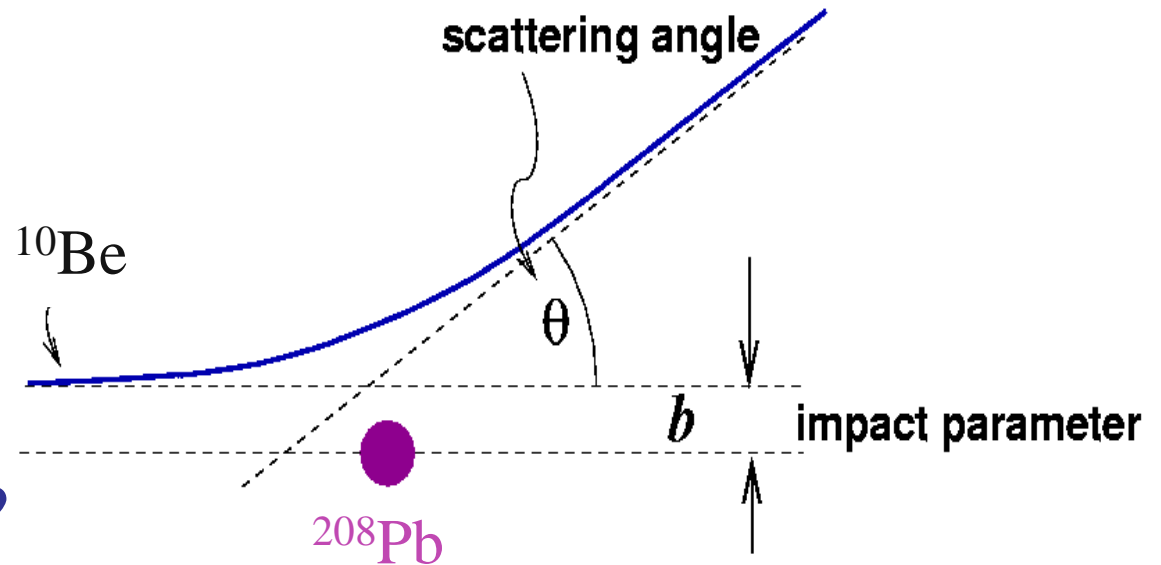
$Q = 11.418 \text{ MeV}$



~17% of  $\gamma$ -rays in Li  
~60% of  $\gamma$ -rays in Au  
~23% of  $\gamma$ -rays in vacuum

# An alternate approach: $^{10}\text{Be}$ Coulomb Excitation

$^{10}\text{Be}$  was produced in a reactor using  $^9\text{Be}(n,\gamma)^{10}\text{Be}$  and mass separated at ORNL. It is now being prepared as a beam, with expected intensity  $>10^9$ pps, as it has  $T_{1/2}=1.5\text{Myr}$ .



Is this a better approach?

Is it more precise?

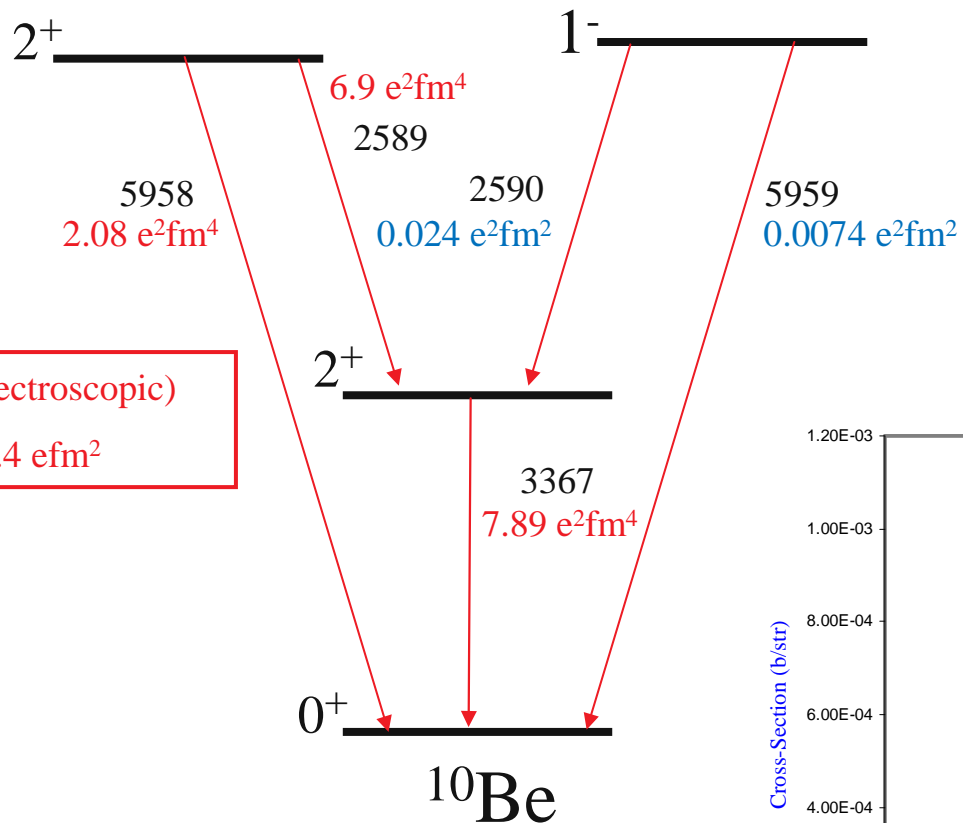
Can one investigate the upper  $5960\text{keV } J^\pi=2^+, 1^-$  doublet?

Can we cross-check our DSAM results?

## A QUESTION OF NUMBERS

# Coulomb Excitation: 45 MeV $^{10}\text{Be}$ on $^{206}\text{Pb}$ target

$Q_{22}$  (spectroscopic)  
+3.7 efm<sup>2</sup>



$Q_{22}$  (spectroscopic)  
-5.4 efm<sup>2</sup>

Ab-initio matrix elements (Wiringa / Pieper) into DeBoer / Winter inelastic excitation code (Esbensen).

Tough experiment:

Cross-sections are **VERY** low

Hard to get statistics to beat DSAM...but can confirm.

IF the beam  $>10^9$  pps **AND** the B(E1) are large then the upper doublet might be investigated.

